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Is the "Conservative" frequentist upper limit always worse than the "Optimistic" Bayesian upper limit ?
No-A Simple Counterexample

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1 Model

Take a one dimensional Gaussian distribution for $X \ge 0$ with a parameter $A \ge 0$.

$$P(X|A) = N \exp^{-(X-A)^2} \tag{1}$$

$$N^{-1} = \frac{\sqrt{\pi}}{2} (1 + \text{Erf}(A)) \tag{2}$$

Suppose a relation A = HD where D is not known but we know D is bounded: $0 < W \le D \le B < \infty$. W (worst) and B (best) denote arbitrary finite real constants. (In this model, A is akin to gravitational-wave strain at the detector, projected onto the detector's antenna pattern, while H is akin to the source's intrinsic strength h_0 , and D is the uncertain, angle-dependent factor that relates them.) Now suppose we want to set an upper limit on H, given a data x.

Given a significance level p (and thus the confidence level 1-p), the frequentist upper limit on A, which we denote by $\hat{A}_f = \hat{A}_f(x, p)$, is given implicitly by

$$p = \int_0^x P(X|\hat{A}_f) dX = (\text{Erf}(\hat{A}_f) - \text{Erf}(\hat{A}_f - x)) / (1 + \text{Erf}(\hat{A}_f)).$$
 (3)

The conservative frequentist upper limit on H, which we call \hat{H}_{fw} , is then $\hat{H}_{fw} = \hat{A}_f/W$.

Given the uniform priori probability $P(A) = \text{const } (0 \le A < \infty)$ the Bayesian upper limit on A denoted by $\hat{A}_b = \hat{A}_b(x, p)$ is given by A

$$1 - p = \int_0^{\hat{A}_b} P(A|x) \, dA = N' \int_0^{\hat{A}_b} P(x|A) P(A) \, dA =$$

$$= N'' \int_0^{\hat{A}_b} P(x|A) \, dA = N'' N \int_0^{\hat{A}_b} \exp(-(x-A)^2) \, dA =$$

$$= (\text{Erf}(\hat{A}_b - x) + \text{Erf}(x)) / (1 + \text{Erf}(x)). \tag{4}$$

where we used

$$(N''N)^{-1} = \frac{\sqrt{\pi}}{2}(1 + \text{Erf}(x)).$$
 (5)

The optimistic Bayesian upper limit on H, which we call \hat{H}_{bo} , is then just given by $\hat{H}_{bo} = \hat{A}_b/B$.

2 Two examples

Suppose (W,B) = (0.5,1) and we take p = 0.05.

¹This is not quite the procedure we followed in the Bayesian pulsar analysis. There we assumed the uniform priori probabilities for $h_0 > 0$, for the cosine of the inclination angle $\cos \iota$ that is the angle between the angular momentum of the pulsar and the line of the sight, the gravitational-wave polarization angle ψ , and the gravitational-wave initial phase Φ_0 . Then we marginalised the probability with respect to $\cos \iota$, ψ , and Φ_0 to set an upper limit on h_0 . Analogously, in this toy model, one may assume P(H) = const. (0 ≤ $H < \infty$), P(D) = const. ($W \le D \le B$), and then marginalise with respect to D to obtain the marginalised probability for H. Here, for simplicity, we do not marginalise the probability, but simply set D=B to yield the most optimistic limit as stated below by Eq. (5).

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Suppose x = 1.4, then
$$(\hat{A}_f,\hat{A}_b)=(2.6,2.6)$$
 and correspondingly, $(\hat{H}_{fw},\hat{H}_{bo})=(5.2,2.6)$.

case 2.

Suppose x = 0.1, then
$$(\hat{A}_f, \hat{A}_b) = (0.62, 1.5)$$
 and correspondingly, $(\hat{H}_{fw}, \hat{H}_{bo}) = (1.24, 1.5)$.

Obviously, $\hat{H}_{fw} < \hat{H}_{bo}$ in case 2. Note that the ratio $\hat{H}_{fw}/\hat{H}_{bo}$ depends on the data, not just on the methods.

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