Scaling law in signal-recycled interferometers

Alessandra Buonanno 1,2 and Yanbei \mbox{Chen}^2

¹ Institut d'Astrophysique de Paris (CNRS)
² California Institute of Technology

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Signal-recycled interferometer of Advanced LIGO



The Presence of a Signal-Recycling Mirror (SRM)

[At all optical powers]

• Modifies optical resonant frequency and storage time. [Drever, 82; Meers, 88; Mizuno, 95.]

[At high enough optical power, e.g., in LIGO-II]

• Creates optomechanical coupling that changes test-mass dynamics. [Buonanno and Chen 01–02]

• Allows the interferometer to beat the SQL [Buonanno and Chen 01-02]

Exploring the parameter space: *scaling law?* [Buonanno and Chen 01–02, Fritschel, Shoemaker, Strain ...]

1×10⁻²

1×10⁻²²



Example given by Buonanno and Chen:

 $I_{\rm BS} = 10 \,\mathrm{kW}$ T = 0.033, $\gamma_{\rm arm}$ = 2π \times 100 Hz $ho = 0.9, \, \phi = \pi/2 - 0.47$



Reference design by LIGO experimentalists:

 $I_{\rm BS} = 2.5 \,\mathrm{kW}$ T = 0.005, $\gamma_{\rm arm} = 2\pi \times 15 \, {\rm Hz}$ $\rho = 0.964, \, \phi = \pi/2 - 0.06$



The use of the characteristic description [Buonanno and Chen, 02]

• Previous results of SR interferometers at high powers [Buonanno and Chen, 00–01] written in characteristic parameters: equations much simpler, more physical. For example, the *optical spring* constant:

$$\frac{K_{\rm opt}(\Omega)}{m/4} = \frac{-\lambda\iota_c}{(\Omega - \lambda + i\epsilon)(\Omega + \lambda + i\epsilon)} \qquad \dots = \Omega^2$$

See Buonanno and Chen, gr-qc/0208048 for more nice results.

- Scaling laws:
 - Optical, confirmed that, with I_c , m, L fixed, scaling laws in the low power regime still valid here: $(\lambda, \epsilon) = \text{const.}$ Can be used for LIGO-II optimization.
 - Opto-Mechanical, one more scaling relation: $\iota_c \equiv \frac{8\Omega_0 Ic}{mLc} = \text{const.}$ Relating LIGO-II to experiments in other regimes.

Using the Optical scaling law: simplifying LIGO-II optimization

For configurations with the same I_c , m and L, fix λ and ϵ , and:



A: $(\lambda, \epsilon) = 2\pi \times (195 \text{ Hz}, 25 \text{ Hz})$, contains $(T, \rho, \phi) = (0.033, 0.9, \pi/2 - 0.47)$. [Buonanno and Chen, 01–02]

B: $(\lambda, \epsilon) = 2\pi \times (228 \text{ Hz}, 69 \text{ Hz})$, contains $(T, \rho, \phi) = (0.005, 0.964, \pi/2 - 0.06)$. [LIGO-II Reference Design.]

Using the full Optomechanical scaling law

Linking LIGO-II and table-top experiments

Table-top experiment proposed by Braginsky, Khalili and Volikov (2001), single detuned cavity with one movable mirror. [Directly analyzed by Khalili, (2001)]

. . . In a very different regime from LIGO



now fits into the same formalism as LIGO-II:

$$\begin{array}{c|c} \epsilon \sim \lambda / \sqrt{3} & \iota_c^{1/3} & \Omega \\ \hline 5 \times 10^7 \, \mathrm{s}^{-1} & 10^5 \, \mathrm{s}^{-1} & 10^4 \, \mathrm{s}^{-1} \end{array}$$

Approximations made in the quantum optical analysis of LIGO interferometers

So far: leading order in T and $\Omega L/c$.

Errors associated:

• Difference in the ITM and ETM contributions to radiation-pressure effects ignored. *Still to be analyzed.*



• Propagations of optical fields less accurate. *Resolved by this work*.

Leading-order in T vs all orders in T



 $T = 0.033, \ \rho = 0.9, \ \phi = \pi/2 - 0.47$ $T = 0.005, \ \rho = 0.964, \ \phi = \pi/2 - 0.06$

Summary

- SR interferometer *with high laser power* (with very short SR cavity) mapped into a single detuned cavity
- Scaling law among experimental parameters
- Scaling law helps:
 - presenting calculation results in a simpler way, more physical (!)
 - optimization for LIGO-II
 - relating LIGO-II to table-top experiments
- Remaining issue: difference in the contributions of ITM and ETM to raditaion-pressure effects.