

# Quantum Non-Demolition in Advanced Gravitational Wave Interferometers

David Ottaway  
LIGO MIT  
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# People Involved

Nergis Mavalvala, David Ottaway, Keisuke Goda,  
Thomas Corbitt (MIT)

Stan Whitcomb  
(Caltech)

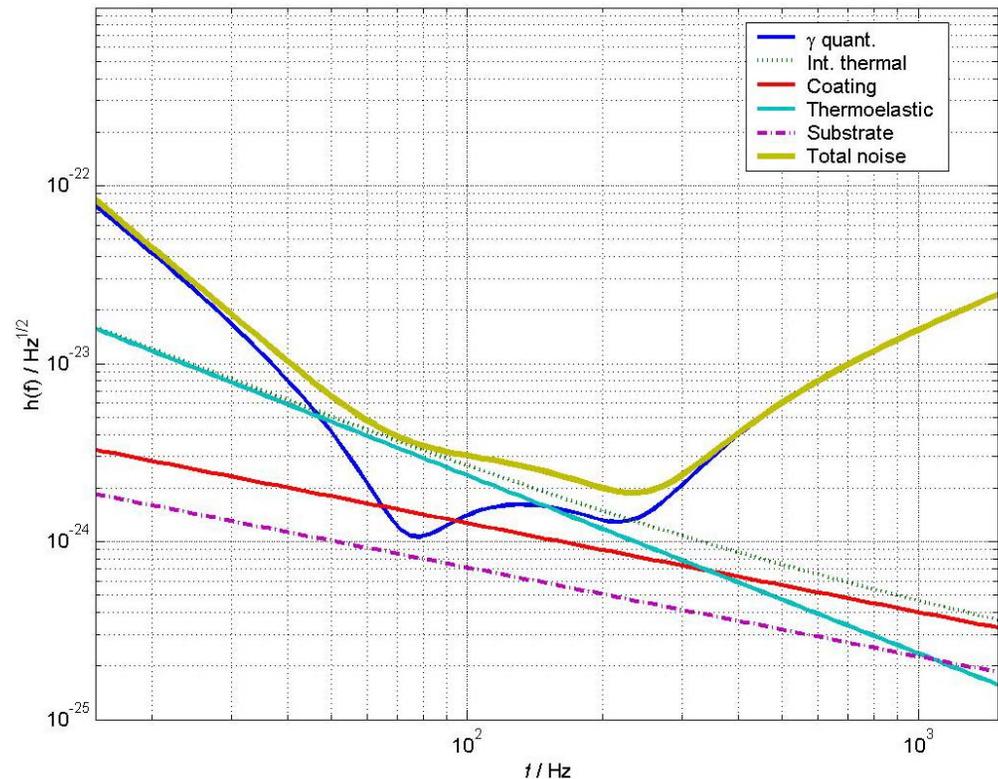
Collaborators: Pin Koy Lamb, David McClelland  
(ANU)

# Overview of Talk

- Advanced LIGO and the implications of quantum noise
- The Standard Quantum Limit (SQL)
- A discussion of quantum noise
- Ideas for beating quantum noise
- Current limits on QND performance
- Ideas and experiments to explore and beat the standard quantum limit with bench top experiments
- The potential gain for Advanced LIGO

# Advanced LIGO Sensitivity

- Predicted sensitivity for Sapphire interferometer
- Predominantly limited by thermoelastic noise and quantum noise
- Assumes a solution to coating thermal noise is found
- Low frequency optical noise  $\Rightarrow$  Radiation pressure noise
- High frequency optical noise  $\Rightarrow$  Shot Noise



# Quantum Noise

- Measurement process
  - Interaction of light with test masses
  - Counting signal photons with a PD
- Noise in measurement process
  - Poissonian statistics of force on test mass due to photons → radiation pressure noise (RPN)
  - Poissonian statistics of counting the photons → shot noise (SN)

# Strain sensitivity limit due to quantum noise

- Shot Noise

- Uncertainty in number of photons detected  $\Rightarrow$

$$h_{shot}(f) = \frac{1}{L} \sqrt{\frac{hc\lambda}{8F^2 P_{bs}}} \frac{1}{T_{ifo}(\tau_s, f)}$$

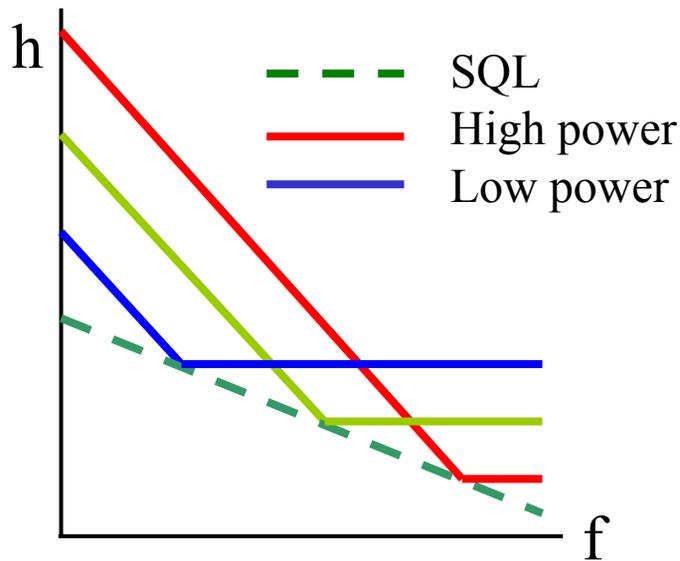
- (Tunable) interferometer response  $\rightarrow T_{ifo}$  depends on light storage time of GW signal in the interferometer

- Radiation Pressure Noise

- Photons impart momentum to cavity mirrors
  - Fluctuations in the number of photons  $\Rightarrow$

$$h_{RP}(f) = \frac{2F}{ML} \sqrt{\frac{2\hbar P_{bs}}{\pi^3 c \lambda}} \frac{T_{ifo}(\tau_s, f)}{f^2}$$

# Standard Quantum Limit (SQL)



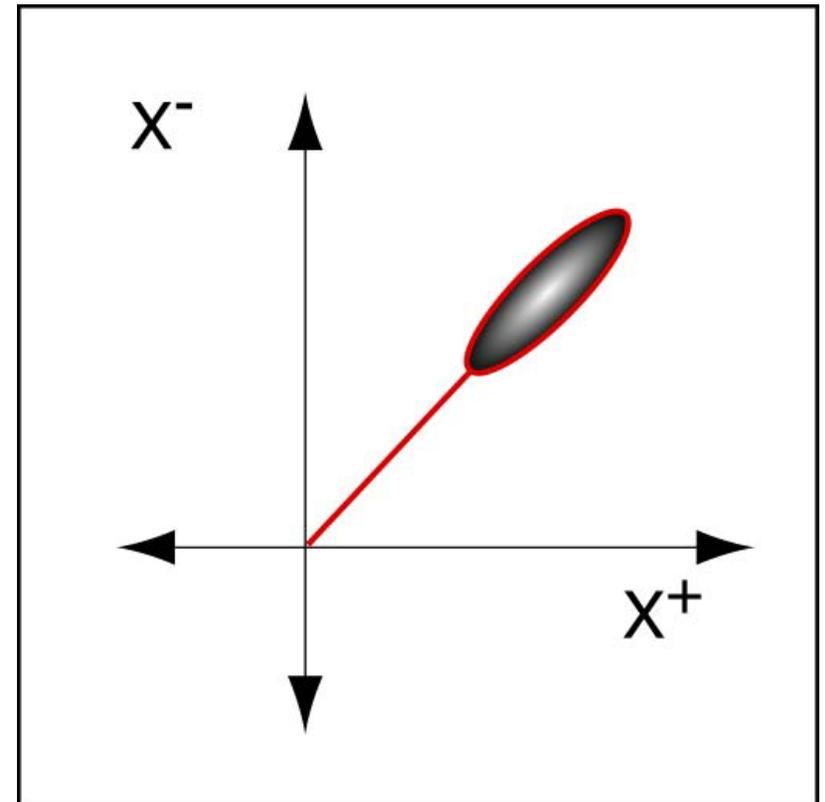
$$h_{SQL} = \frac{1}{L} \sqrt{\frac{h}{2M\pi^{3/2}}} \frac{T_{ifo}(\tau_s, f)}{f}$$

$$P_{bs} = \frac{c\lambda\pi^{3/2}}{8F^2} \frac{f^2}{T_{ifo}(\tau_s, f)^2}$$

- Minimum noise achieved when  $h_{\text{shot}} = h_{\text{rp}}$  for a given frequency  $f$
- Assumes that radiation pressure noise and shot noise add in quadrature
- No correlations and hence no noise cancellations considered

# 'Ball on stick' representation

- Analogous to the phasor diagram
- Stick  $\rightarrow$  dc term
- Ball  $\rightarrow$  fluctuations
- Common states
  - Coherent state
  - Vacuum state
  - Amplitude squeezed state
  - Phase squeezed state



# How to increase sensitivity beyond Advanced LIGO ??

Two options for increasing the sensitivity:

Classical approach: Get a bigger hammer !!

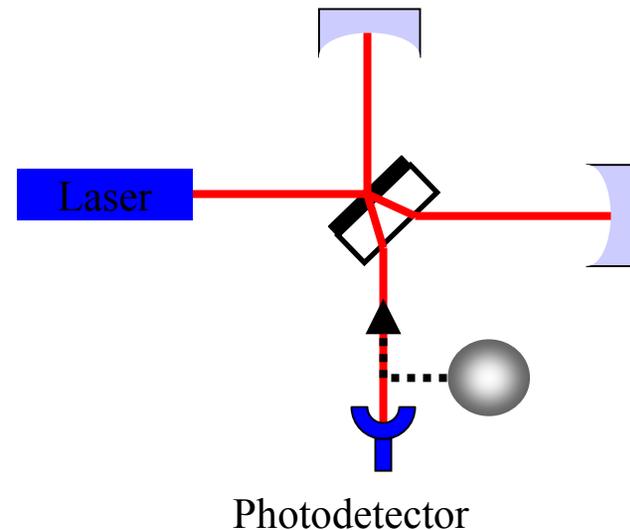
- Increase the mass of the test mass and increase the laser power

QND approach:

- Plug open output port of the beam splitter with squeezed vacuum
  - => Improved high frequency noise without power increase
- Cancel out the radiation pressure fluctuations effects on TM
- Combination of both

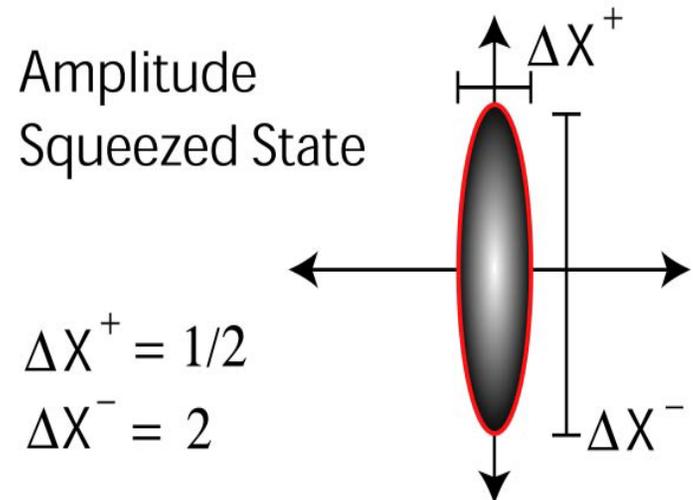
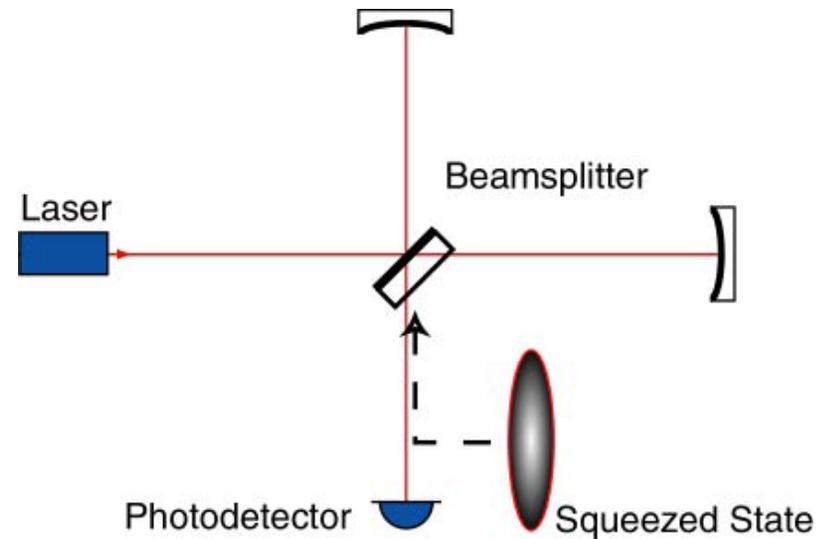
# Vacuum State in a Michelson

- Michelson on dark fringe  
→ laser contributes only correlated noise
- GW signal from anti-correlated effect  
→ phase quadrature
- Vacuum noise couples in at the beamsplitter
  - Phase quadrature fluctuations: **Shot Noise**
  - Amplitude quadrature fluctuations: **Radiation Pressure Noise**



# Squeezed State in a Michelson

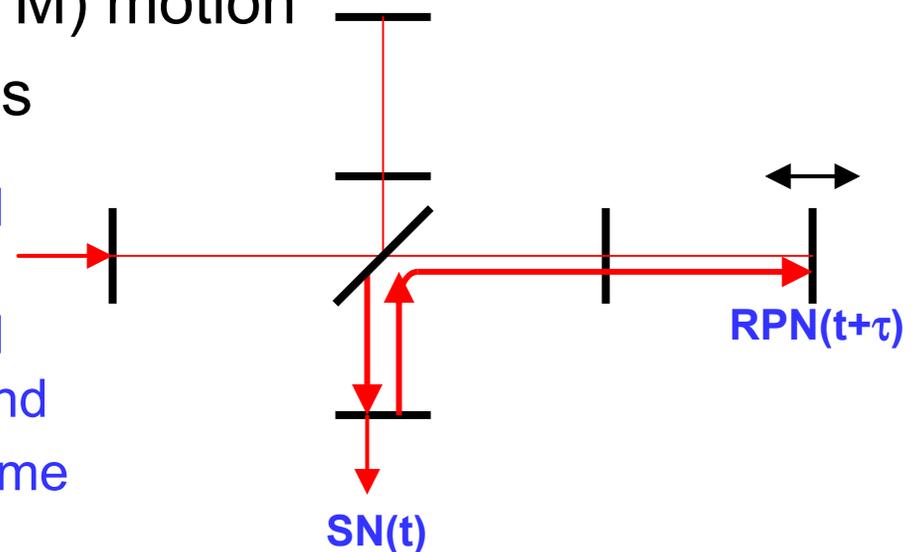
- Inject squeezed state into the dark port of Michelson to replace vacuum
- Amplitude squeezed state oriented to reduce noise in the signal (phase) quadrature
- GW signal in the phase quadrature



# Signal recycling mirror → quantum correlations

- Shot noise and radiation pressure (back action) noise are **correlated** (Buonanno and Chen, PRD 2001)
  - Optical field (which was carrying mirror displacement information) returns to the arm cavity
  - Radiation pressure (back action) force depends on history of test mass (TM) motion
  - Dynamical correlations

Part of the light leaks out the SRM and contributes to the shot noise  
**BUT** the (correlated) part reflected from the SRM returns to the TM and contributes to the RPN at a later time

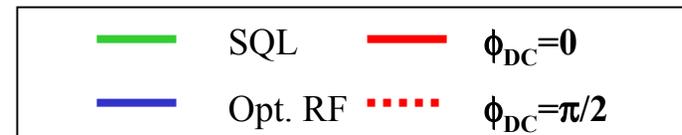
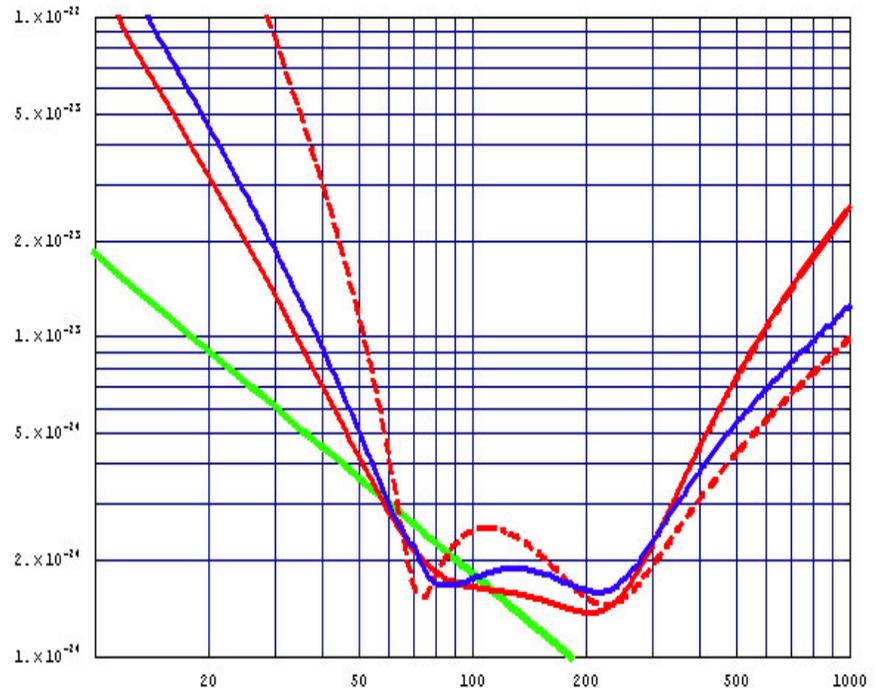


# New quantum limits

- Quantum correlations →
  - SQL applies to free particles
  - Test masses are connected by radiation pressure
  - Optomechanical resonance (“optical spring”)
  - Noise cancellations possible
- Quantization of TM position **not** important  
(Pace, et. al, 1993 and Braginsky, et. al, 2001)
  - GW detector measures position *changes* due to classical forces acting on TM
  - No information on quantized TM position extracted

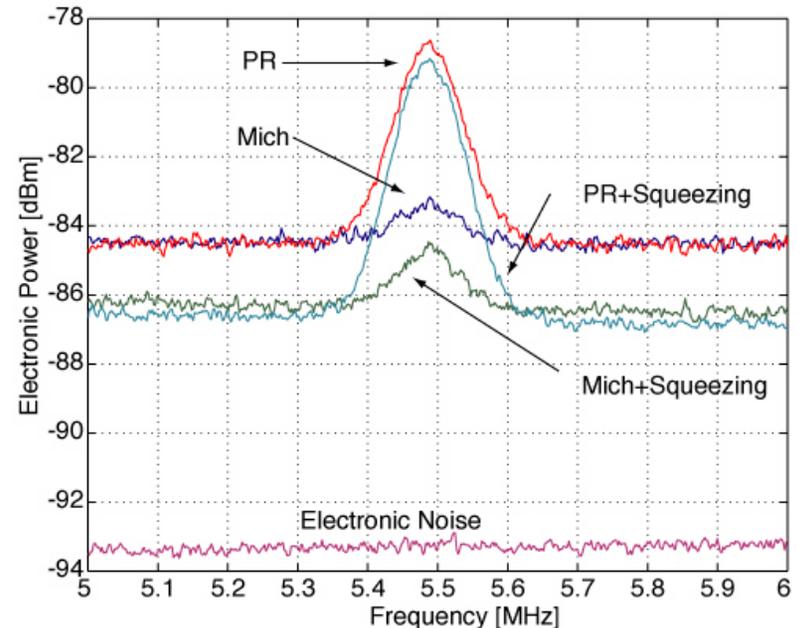
# Quantum Manipulation: AdLIGO as an example

- “Control” the quantum noise
- Many knobs to turn: Optimize ifo sensitivity with
  - Choice of homodyne (DC) vs. heterodyne (RF) readout
  - RSE detuning  $\rightarrow$  reject noise one of the SB frequencies
  - Non-traditional modulation functions
  - Non-classical light?



# Quantum manipulation: Avenues for AdLIGO+

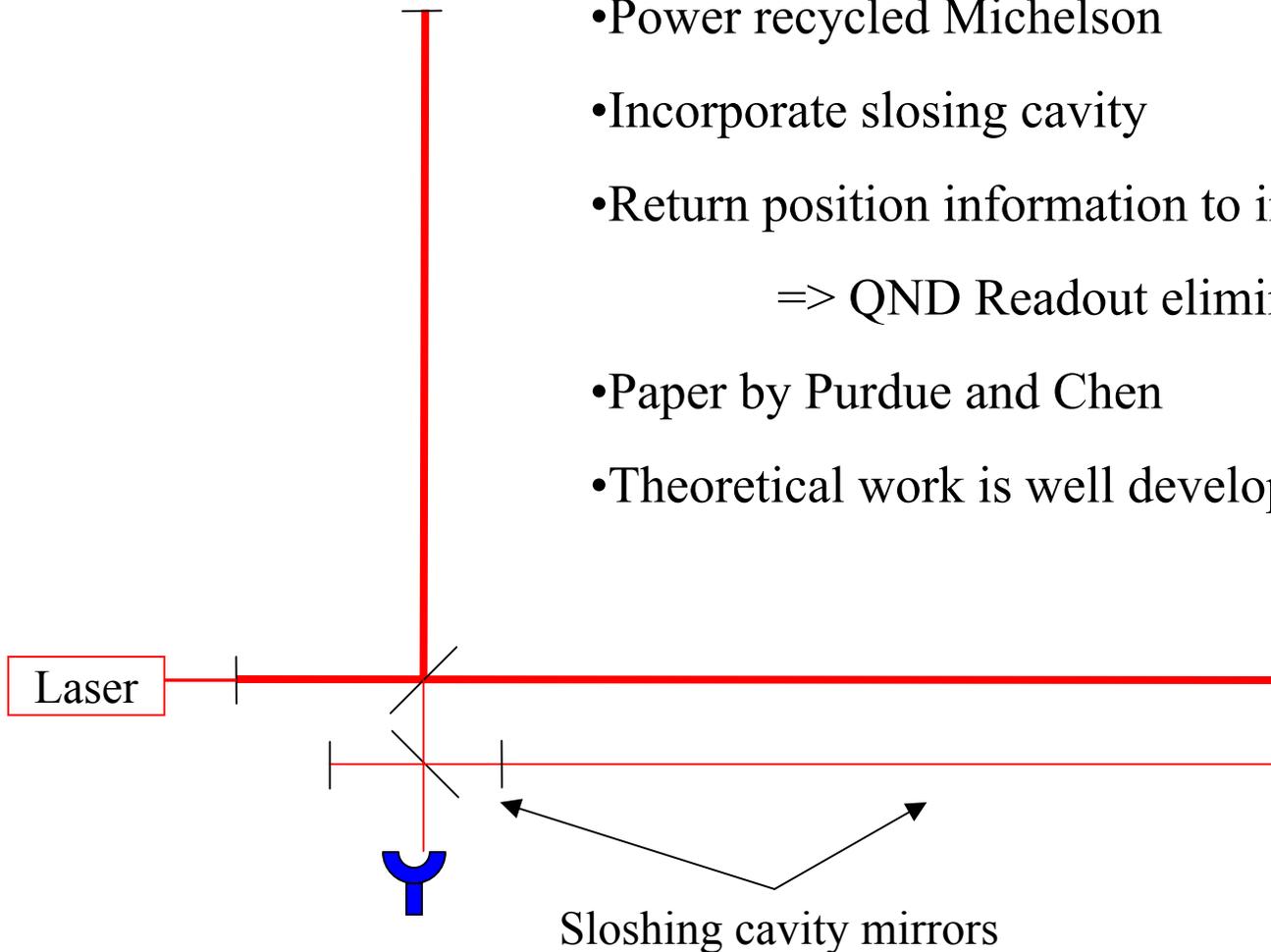
- Squeezed light
  - Increased squeeze efficiency
    - Non-linear susceptibilities
    - High pump powers
    - Internal losses
    - Low (GW) frequencies
- QND readouts
  - Manipulation of sign of SN-RPN correlation terms
  - Manipulation of signal vs. noise quadratures (KLMTV, 2000)
  - Squeezed vacuum into output port



ANU, 2002

# Speed Meter Example

- Power recycled Michelson
- Incorporate sloshing cavity
- Return position information to interferometer  
=> QND Readout eliminates back action
- Paper by Purdue and Chen
- Theoretical work is well developed



# Squeezed source applications:

## ▪ Advance LIGO

- Signal tuned
- Quantum noise limited
- SEM  $\rightarrow$  SN and RPN correlated  $\rightarrow$  QND
- Squeezed input
- Evaluation in the presence of other noise limits and technical noise

## ▪ AdLIGO +

- Speedmeters
- White light interferometers
- All reflective interferometers
- Ponderomotive squeezing

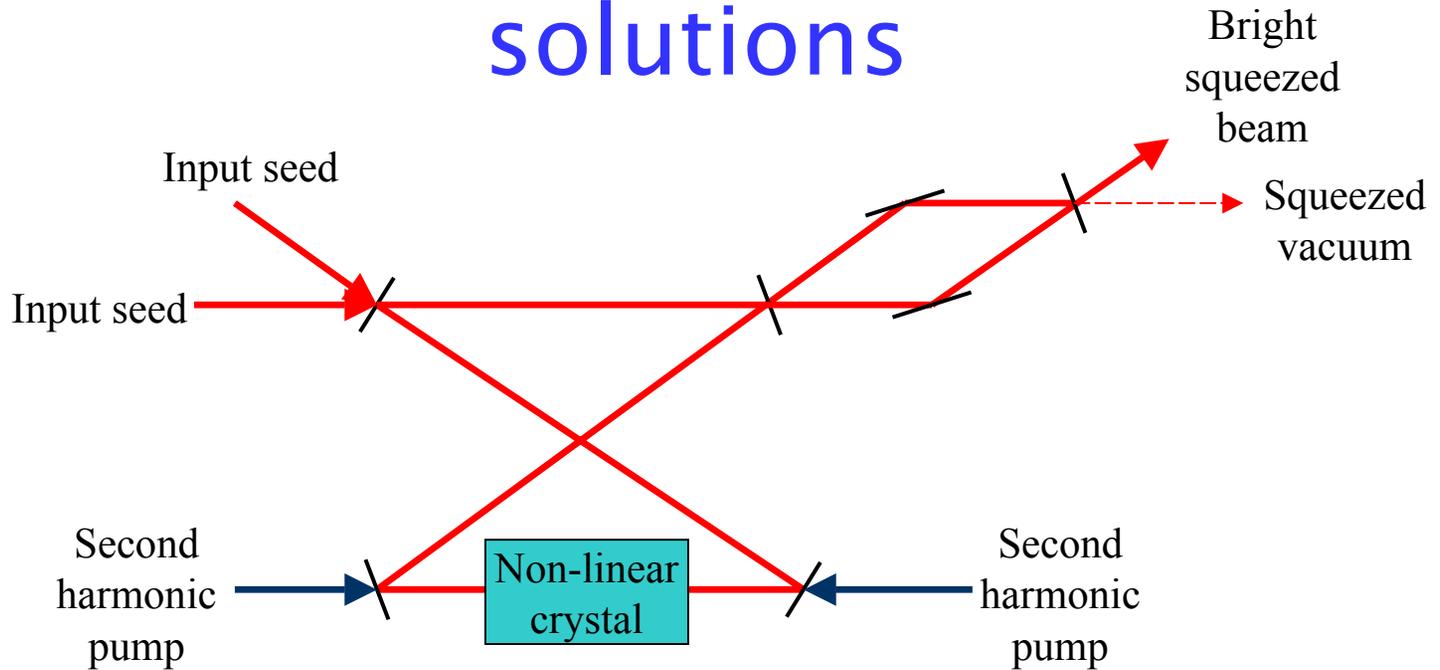
# Current issues in Squeezing

- Current squeezing is limited to around 7dB
- Current squeezing is limited to above the near MHz region
- Squeezing required at the level of 10 dB at 100 Hz
- How ????

## 10 dB Possible Solutions

- Internal losses and interferometer injection losses
- Larger  $\kappa^2$  and  $\kappa^3$  non-linearities
- Increased pump power
- Improved photo-detection efficiency (different wavelength)
- 10 dB Consortium to obtain improved crystals

# Low frequency squeezing solutions



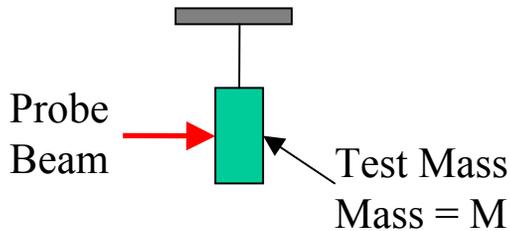
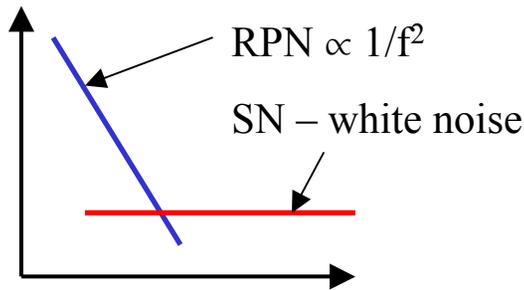
- Uses two squeezers to cancel the technical noise on the seeds – shown to produce squeezing to 200 kHz by Bowen *et al.*
- Common OPO to cancel OPO technical noise
- Proposed by Bowen *et al.*



# Bench top sub-SQL experiments

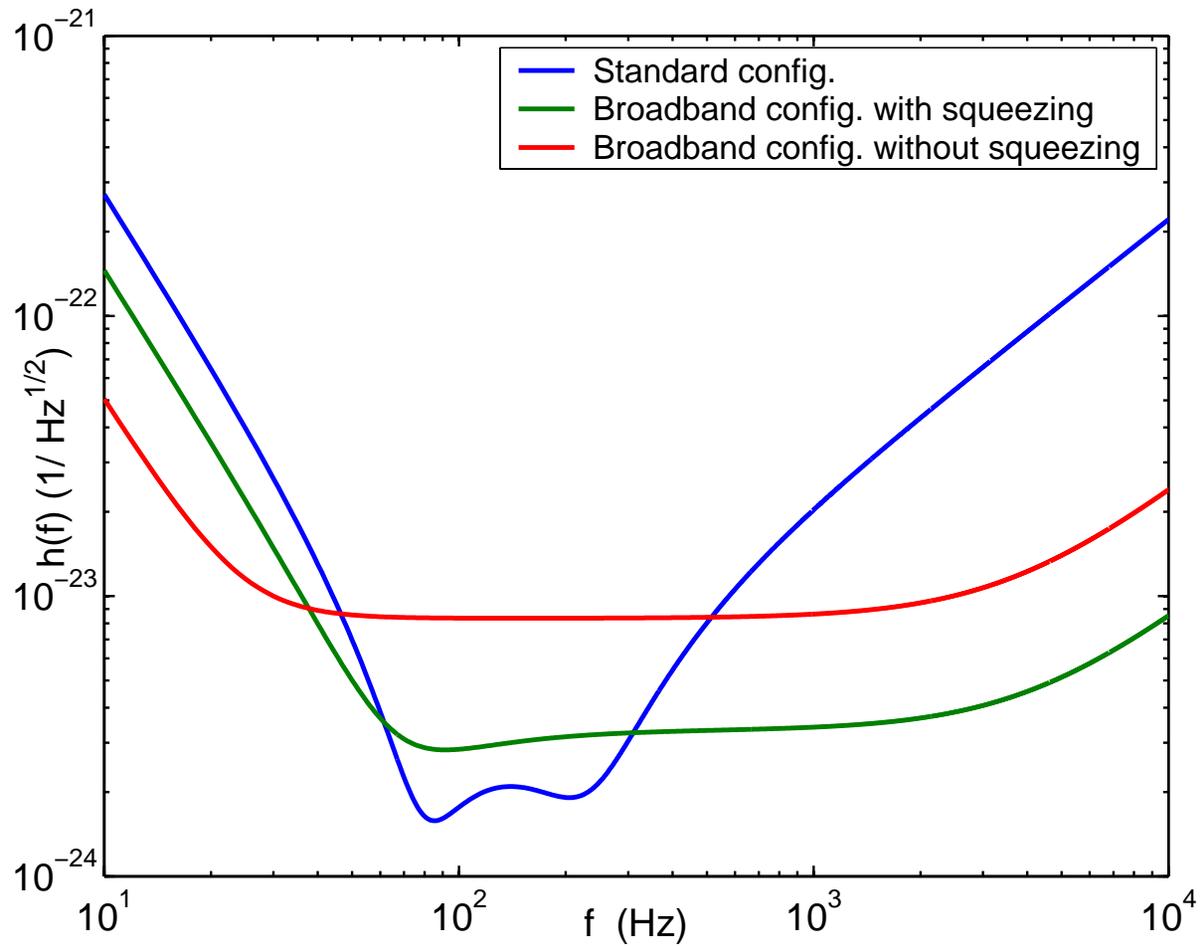
▪ Remember:

$$x_{shot}(f) = \frac{1}{L} \sqrt{\frac{hc\lambda}{8F^2 P_{bs}}} \frac{1}{T_{ifo}(\tau_s, f)} \quad x_{RP}(f) = \frac{2F}{ML} \sqrt{\frac{2hP_{bs}}{\pi^3 c\lambda}} \frac{T_{ifo}(\tau_s, f)}{f^2}$$



- Increase SQL noise so it is larger than other noise sources at some frequency
- Increasing power decreases shot noise whilst increasing RP noise
- Make mirrors very light increases radiation pressure noise, shot noise is unchanged
- Light mirrors increases suspension thermal noise, potential solution, very low frequency pendulum
- Another knob, choice of  $T_{ifo}$

# If we succeed...



Calculation done by Corbitt using derivation by Buonanno and Chen