



# Development of a Stable Low-Frequency Squeezed Vacuum Source for Gravitational Wave Interferometers

**Keisuke Goda, Eugeny Mikhailov, Thomas Corbitt,  
Christopher Wipf, David Ottaway, Stan Whitcomb, Nergis Mavalvala**

Quantum Measurement Group, LIGO Laboratory  
Kavli Institute for Astrophysics and Space Research  
Massachusetts Institute of Technology

## Collaborators

Kirk McKenzie, Ping Koy Lam, Mal Gray, David McClelland  
The Australian National University

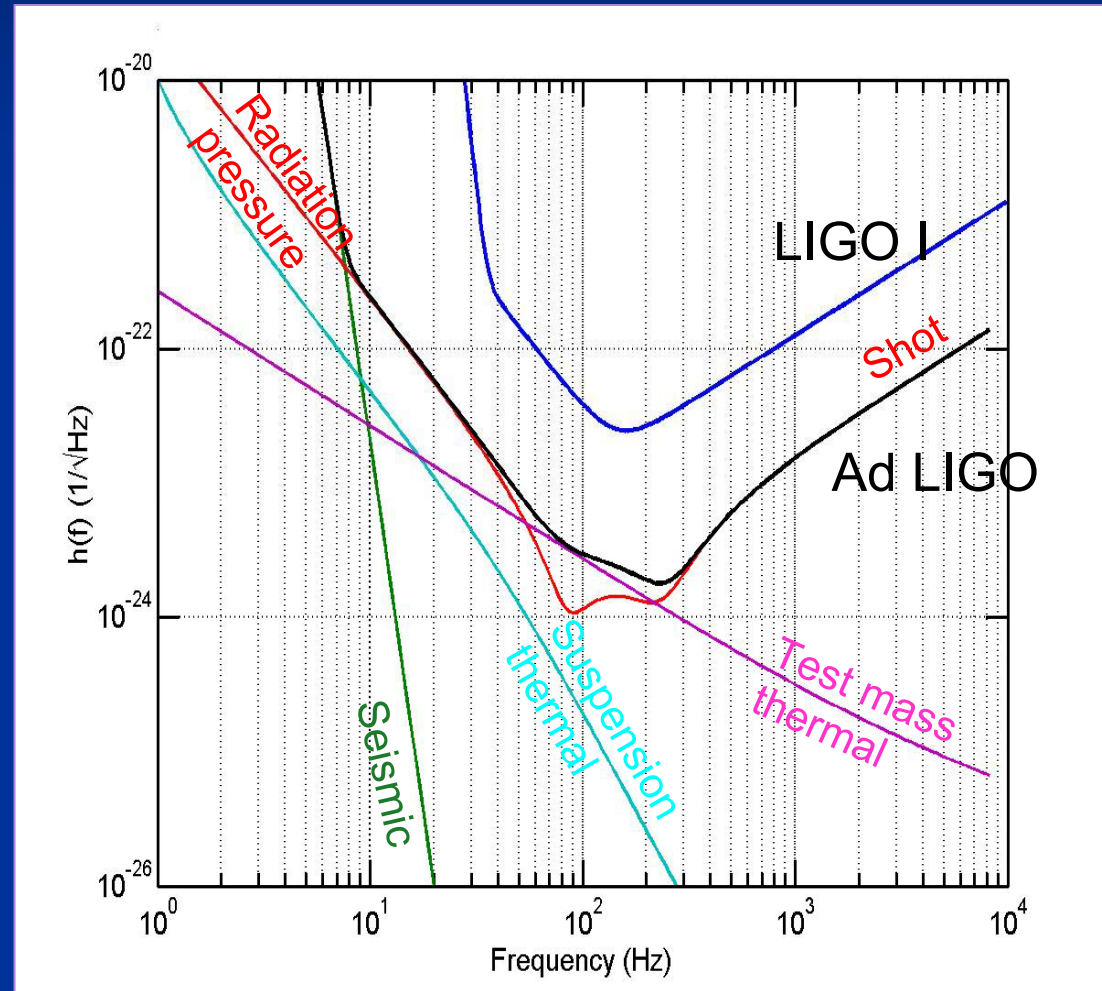
Caltech Special LIGO Seminar  
July 29, 2005

# Outline/Agenda

- ❖ Motivation and Goal
- ❖ Introduction to Quadrature Squeezing in OPO/OPA
- ❖ Some Results
- ❖ Requirements for Squeezing in GW Interferometers
- ❖ Work so far (by me, the group, and collaborators)
- ❖ Summary and Future Plans

# Motivation and Goal

- ❖ The sensitivity of the next generation GW interferometers will be limited by **quantum noise** (and thermal noise).
- ❖ To go beyond that, must reduce the quantum noise
- ❖ Quantum noise:
  - **Radiation pressure noise**  
at low frequencies  
(below  $\sim 100\text{Hz}$ )
  - **Shot noise**  
at high frequencies  
(above  $\sim 100\text{Hz}$ )



# Quantum Noise in GW Interferometers

- ❖ The sensitivity of next generation GW interferometers will be limited by **quantum noise**.

## ❖ Shot Noise

- Uncertainty in number of photons detected  $\Rightarrow$
- Higher input power  $P_{bs}$   
 $\Rightarrow$  need low optical losses
- Frequency dependence  
 $\Rightarrow$  depends on light (GW signal) storage time in the interferometer
- **Important at high frequencies (>100Hz)**

$$h(f) \propto \sqrt{\frac{1}{P_{bs}}}$$

## ❖ Radiation Pressure Noise

- Photons impart momentum to cavity mirrors
- Fluctuations in the number of photons  $\Rightarrow$
- Lower input power,  $P_{bs}$
- **Important at low frequencies (<100Hz)**

$$h(f) \propto \sqrt{\frac{P_{bs}}{Mf^4}}$$

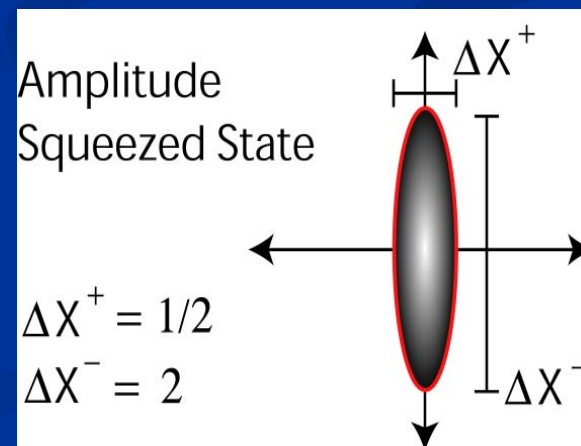
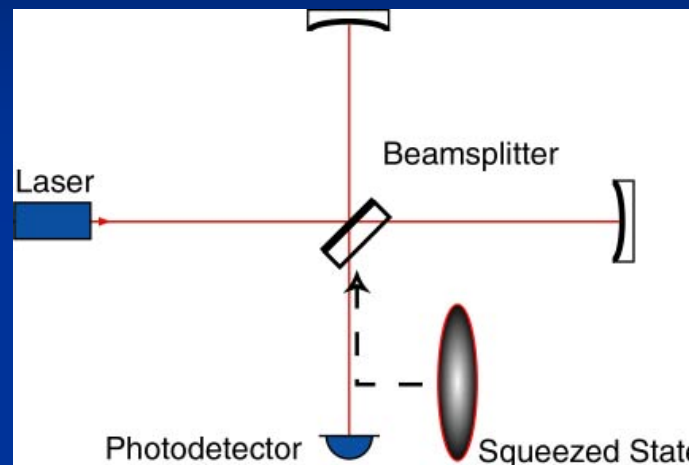
# Squeezed State in a Michelson

How to increase the sensitivity of GW interferometers?



- ❖ **Inject** a squeezed field into the **dark port** of a Michelson interferometer to replace vacuum noise

C.M. Caves, Phys. Rev. D **23**, 1693 (1981)



K. McKenzie

Before going into the details of the issues..

# **Introduction to Quadrature Squeezing in Optical Parametric Oscillation/Amplification (OPO/OPA)**

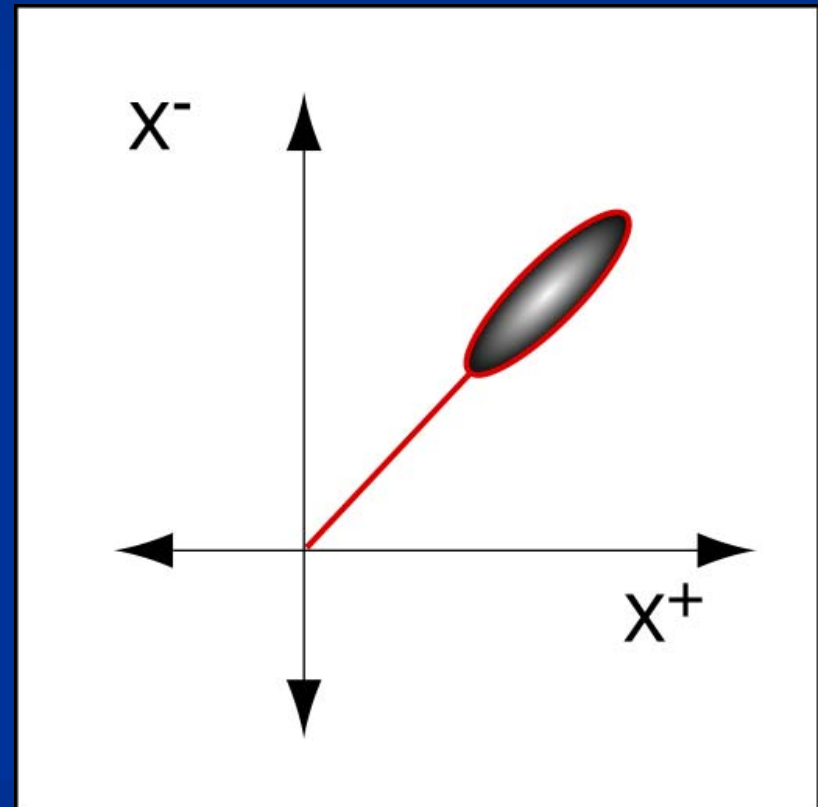
# Quantum States of Light/Vacuum

- Analogous to the phasor diagram
- Stick  $\rightarrow$  dc term
- Ball  $\rightarrow$  fluctuations
- Common states
  - **Coherent state**
  - **Vacuum state**
  - **Amplitude squeezed state**
  - **Phase squeezed state**

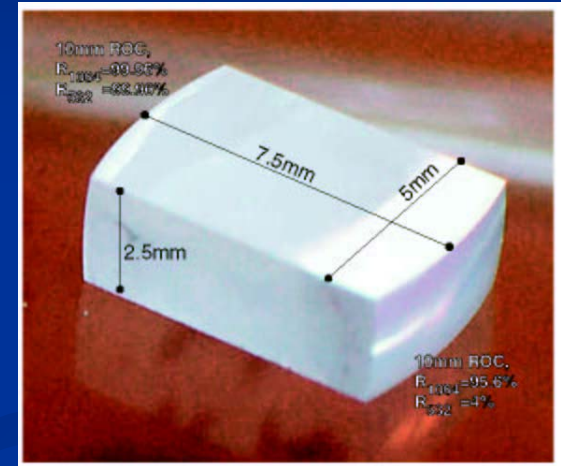
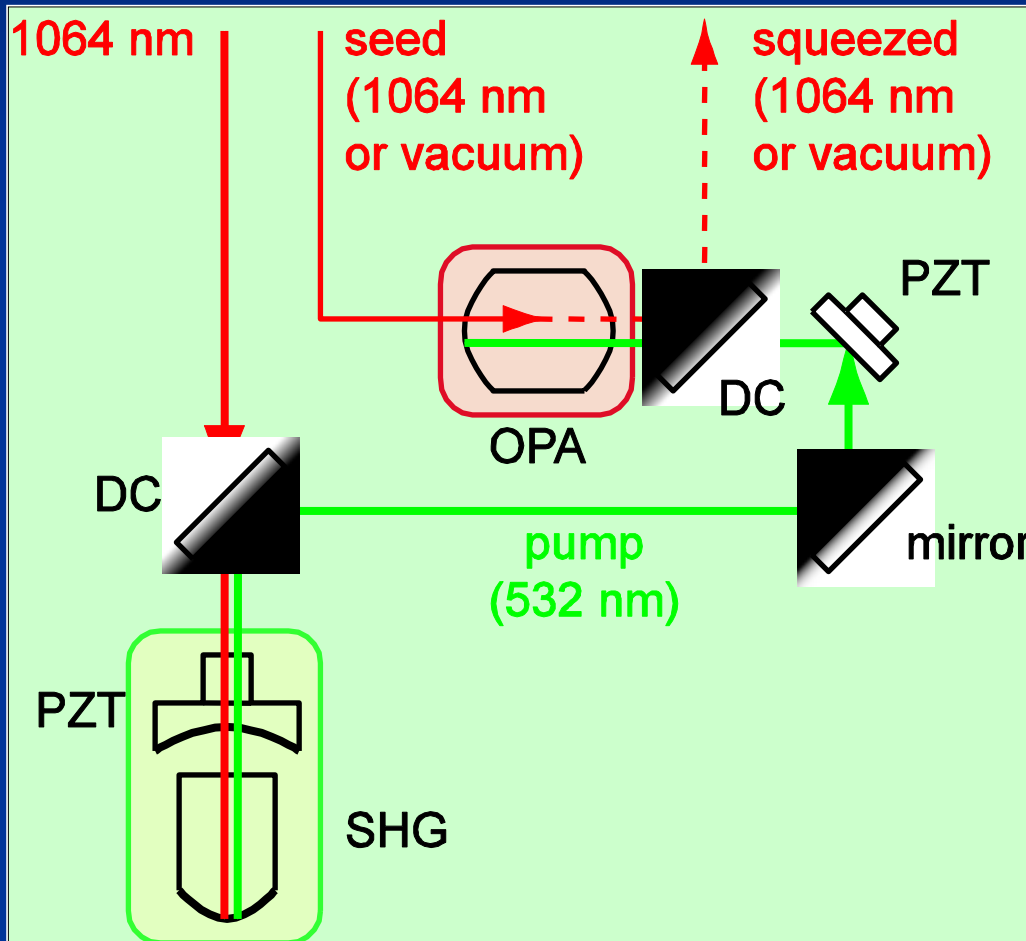
$$\begin{cases} X_1 = a + a^+ \\ X_2 = i(a - a^+) \end{cases}$$

1 or + = amplitude

2 or - = phase



# Generation of Squeezed States in Optical Parametric Amplification (OPA)



OPA crystal (MgO:LiNbO<sub>3</sub>)

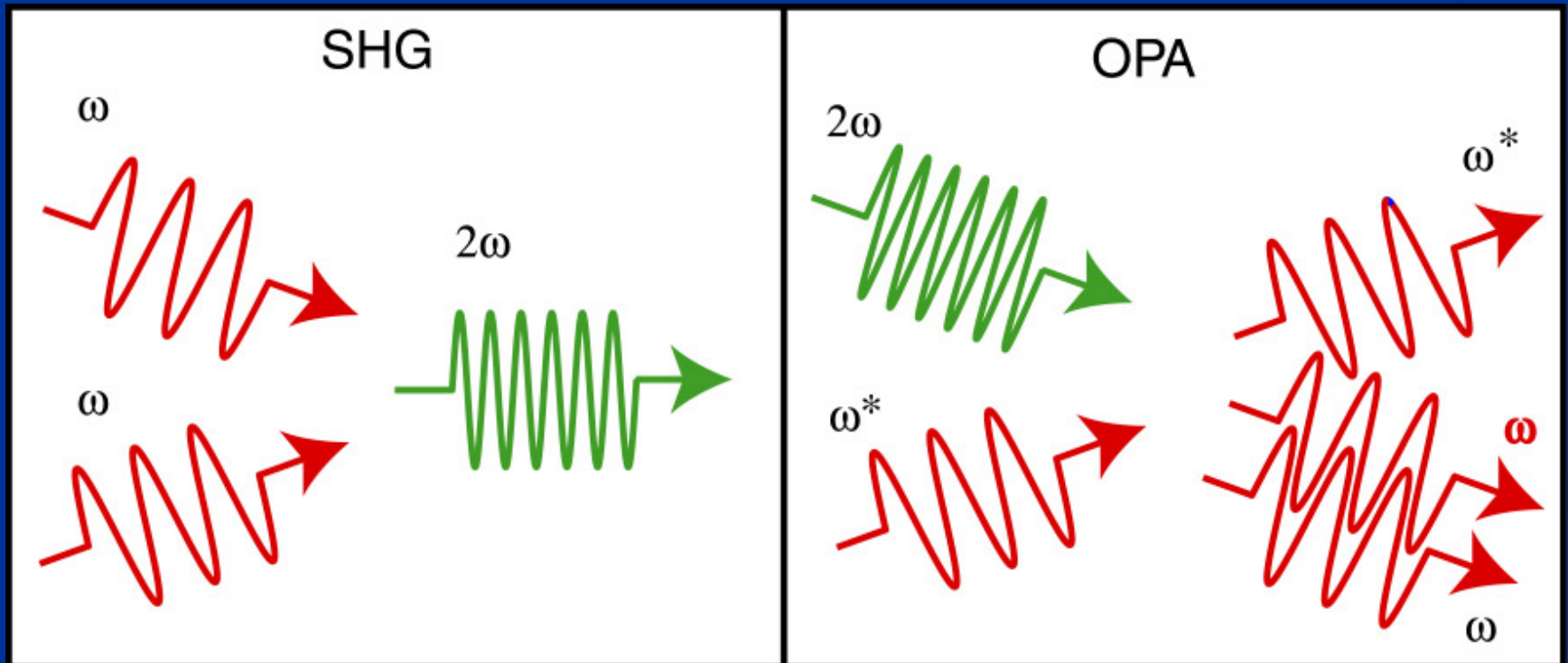
Typically, the OPA is a cavity at 1064nm and single/double pass at 532nm



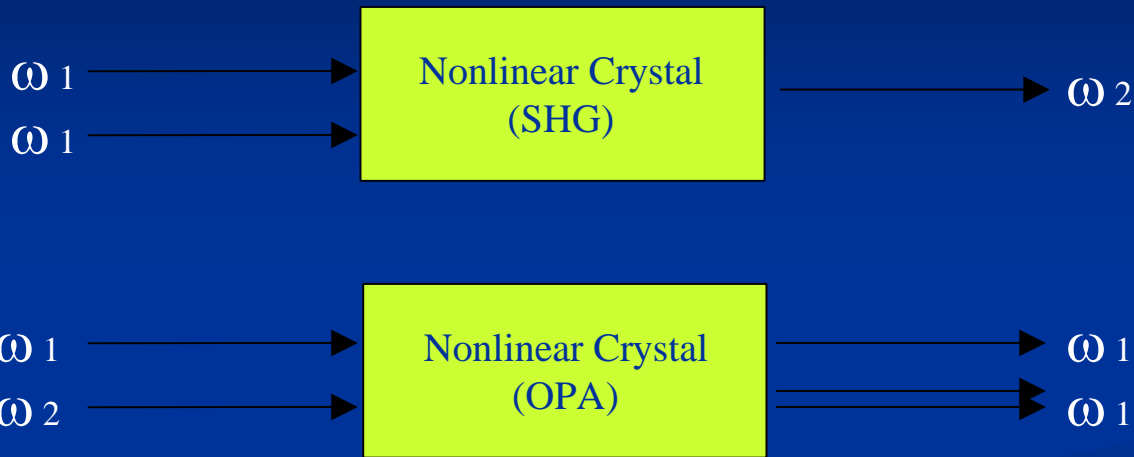
# Before Squeezing, A Little Introduction to Nonlinear Optics

- Optical parametric amplification (OPA)
- Seed (1064nm): field “a”
- Pump (532nm): field “b”

$$H = \hbar\omega_1 a^\dagger a + \hbar\omega_2 b^\dagger b + \frac{1}{2} i\hbar\kappa (ba^{+2} - b^+a^2)$$



# Required Conditions

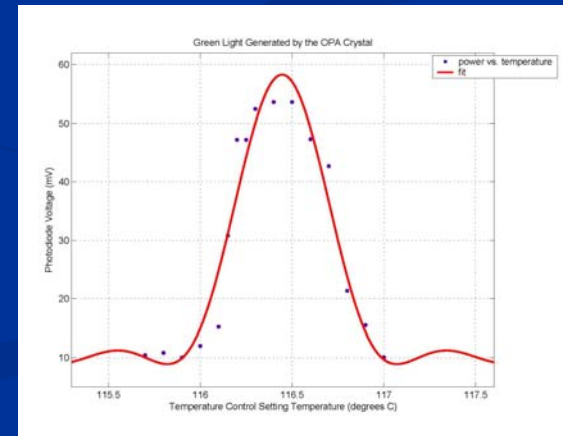


- Energy conservation

$$2\omega_1 = \omega_2$$

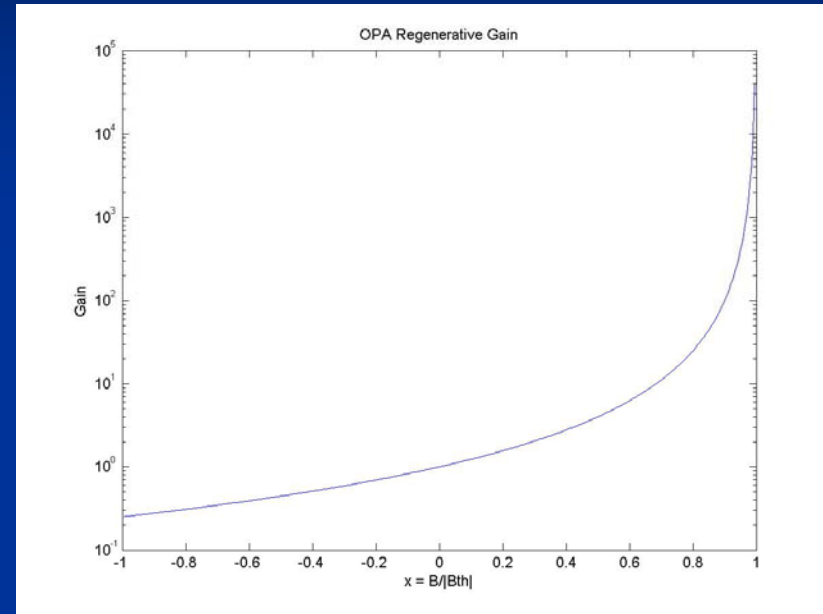
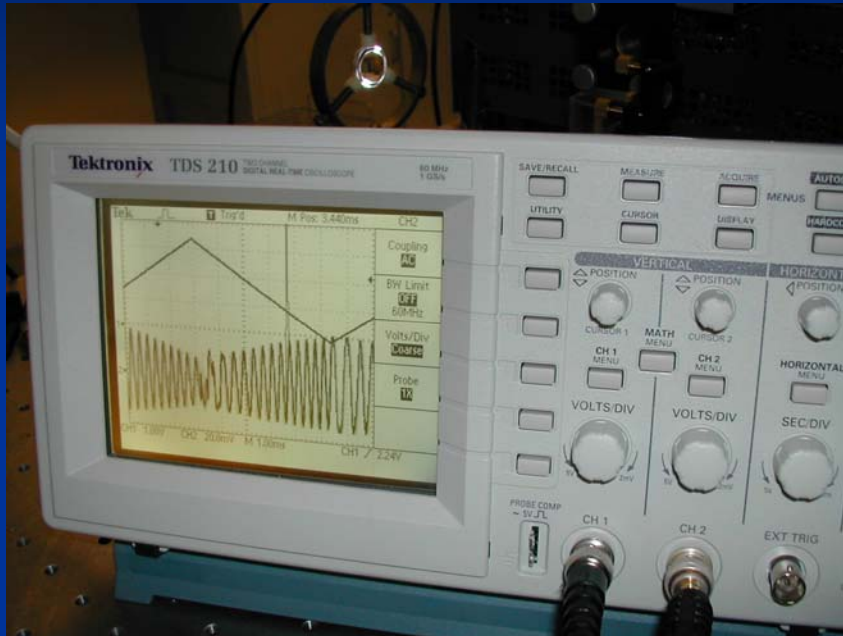
- Momentum conservation (phase-matching condition)

$$2k_1 = k_2$$



The phase-matching parameter is temperature-dependent, and therefore the crystal requires temperature control.

# Parametric Amplification/De-amplification



- Ramp signal driving the phase-matching PZT (upper)
- Parametric amplification and de-amplification of the seed beam (lower)
- Parametric amplification or deamplification is determined by the relative phase between the seed and pump

$$G = \frac{P_{(\varepsilon)}^{out}}{P_{(\varepsilon=0)}^{out}} = \left( \frac{1+x}{1-x^2} \right)^2, \text{ where } x = \frac{B_{in}}{|B_{in}^{threshold}|}$$

# Simple Model of Squeezing 1

$$H = \hbar\omega_1 a^\dagger a + \hbar\omega_2 b^\dagger b + \frac{1}{2}i\hbar\kappa(ba^{\dagger 2} - b^\dagger a^2)$$

In the Heisenberg picture, using

$$\frac{d}{dt}O(t) = \frac{i}{\hbar}[H, O(t)] \quad \text{and} \quad [a(t), a^\dagger(t)] = 1$$

The Heisenberg equations of motion for the signal mode:

$$\begin{cases} \frac{d}{dt}a(t) = -i\kappa\beta a^\dagger(t) - \frac{\Gamma}{2}a + N(t) \\ \frac{d}{dt}a^\dagger(t) = i\kappa\beta a(t) - \frac{\Gamma}{2}a^\dagger + N^\dagger(t) \end{cases}$$

$\Gamma$  = the cavity decay rate,  $N(t)$  = the associated noise operator.

# Simple Model of Squeezing 2

For the signal initially in a vacuum state, the expectation values in the steady state are

$$\begin{aligned}\langle a(t) \rangle_{ss} &= \langle a^+(t) \rangle_{ss} = 0 \\ \langle a^2(t) \rangle_{ss} &= \langle a^{+2}(t) \rangle_{ss} = \frac{-\Gamma \kappa \beta}{2(\Gamma^2 - (\kappa \beta)^2)} \\ \langle a(t)a^+(t) + a^+(t)a(t) \rangle_{ss} &= \frac{\Gamma^2}{2(\Gamma^2 - (\kappa \beta)^2)}\end{aligned}$$

The field amplitudes in the amplitude and phase quadratures are defined by  $X_1$  and  $X_2$  (1 = amplitude, 2 = phase). The noise variances of the field amplitudes in the steady state are given by

$$\begin{cases} X_1 \equiv a + a^+ \\ X_2 \equiv i(a - a^+) \end{cases}$$

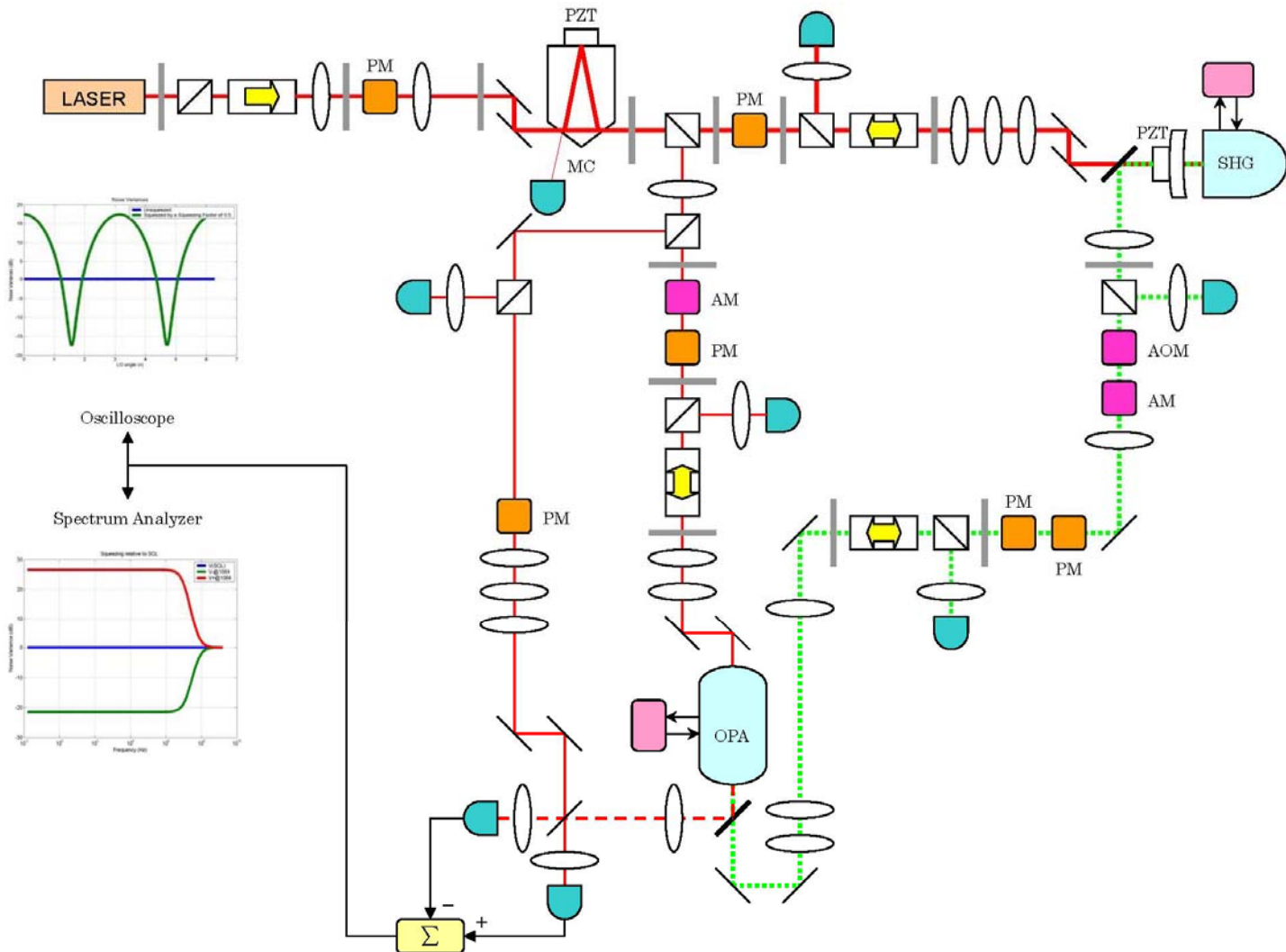
$$\begin{cases} \Delta X_1 \equiv \sqrt{\langle X_1^2 \rangle - \langle X_1 \rangle^2} = \sqrt{\frac{\Gamma}{\Gamma + \kappa \beta}} < 1 \\ \Delta X_2 \equiv \sqrt{\langle X_2^2 \rangle - \langle X_2 \rangle^2} = \sqrt{\frac{\Gamma}{\Gamma - \kappa \beta}} > 1 \end{cases}$$

**$\Delta X_1$  squeezed!!**

The level of squeezing increases with the coupling constant,  $\kappa$ .

# Squeezer Layout

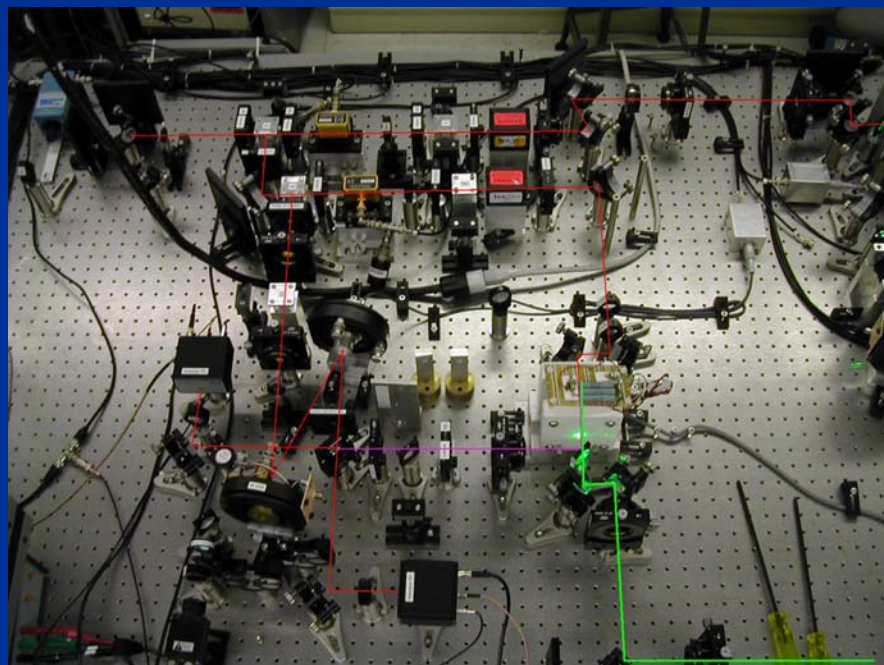
## Squeezer Layout



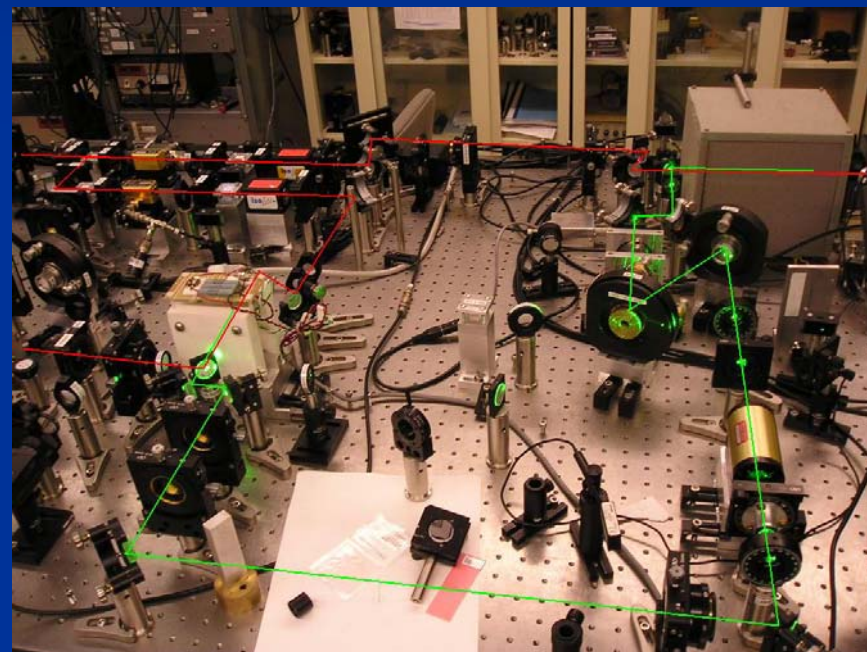
# Squeezer @ MIT



# Squeezer @ MIT



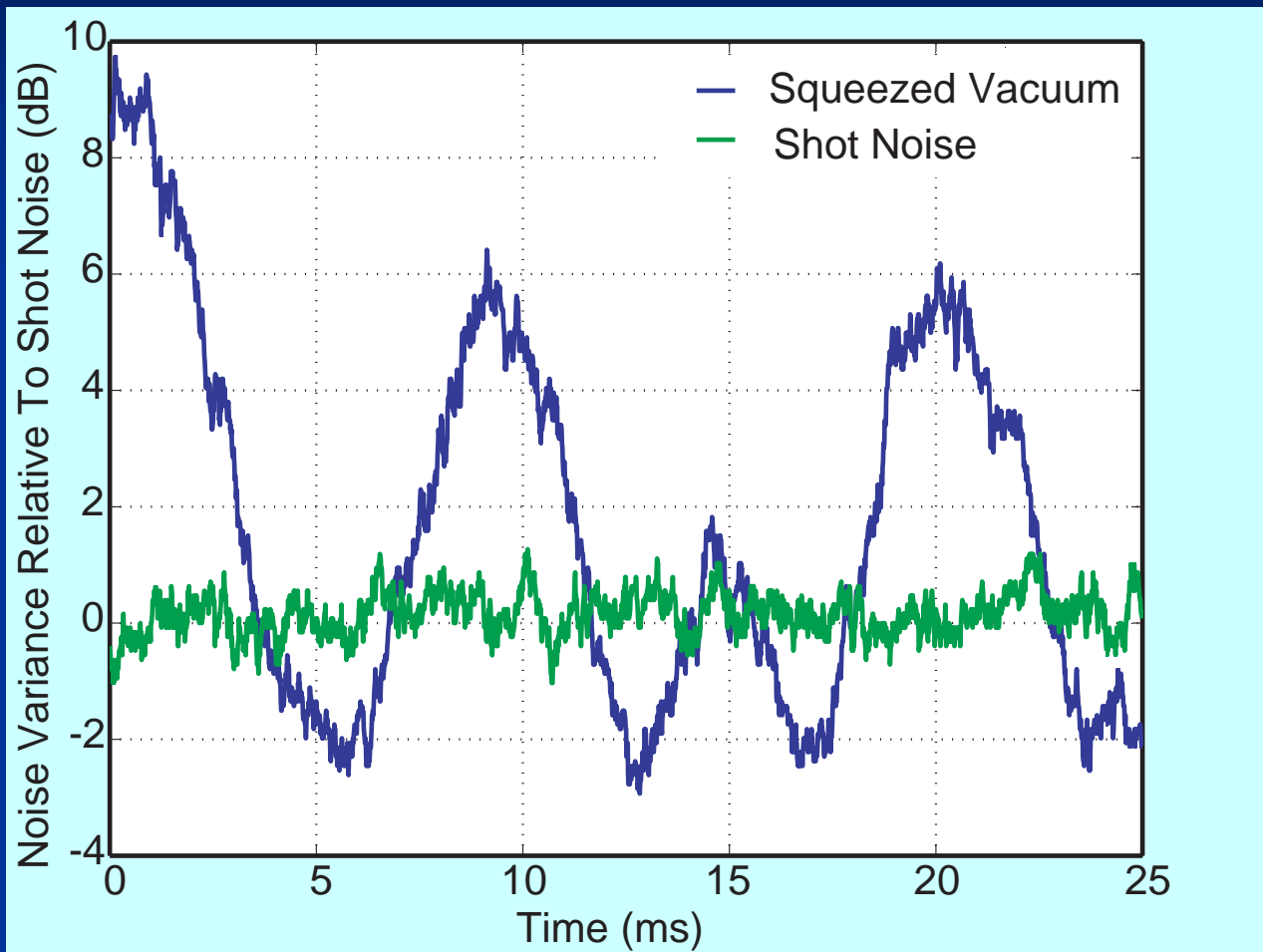
Homodyne Detector



OPA and SHG



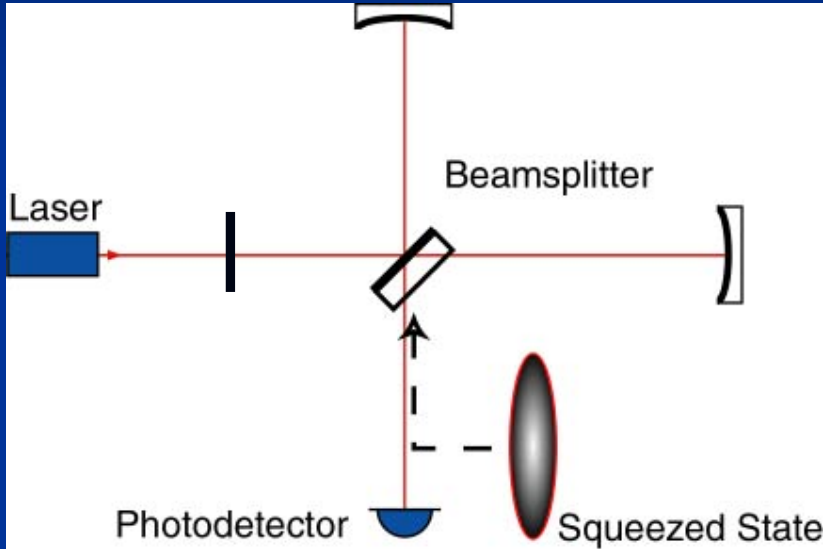
# Squeezing Result



K. Goda

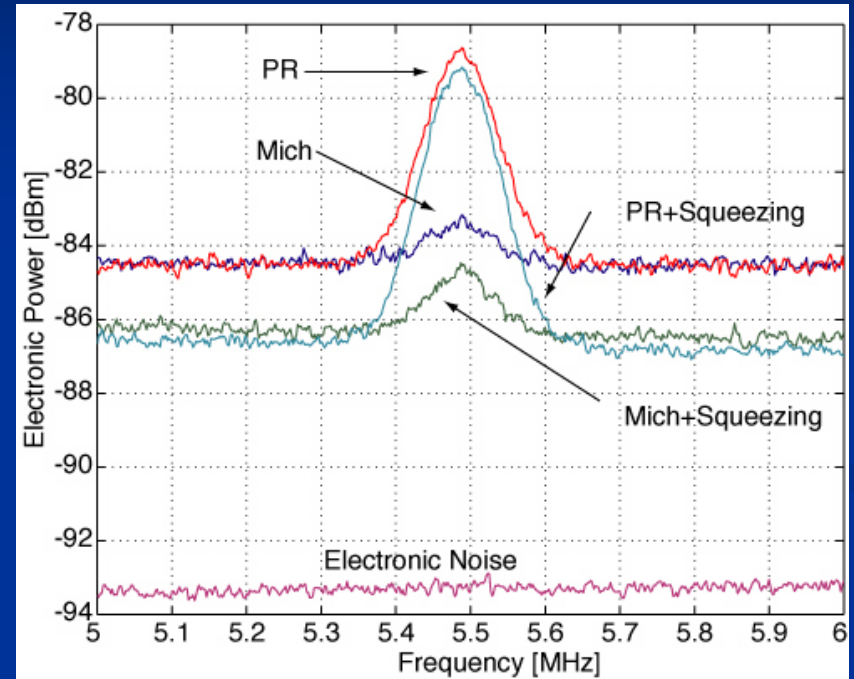
Roughly 10dB of squeezing was created in the OPO cavity, but it was reduced down to 3dB of squeezing due to losses.

# Table-Top Demonstration



K. McKenzie

table-top demonstration of squeezing in a Michelson interferometer with a power recycling mirror



K. McKenzie

Injection of squeezed **light** into the dark port

K. McKenzie, B.C. Buchler, D.A. Shaddock, P.K. Lam, and D.E. McClelland, Phys. Rev. Lett. 88, 231102 (2002)



# Requirements for Squeezing in GW Interferometers (practical issues and what has been achieved)

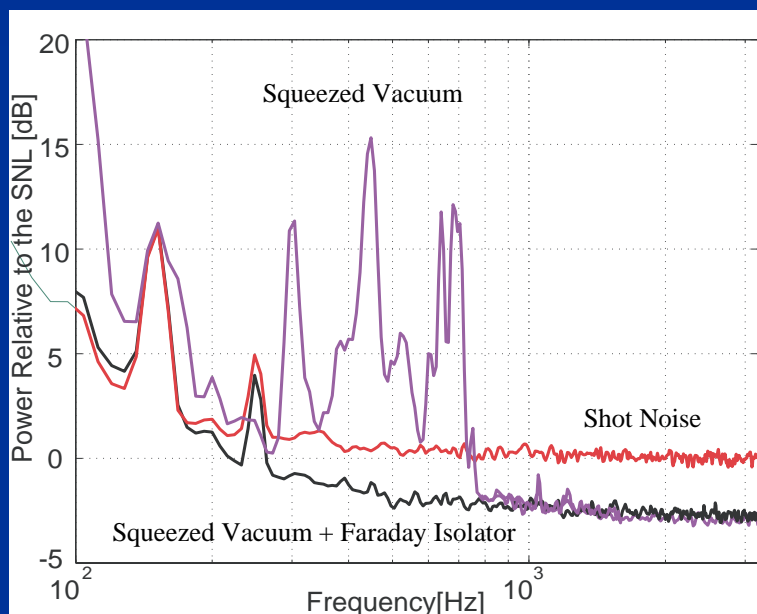
# Requirements for Squeezing in GW Interferometers

- ❖ Squeezing in the GW Band (10Hz – 10kHz)
  - Squeezed vacuum preferred as opposed to squeezed light
  - Photothermal noise, laser excess noise, etc..
- ❖ Increased Level of Squeezing (ideally 10dB or higher)
  - Again, squeezed vacuum preferred to have a shot noise limited seed
  - Use better crystals, change the OPO cavity configuration, etc..
- ❖ Long-Term Stability for Operation
  - Control signal required for phase-locking squeezed vacuum
  - Noise-locking, frequency-shifted sub-carrier, etc..
- ❖ Frequency-Dependent Squeeze Angle
  - Phase-squeezing at high frequencies ( $<100\text{Hz}$ ) and amplitude-squeezing at low frequencies ( $>100\text{Hz}$ ) required
  - Long squeeze angle rotation cavities, squeeze amplitude filter cavities, etc..

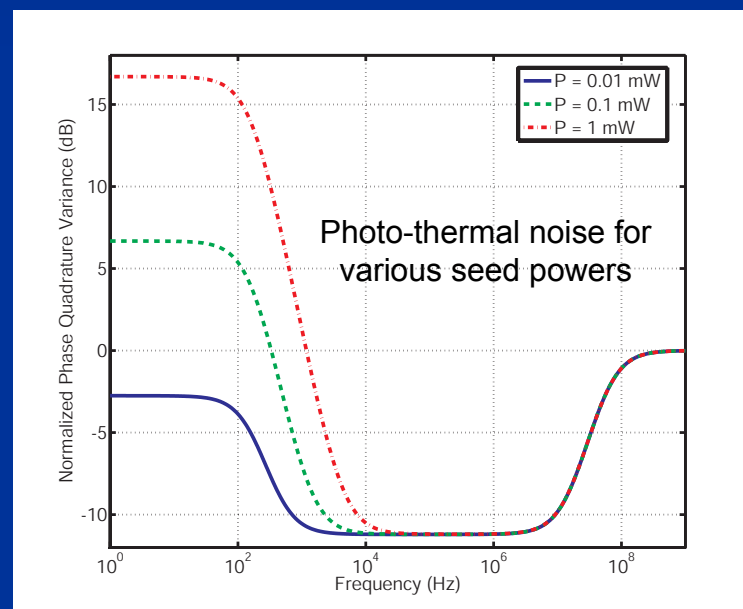


# Squeezing in the GW Band (10Hz – 10kHz)

- ❖ Historically, squeezing has been done at  $\sim$ MHz.
- ❖ Squeezed light has been used for applications.
- ❖ However, for application to GW interferometers, **squeezed vacuum** is preferred as opposed to squeezed light.



Noise-locked squeezed vacuum  
down to frequencies as low as 280Hz  
K. McKenzie, Nicolai Grosse, W.P. Bowen, S.E. Whitcomb, M.B. Gray, D.E. McClelland, and P.K. Lam, Phys. Rev. Lett. **93**, 161105 (2004)

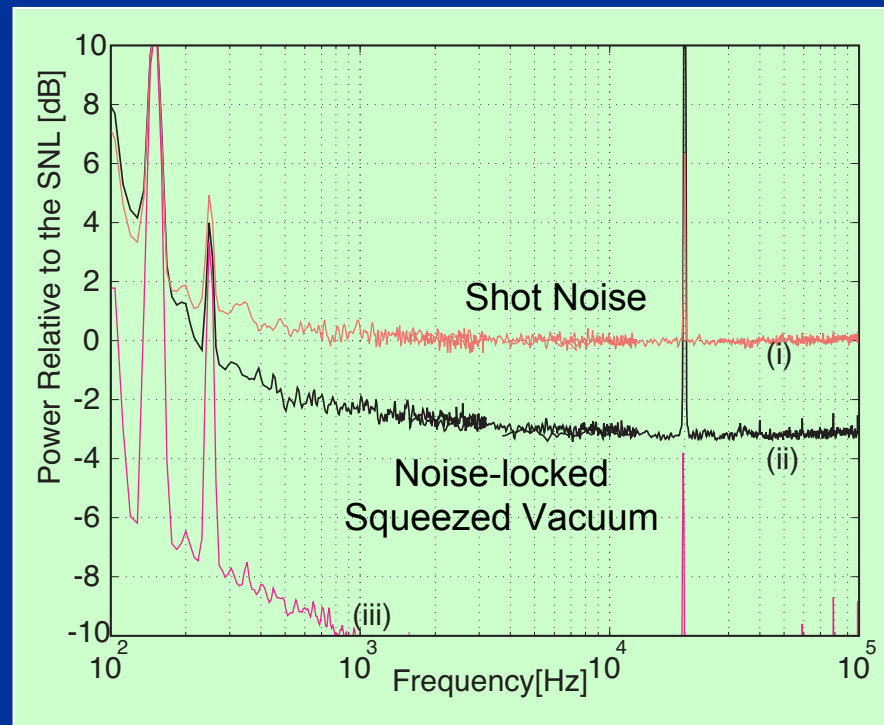
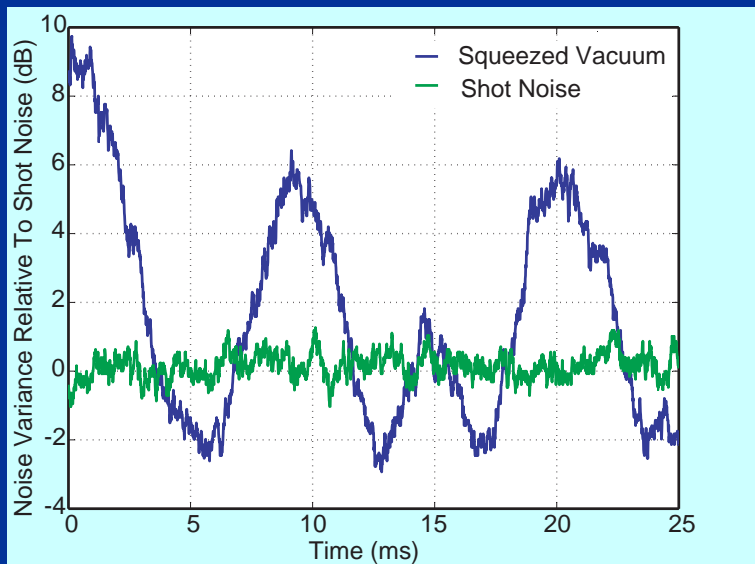


Limiting noise sources at low frequencies  
(photo-thermal noise, laser noise, etc..)  
K. Goda, K. McKenzie, E. Mikhailov, P.K. Lam, D. McClelland, and N. Mavalvala,  
submitted to Phys. Rev. A, quant-ph/0505154



# Phase-Locking Squeezed Vacuum for Long-Term Stability

- ❖ Squeezed vacuum is hard to control since it has no carrier light.
- ❖ Noise-locking and/or frequency-shifted sub-carrier technique can be used  
K. McKenzie, E. Mikhailov, K. Goda, P.K. Lam, N. Grosse, M.B. Gray, N. Mavalvala, and D.E. McClelland, *J. Opt. B: Quantum Semiclass. Opt.*, quant-ph/0505164



**Unlocked** squeezed vacuum..

If the squeezing spectrum itself is modulated to obtain a control signal..

**Noise-locked** squeezed vacuum

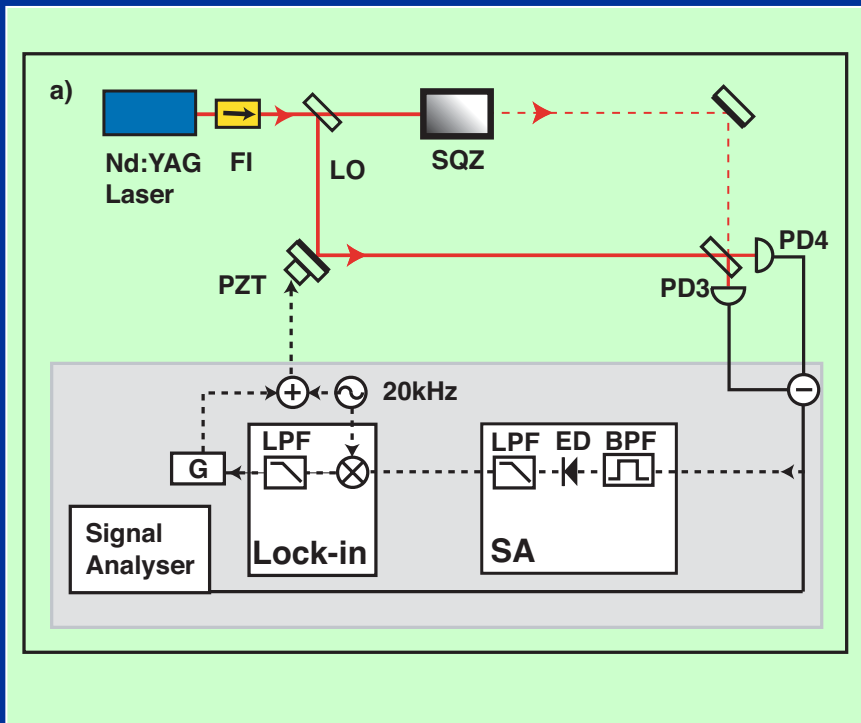
# Stability of Noise-Locking

- ❖ The stability of noise-locking depends on the level of squeezing.
- ❖ The more squeezing you have, the more stable your noise-locking is.

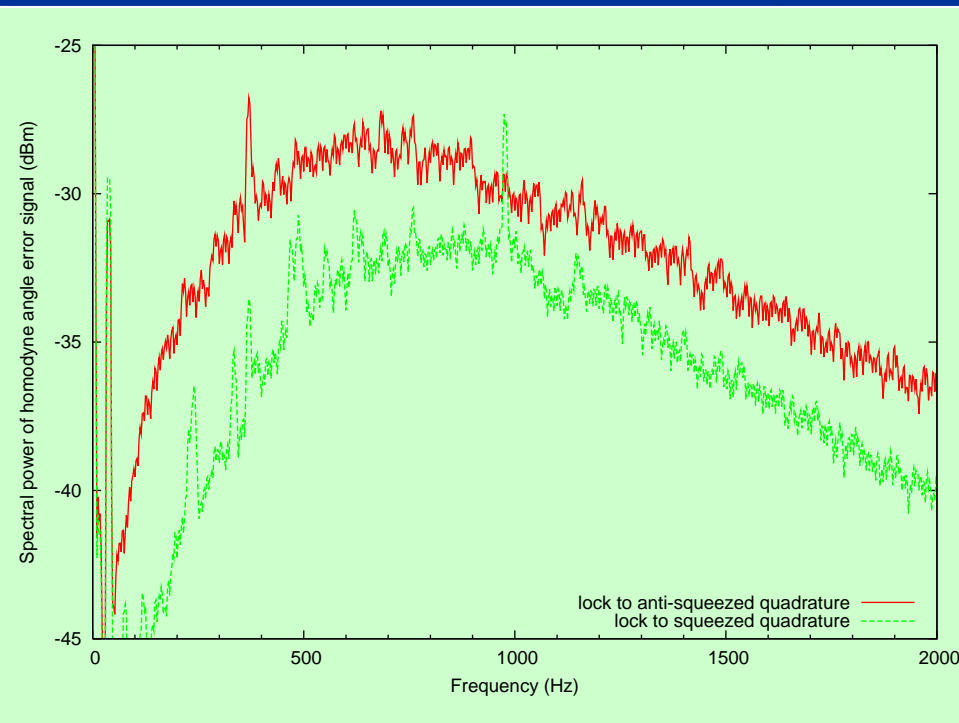
Instability or Fluctuations in Noise-Locked Angle

$$\Delta\theta \sim \frac{1}{\sqrt{e^{4R} - 1}} \left( \frac{1}{\Delta\omega} \right)^{1/4}$$

$\Delta\omega$  = detection bandwidth  
R = squeeze factor



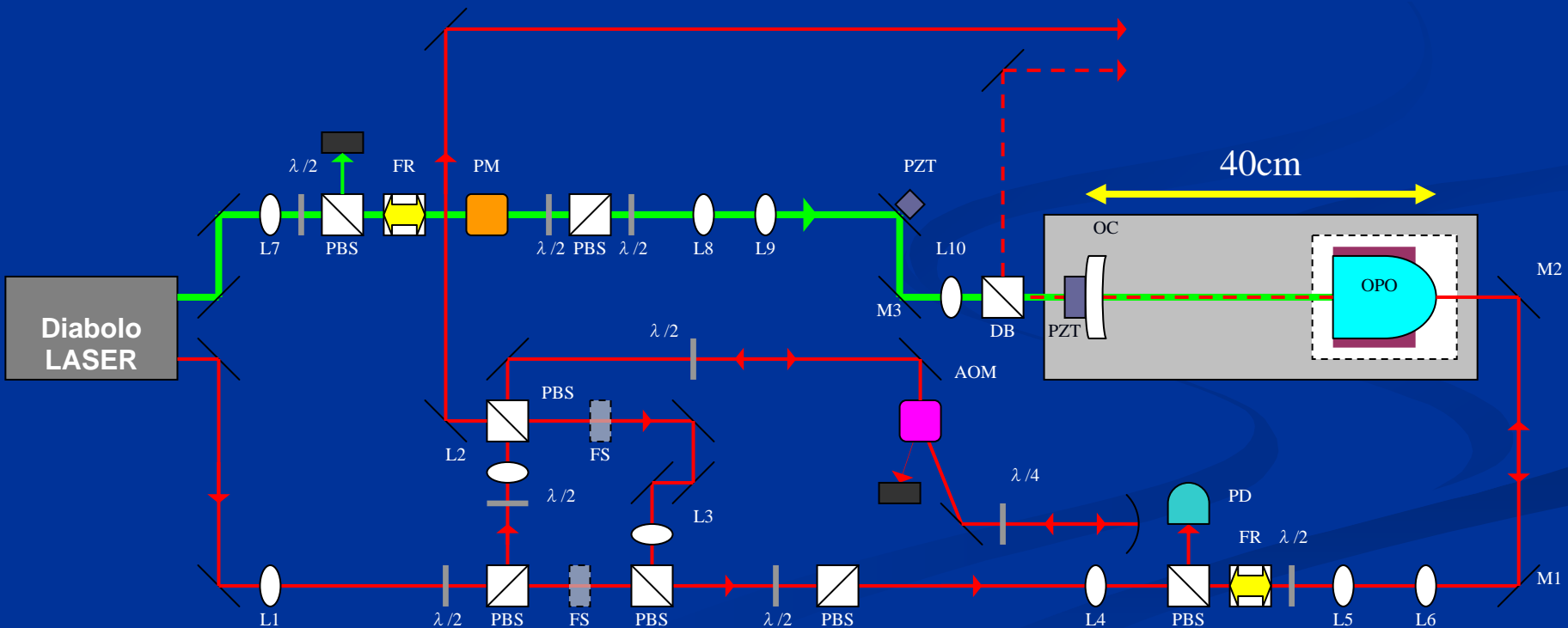
Schematic of the Noise-Locked Squeezer



Noise Power of the Noise-Locking Error Signal

# Locking Squeezed Vacuum with Coherent Light

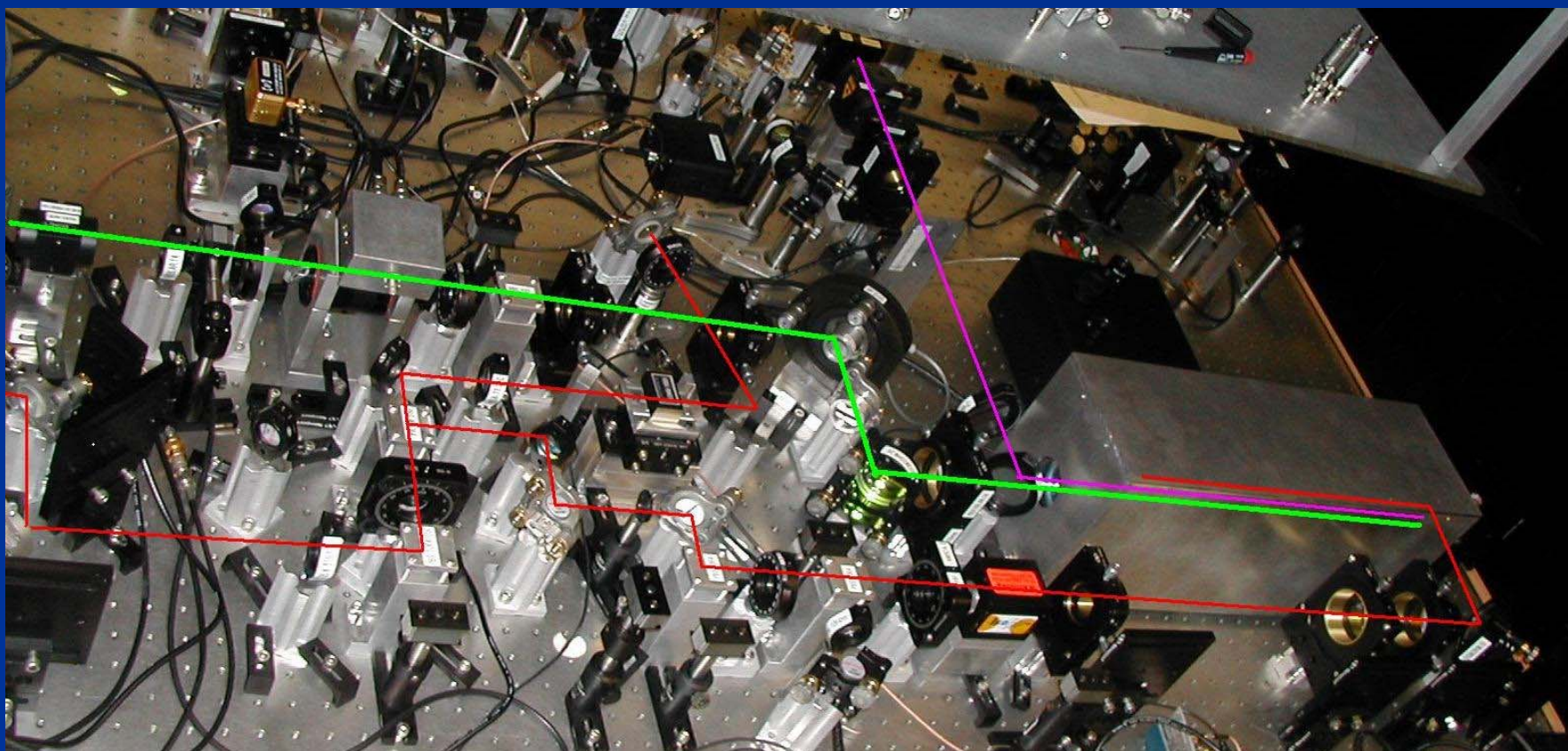
- ❖ Create bright sidebands on top of squeezed vacuum
- ❖ Extract the squeeze angle information by measuring the optical parametric amplification of the signal and idler beams which are at the carrier frequency  $\pm$  FSR





# Squeezed Vacuum with Bright Sidebands

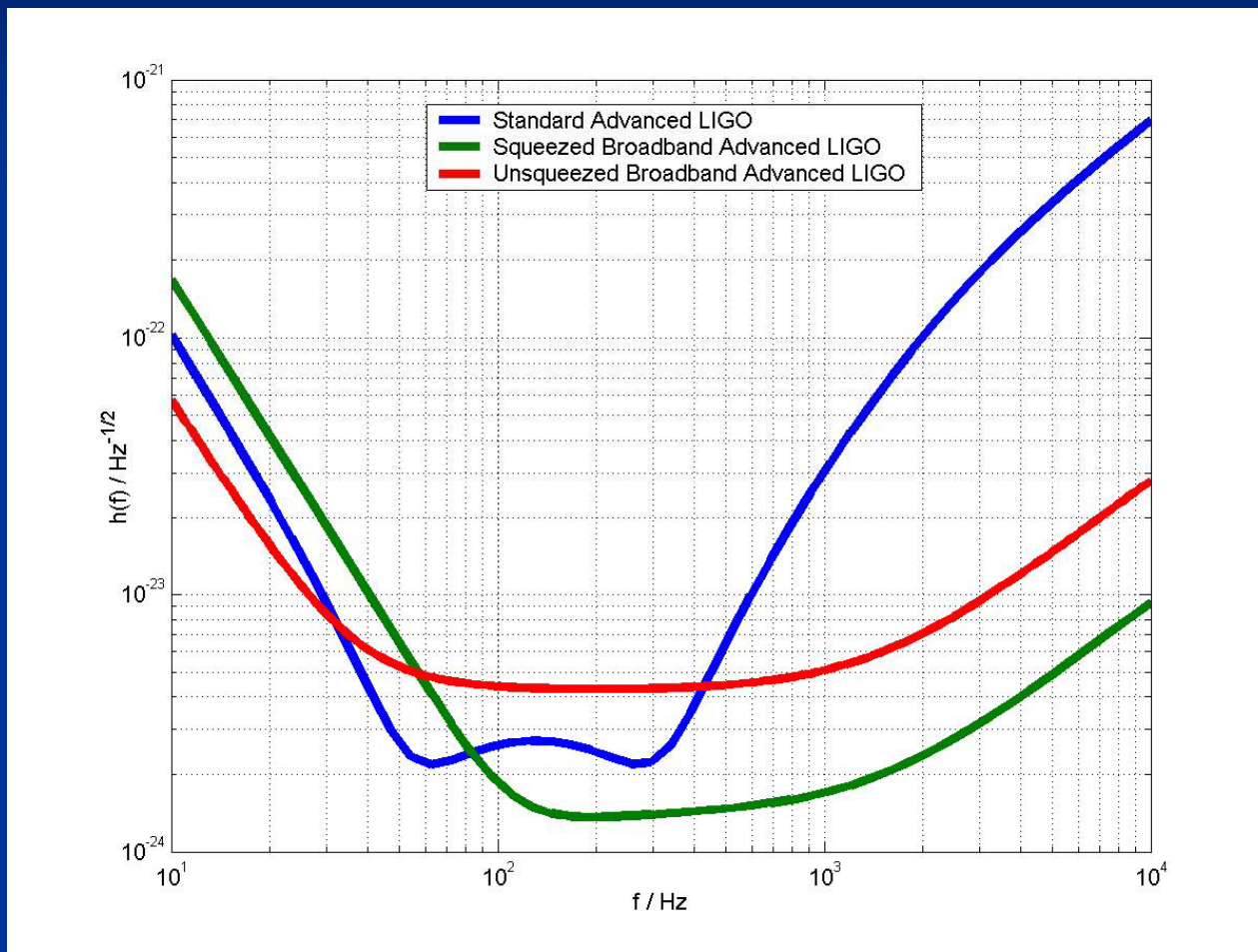
- $\sim 10\text{dB}$  of squeezing (before detection)
- Bright sidebands used to control the squeeze angle





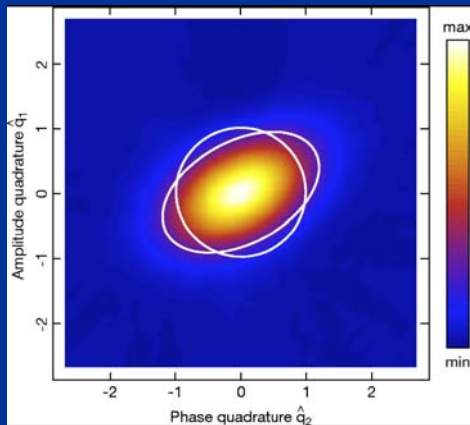
# Advanced LIGO

(The need for Frequency-Dependent Squeeze Angle)

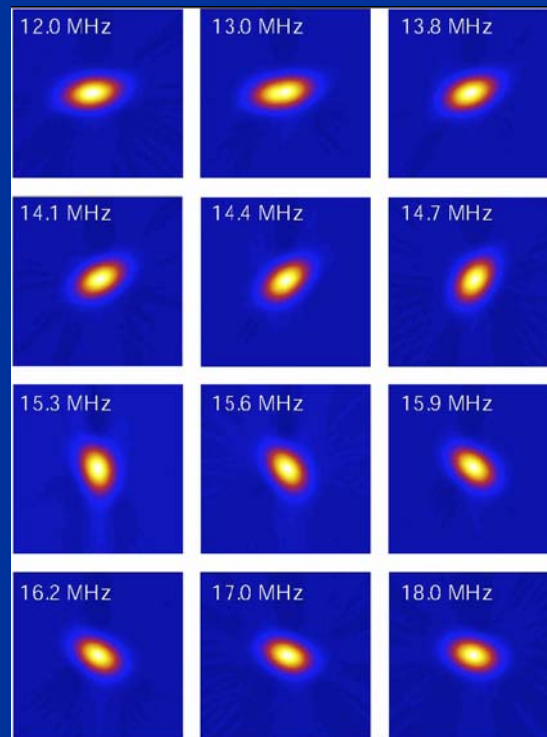


# Frequency-Dependent Squeeze Angle

- ❖ For broadband Advanced LIGO, the detector sensitivity will be limited by **phase/shot noise at high frequencies** ( $>100\text{Hz}$ ) and **amplitude/radiation-pressure noise at low frequencies** ( $<100\text{Hz}$ ).
- ❖ Need frequency-dependent squeeze angle, i.e., phase squeezing at high frequencies and amplitude squeezing at low frequencies



Wigner function



- ❖ Squeeze angle rotated by a filter cavity

S. Chelkowski, H. Vahlbruch, B. Hage, A. Franzen, N. Lastzka, K. Danzmann, and R. Schnabel, *Phys. Rev. A* **71**, 013806 (2005)

- ❖ However, this was done at  $\sim\text{MHz}$ , not ideal for  $10\text{Hz} - 10\text{kHz}$  for GW interferometers.

# Summary and Future Work

- ❖ Low-frequency squeezing with the noise-locking technique achieved
  - ❖ K. McKenzie, Nicolai Grosse, W.P. Bowen, S.E. Whitcomb, M.B. Gray, D.E. McClelland, and P.K. Lam, Phys. Rev. Lett. **93**, 161105 (2004)
  - ❖ K. Goda, K. McKenzie, E. Mikhailov, P.K. Lam, D. McClelland, and N. Mavalvala, submitted to Phys. Rev. A, quant-ph/0505154
  - ❖ K. McKenzie, E. Mikhailov, K. Goda, P.K. Lam, N. Grosse, M.B. Gray, N. Mavalvala, and D.E. McClelland, J. Opt. B: Quantum Semiclass. Opt., quant-ph/0505164
- ❖ Experimental demonstration of noise-locking
  - ❖ K. McKenzie, E. Mikhailov, K. Goda, P.K. Lam, N. Grosse, M.B. Gray, N. Mavalvala, and D.E. McClelland, J. Opt. B: Quantum Semiclass. Opt., quant-ph/0505164
- ❖ A few different control schemes of squeezed vacuum under way
  - ❖ to be submitted to Phys. Rev. A
- ❖ Theoretical and experimental work on the increased level of squeezing under way
  - ❖ to be submitted to Phys. Rev. A or JOSA
- ❖ Theoretical work on frequency-dependent squeeze angle rotation and increased level of squeezing under way
  - ❖ to be submitted to Phys. Rev. Lett.
- ❖ **Future Plans**
  - Testing DC readout-compatible squeezing with 40M