

Search for Stochastic Gravitational Waves with LIGO

XLIInd Rencontres de Moriond

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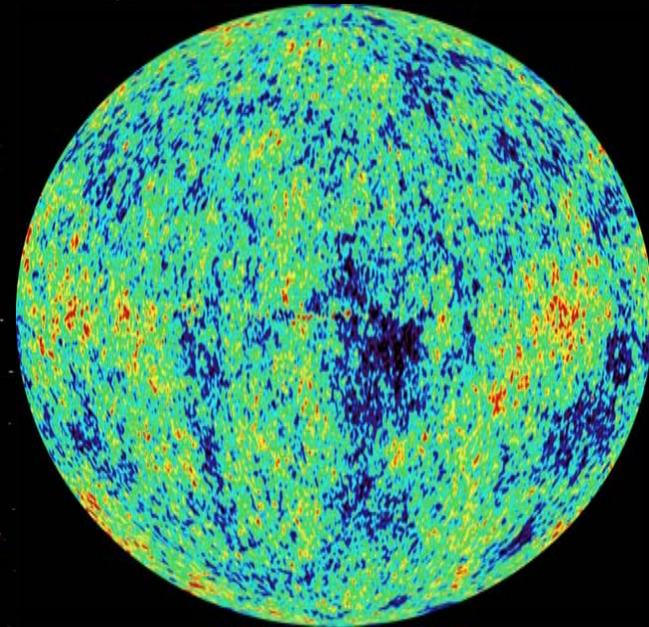
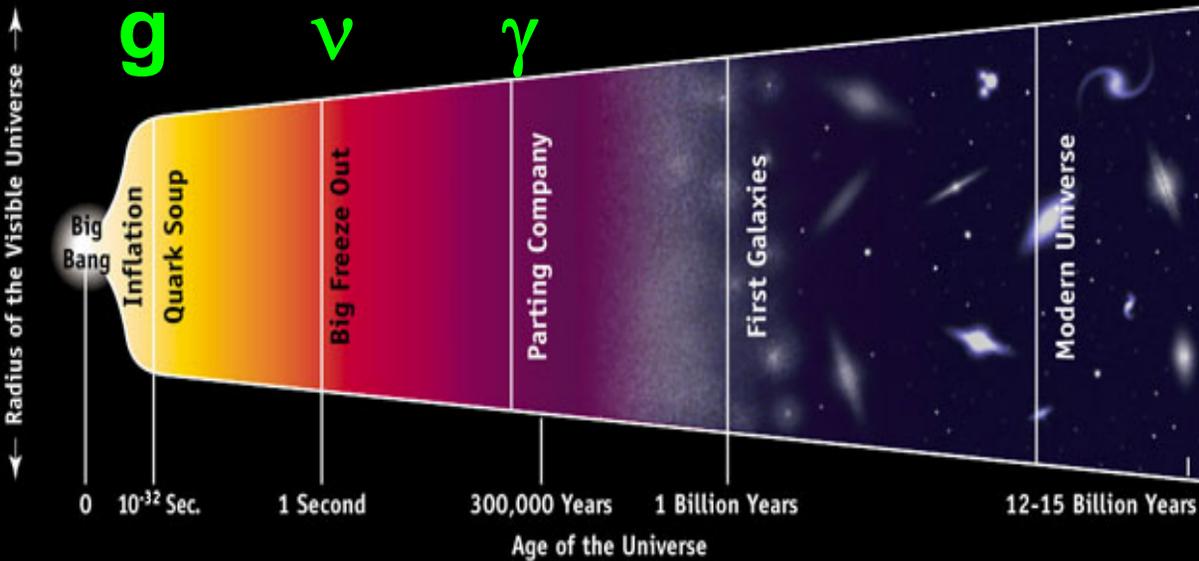
(on behalf of the LIGO Scientific Collaboration)

La Thuile, Val d'Aosta, Italy

March 11-18, 2007

- Stochastic gravitational waves
 - Sources
 - Characterization of strain
- Search techniques
 - All-sky averaged search
 - Spatially resolved map
- Recent observational results
- Prospects

- GWs are the able to probe the very early universe



Analog from cosmic microwave background -- WMAP 2003

- Search for a GW background by cross-correlating interferometer outputs in pairs
 - US-LIGO : Hanford, Livingston
 - Europe: GEO600, Virgo
 - Japan: TAMA
- Good sensitivity requires:
 - $\bullet_{GW} \geq 2D$ (detector baseline)
 - $f \leq 40$ Hz for LIGO pair over 3000 km baseline
 - Good low-f sensitivity and short baselines

Stochastic signals

Cosmological processes

- Inflation -- flat spectrum (Turner)
- Phase transitions -- peaked at phase transition energy scale (Kamionkowski)
- Cosmic strings -- gradually decreasing spectrum (Damour & Vilenkin)
- Pre big-bang cosmology -- rising strength with f (Buonanno et al.)

Astrophysical Foregrounds

- Incoherent superposition of many signals from various signal classes (Ferrari, Regimbau)
 - Coalescing binaries
 - Supernovae
 - Pulsars
 - Low Mass X-Ray Binaries (LMXBs)
 - Newly born neutrons stars
 - Normal modes - \bar{R} modes
 - Binary black holes
 - Black hole ringdowns
- Gaussian \rightarrow non-Gaussian (“popcorn” noise) depending on rates
Drasco & Flannigan, Phys.Rev. D67 (2003)
- Spectra follow from characteristics of individual sources -- e.g., spatial anisotropy for nearby (foreground) contributors

Characterization of a Stochastic Gravitational Wave Background

- GW energy density given by time derivative of strain

$$\rho_{gw} = \frac{1}{32\pi G} \overline{\dot{h}_{ij}(t)\dot{h}_{ij}(t)}$$

- Assuming SGWB is isotropic, stationary and Gaussian, the strength is fully specified by the energy density in GWs

$$\frac{\rho_{gw}}{\rho_c} \equiv \int_0^\infty \Omega_{gw}(f) d\ln f$$

- $\Omega(f)$ in terms of a *measurable* strain power spectrum, $S_{gw}(f)$:

$$S_{gw}(f) = \frac{3H_0^2}{10\pi^2} \frac{\Omega_{gw}(f)}{f^3}$$

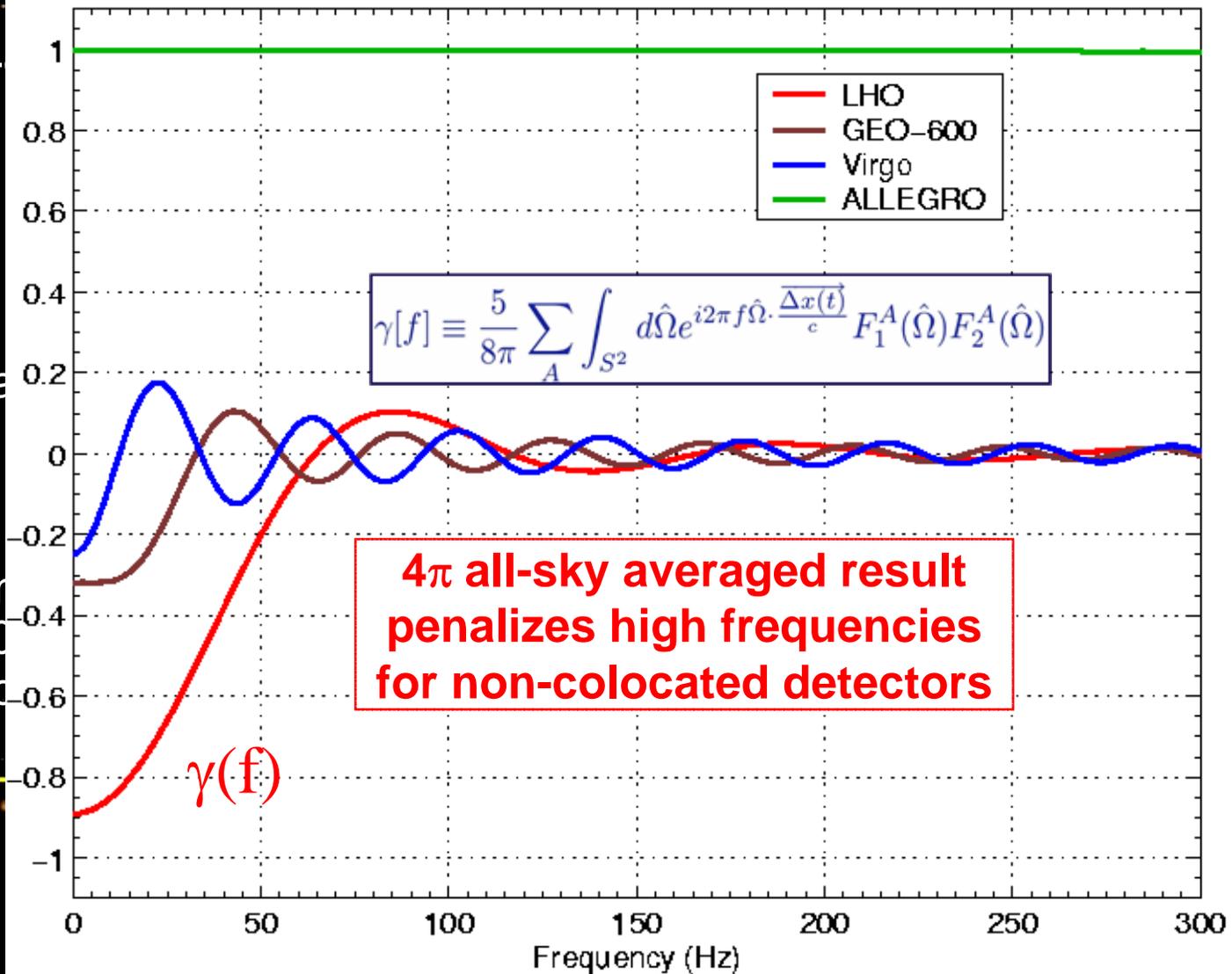
- Strain amplitude scale:

$$\tilde{h}(f) = \sqrt{S_h(f)} = 6.3 \times 10^{-22} \sqrt{\Omega_{gw}(f)} \left(\frac{100\text{Hz}}{f}\right)^{3/2} \text{Hz}^{-1/2}$$

Allen & Romano, Phys.Rev. D59 (1999)

Technique -- Sky-averaged search Time-independent Detection Strategy

Example: Overlap Reduction Function (LLO and other detectors)



■ Cross-

■ Optima

■ Make n
(T = 60
variatio

$$Y_i = \Omega_i T$$

$$\frac{df}{(f)^2}$$

$$\frac{(f)}{(f)}$$

$$1/\sigma_i^2$$



Sensitivities of LIGO interferometers during S4 22 February - 23 March 2005

H1 - L1: 353.9 hr

H2 - L1: 332.7 hr.

Scheduled for publication March 2007 : ApJ 658 (2007)

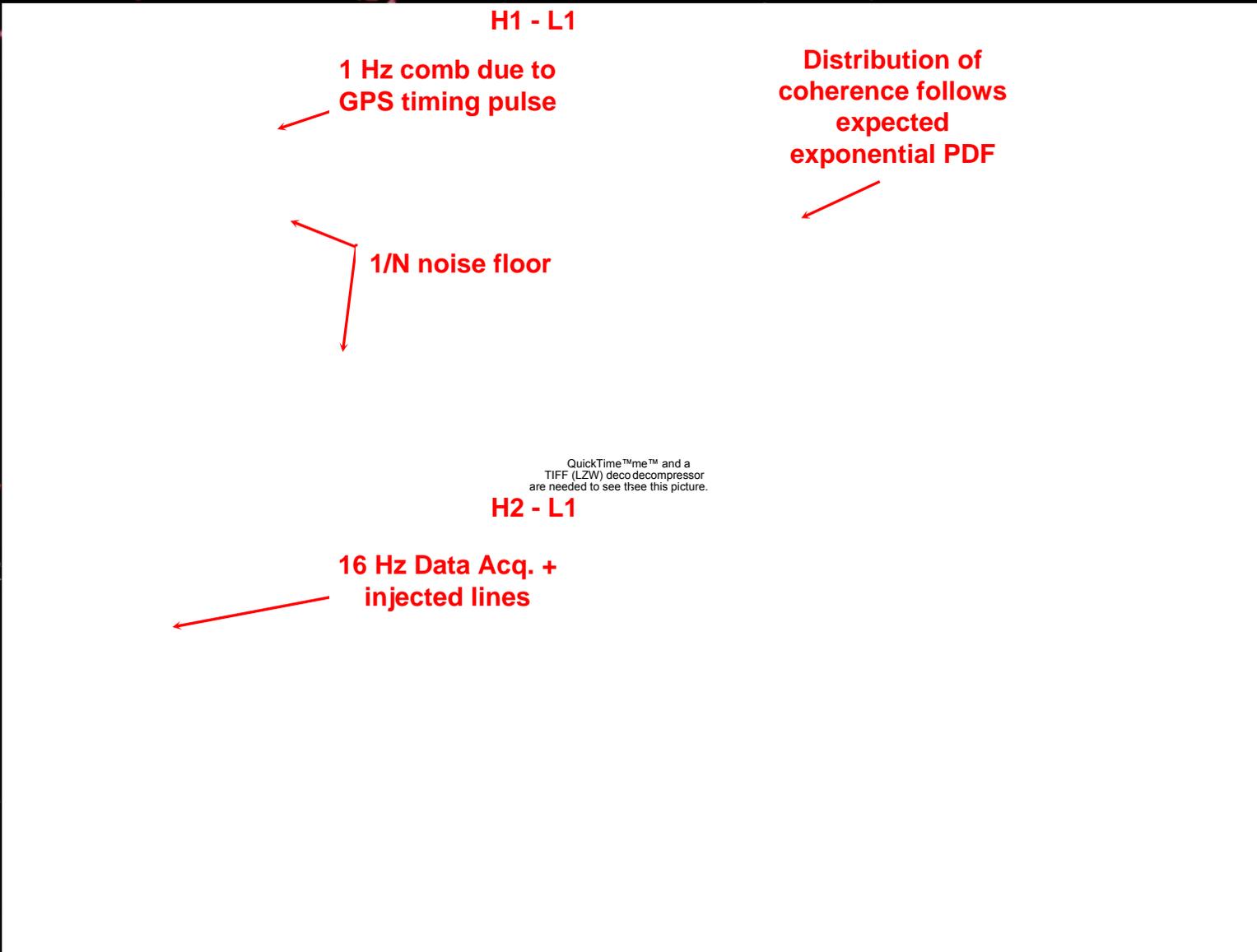
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TIFF (LZW) decompressor
are needed to see this picture.



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*improves sensitivity reach over
individual noise floors.*

Run Averaged Coherences: A measure of intrinsic sensitivity

$$\Gamma(f) \equiv \frac{|P_{12}(f)|^2}{P_1(f)P_2(f)}$$



Histogram of normalized measurement residuals

- Make a measurement every 192s
- Look at statistics of measurements to determine whether they are Gaussian

\tilde{r}_i

$$T_{obs} = 353.9 \text{ hr}$$

$$N_{seg} = 12637$$

(after cuts)

QuickTime™ and a
TIFF (LZW) decompressor
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- Scheduled for publication in ApJ 658 (March 2007)

H1 - L1

Evolution of measurement over time

Optimally combined spectrum: H1 - L1 & H2 - L1

H1 - L1: 353.9 hr; $N_{seg} = 12,637$
 H2 - L1: 332.7 hr; $N_{seg} = 11,849$

Standard deviation of measurement
 Modulation of envelope due to $\gamma(f)$

90% Bayesian upper limit vs. spectral index, α

$$\Omega(f) \equiv \Omega_\alpha \left(\frac{f}{100\text{Hz}} \right)^\alpha$$

S4 result: $\Omega_0 < 6.5 \times 10^{-5}$ (90% U.L.)

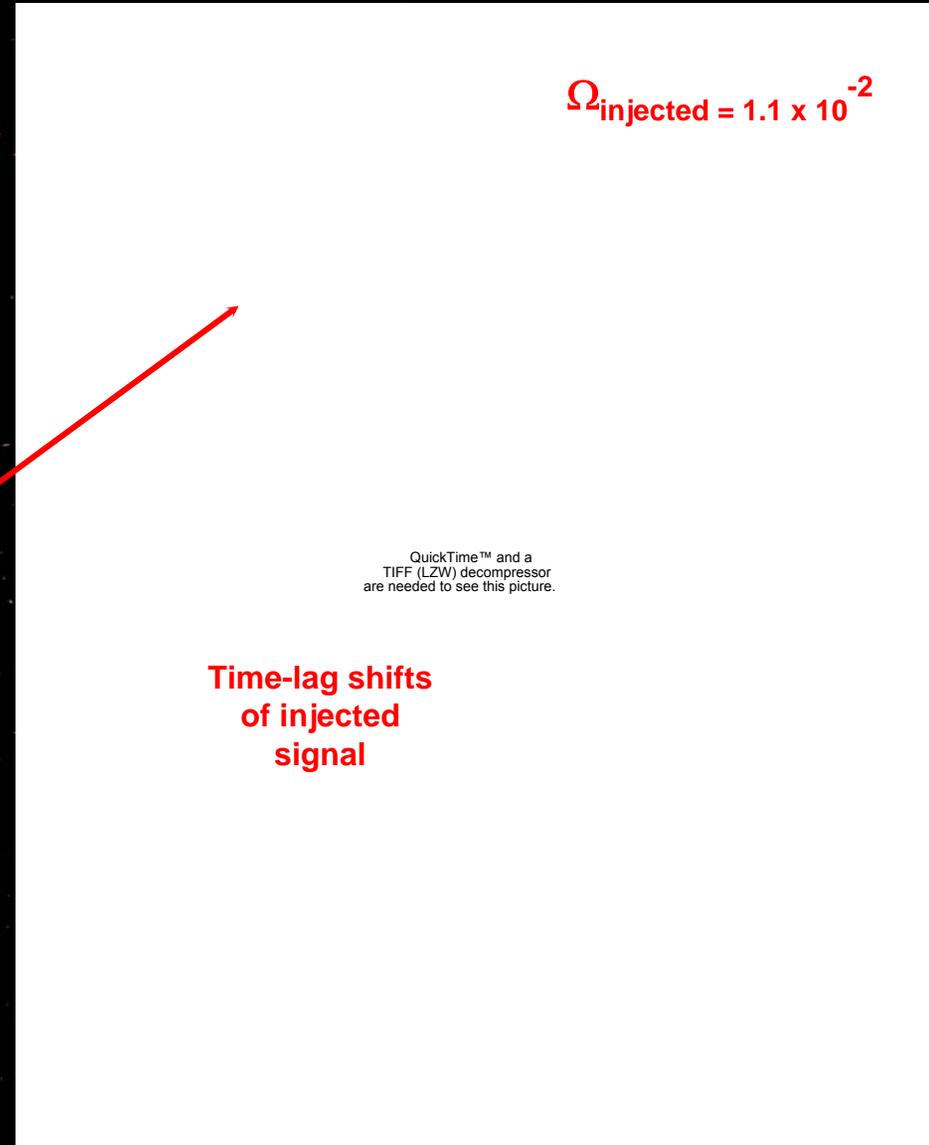
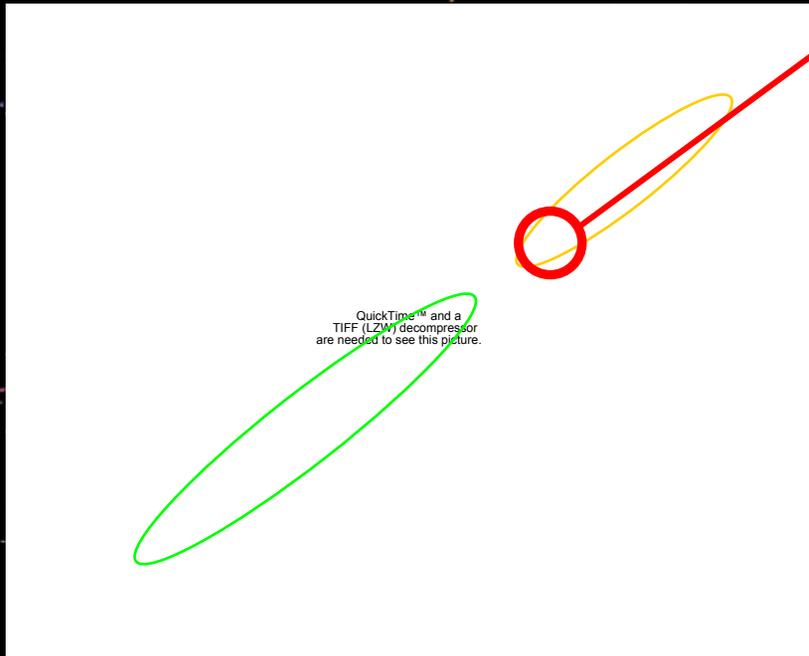
$$\Omega_{expt,S4} = \int Y(f) df = -8 \times 10^{-6}$$

$$\sigma_{expt,S4} = \sqrt{\int \text{sigma}_Y^2(f) df} = 4.3 \times 10^{-5}$$

QuickTime™ and a TIFF (LZW) decompressor are needed to see this picture.

Tests with signal injections Hardware & Software

- HW: Introduce coherent excitation of test masses (WA - LA) with a spectrum simulation a constant $\Omega_{\text{GW}}(f)$ background with different strengths
- SW: Simulated signal added to strain signal



The Stochastic GW Landscape

$H_0 = 72 \text{ km/s/Mpc}$

Armstrong,
ApJ 599 (2003) 806

Smith et al.,
PRL 97 (2006) 021301

Kolb & Turner,
The Early Universe (1990)

Jenet et al.,
ApJ 653 (2006) 1571

QuickTime™ and a
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Allen&Koranda,
PRD 50 (1994) 3713

Damour & Vilenkin
PRD 71 (2005) 063510

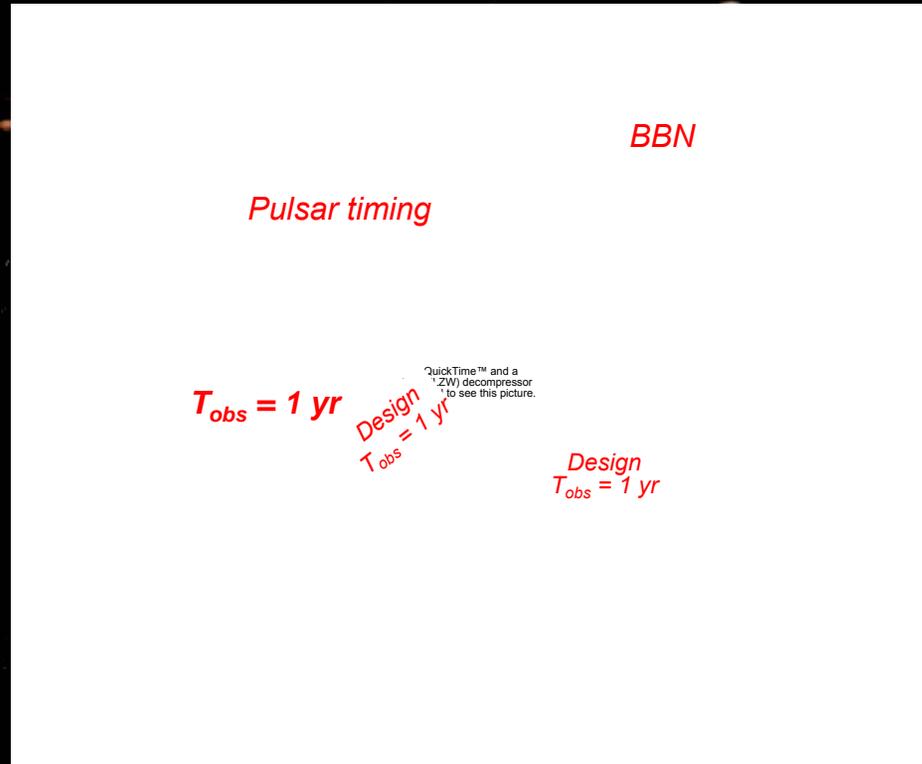
Buonanno et al.,
PRD 55 (1997) 3330

Gasperini & Veneziano,
Phys. Rep. 373 (2003) 1

Turner, PRD 55 (1997) 435

S4 all-sky observations

Implications for Cosmic strings



- Model of Damour & Vilenkin
- p : probability for reconnection
- $G\mu$: string tension - regime accessible by LIGO
- ϵ : loop size - regime accessible by LIGO

Spatially-resolved search Time-dependent Detection Strategy

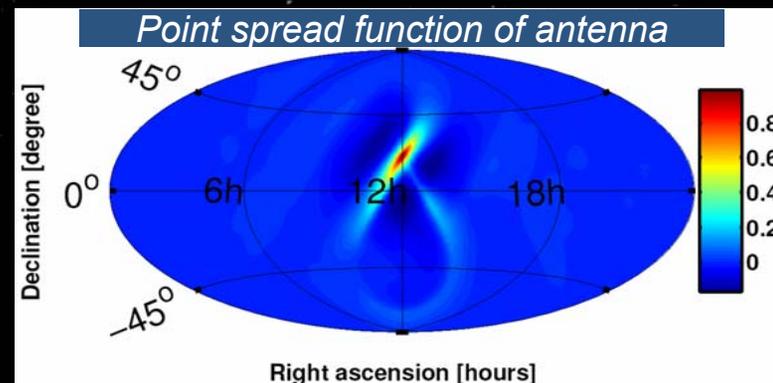
- Time-dependent overlap reduction function tracks a point in the sky over the sidereal day:

$$\gamma_{point}(f, t : \hat{\Omega}) \equiv \frac{1}{2} \sum_A e^{i2\pi f \hat{\Omega} \cdot \frac{\Delta \vec{x}}{c}} F_1^A(\hat{\Omega}) F_2^A(\hat{\Omega})$$

- Time dependent optimal filter enables spatially resolved measurement:

$$\tilde{Q}(f, t : \hat{\Omega}) = \frac{\gamma(f, t : \hat{\Omega}) \Omega_{gw}^{model}(f)}{f^3 P_1(f) P_2(f)}$$

- Characteristic size of point spread function:
 - λ/D
 - $\sim 11^\circ$ @ 500 Hz for $D = 3000$ km
 - Diffraction-limited GW astronomy



Spatially resolved search with S4 Simulations & Validation

Detection of an injected hardware pulsar simulation



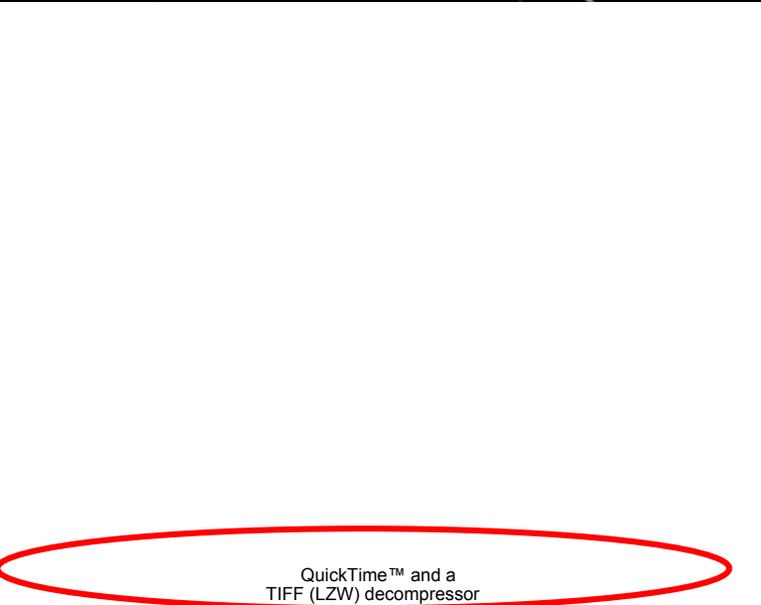
Pulsars Hardware Injection			
Parameter	Pulsar3	Pulsar4	Pulsar8
Freq. during S4	108.86 Hz	1402.20 Hz	193.94 Hz
Estimate Hdf on source	1.74×10^{-46}	4.05×10^{-44}	1.79×10^{-46}
Error bar	1.89×10^{-47}	6.04×10^{-46}	1.73×10^{-47}
SNR	9.2	67.1	10.3
inj. Hdf	1.74×10^{-46}	4.28×10^{-44}	1.54×10^{-46}



Upper limit sky maps

No detection of signal at S4 sensitivity

Submitted to PRD;
arXiv:astro-ph/0701877v1



- $\Omega_{GW}(f) \sim \text{const}$ Results
- Integration over sky yields sky-averaged result that is consistent with the all-sky technique within measurement errors
- Distribution of signal consistent with Gaussian PDF with 100 DOFs (# of independent sky patches)
- **90% Bayesian UL:**
 $\Omega_{GW} < 1.02 \times 10^{-4}$



- $\Omega_{GW}(f) \sim f^3$ Results
- Distribution of signal consistent with Gaussian PDF with 400 DOFs (# of independent sky patches)
- **90% Bayesian UL:**
 $\Omega_{GW}(f) < 5.1 \times 10^{-5} (f/100\text{Hz})^3$





Use spatially resolved stochastic search technique to look for periodic emissions of unknown f : Sco-X1

Template-less search: use signal in one detector as the "template"

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Summary - S4 results

- All-sky measurement: $\Omega_0 h_{72}^2 < 6.5 \times 10^{-5}$
 - 13X improvement over previous (S3) result
 - Still weaker than existing BBN limit
 - S4 (and previous S3) are starting to explore, restrict parameter space of some stochastic models, such as cosmic strings and pre-big bang
 - The S5 data analysis is in progress -- should beat the BBN limits for models in which signal is concentrated in the LIGO band
- Spatially resolved measurements:
 - New technique that exploits phased-array nature of the LIGO site pairs to steer beam and to track sky positions
 - All-sky result follows as a subset of the measurements
 - Approach can be used look for a number of astrophysical foregrounds, by changing the optimal filter to match the source properties
 - Work ongoing to (i) implement deconvolution of antenna point spread function from raw maps; (ii) decompose map into spherical harmonics basis functions, analogous to CMB maps.
- Expected sensitivities with one year of data from LLO-LHO:
 - Initial LIGO $\Omega_0 h^2 < 2 \times 10^{-6}$
 - Advanced LIGO $\Omega_0 h^2 < 7 \times 10^{-10}$

FINIS

Intrinsic Sensitivity: σ_Y Trend for Entire Run

$$\sigma_Y^2 \approx \frac{T}{2} \int_0^\infty df P_1(f) P_2(f) |Q(f)|^2$$

\tilde{S} / \tilde{I}

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