

# Coating Design Optimization for Advanced Interferometers : Minimizing the Total Noise Budget

**Giuseppe Castaldi,**  
**Vincenzo Galdi,**  
**Vincenzo Pierro,**  
**Innocenzo M. Pinto,**  
**Riccardo de Salvo**  
**Juri Agresti**

***TWG, University of Sannio at Benevento***  
***Caltech, LIGO-Lab***

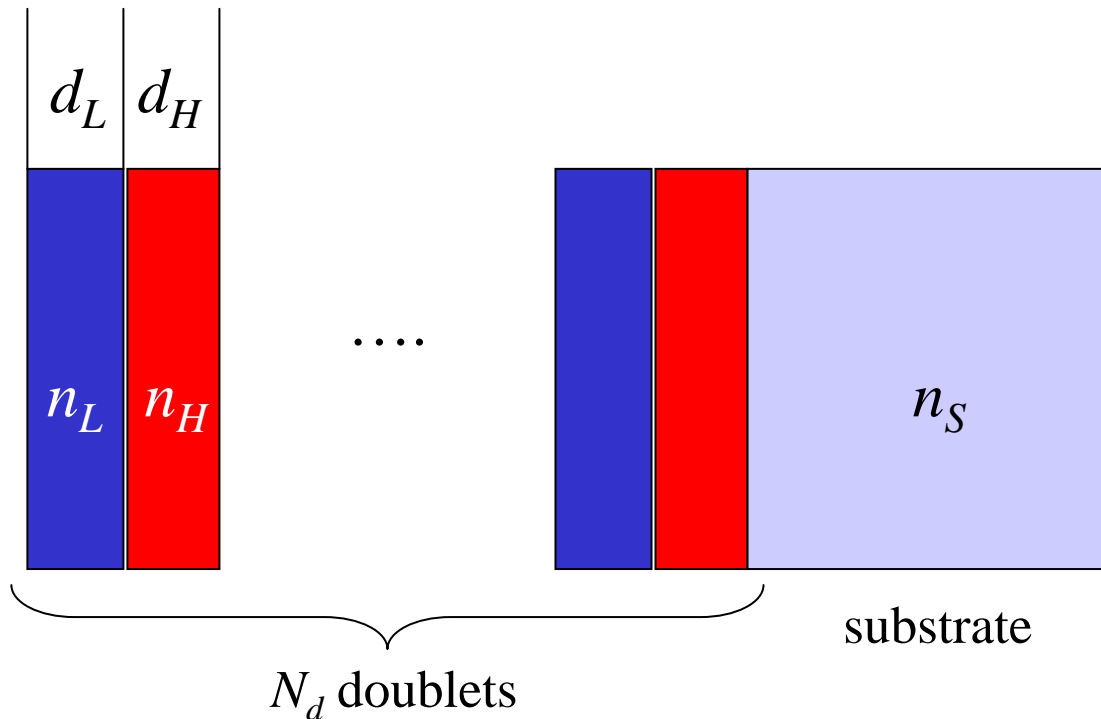




# Outlook



- 
- Iso-reflective Stacked-Doublet Coatings;
  - Minimizing the Coating Brownian Noise;
  - Doped Tantalum;
  - Total Coating Noise Budget Ingredients;
  - Results;
  - Conclusions.



$$d_{L,H} = \left( \frac{\lambda_0}{n_{L,H}} \right) z_{L,H}$$

$$z_{L,H} = \frac{1}{4} \pm \xi$$

For

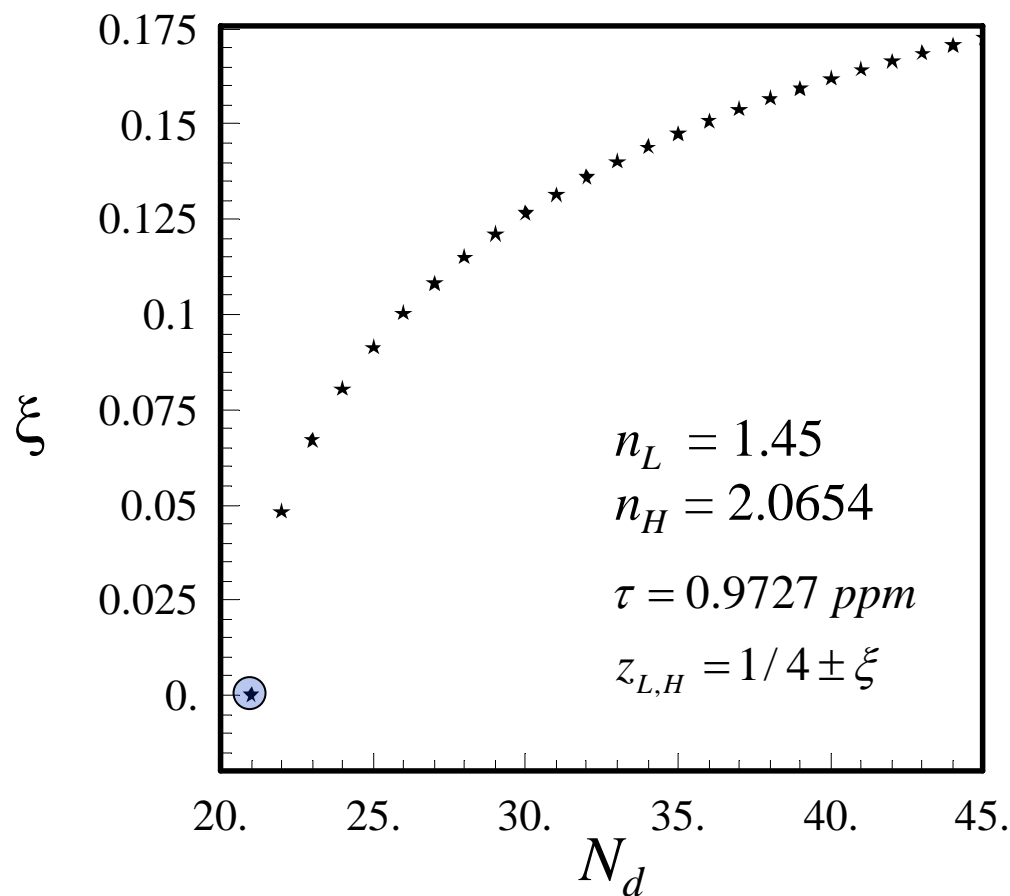
$$n_L = 1.45$$

$$n_H = 2.0654$$

the QWL ( $\xi=0$ ) design  
which goes closest to  
the 1ppm Adv LIGO  
design goal has

$$N_D = 21$$

$$\tau = 0.9727 \text{ ppm}$$



$$S_{\Delta x}^{(B)}(f) = \frac{\overset{\text{Boltzmann}}{\sqrt{2}k_B T} \overset{\text{Poisson ratio}}{(1 - \nu_s^2)}}{\pi^{3/2} f \underset{\text{Beam spot radius}}{r_0} \underset{\text{Young modulus}}{E_s}} \overset{\text{Coating loss angle}}{\phi_c}, \quad \phi_c = N_d(b_L z_L + b_H z_H)$$

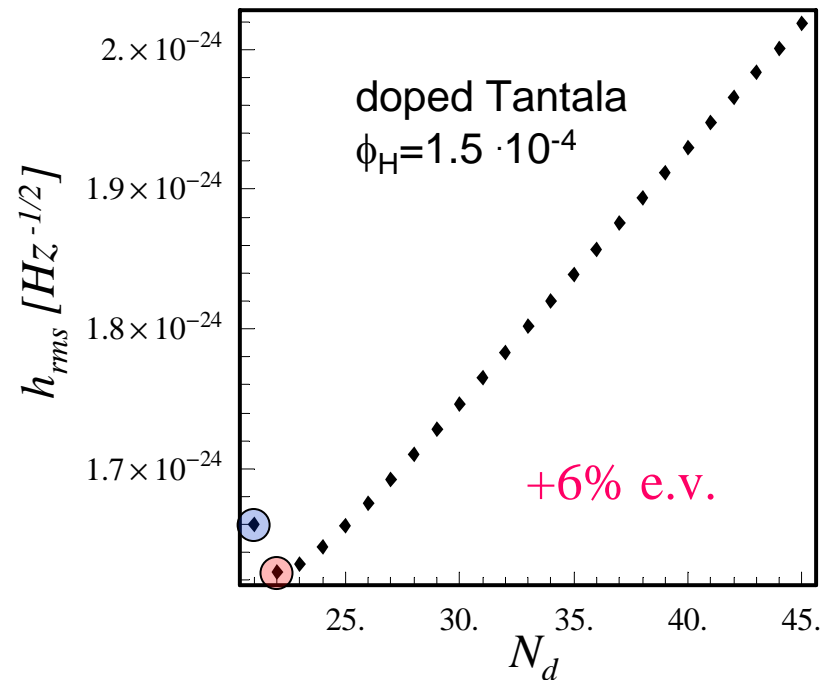
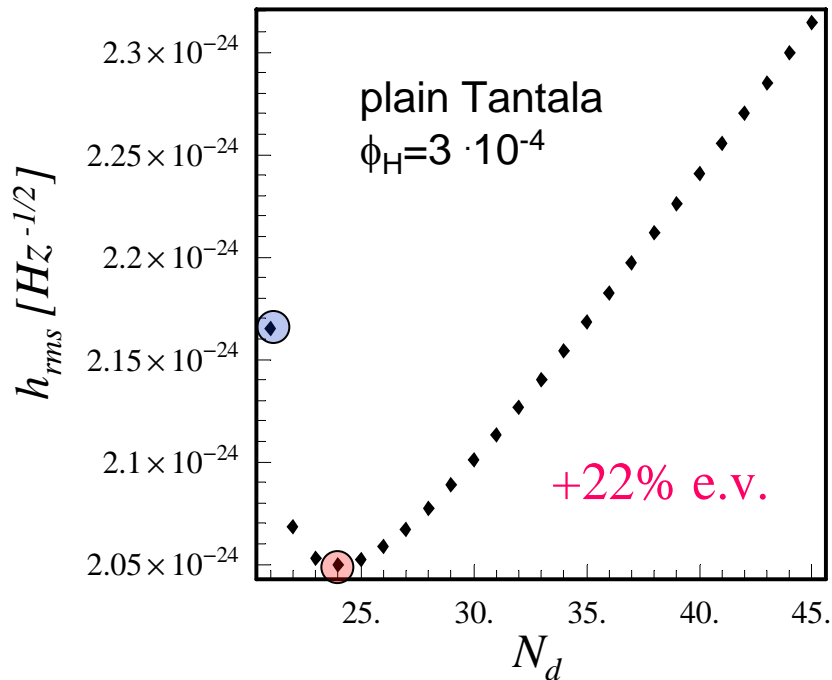
$$b_{L,H} \approx \frac{\lambda_0}{\sqrt{2\pi} r_0} \frac{\phi_{L,H}}{n_{L,H}} \left( \frac{E_{L,H}}{E_s} + \frac{E_s}{E_{L,H}} \right) \quad \nu_{L,H} \ll 1$$

$$b_H/b_L = 5.149 \quad [\text{Tantala (plain) - Silica coatings}]$$

# Plain vs. Doped Tantalum Brownian Noise Only

Brownian Noise Only.  $\tau = 0.9727\text{ppm}$ ,  $f = 100\text{Hz}$

● QWL coating ● Optimized coating



Effective fluctuations of the test-mass (coated mirror) front - face position with respect to the mirror center of mass may occur as an effect of

- Thermal expansion of the coating layers (thermoelastic effect),

$$\Delta x^{(TE)} = \alpha_{eff} d_{tot} \Delta T$$

effective coating expansion coeff.      coating thickness

- Thermal variations of the refraction indexes  $n_{H,L}$  of the coating materials (thermorefractive effect),

$$\Delta x^{(TR)} = \beta_{eff} \lambda_0 \Delta T$$

thermorefractive coefficient      optical wavelength (vacuum)

Power spectral density (PSD) :

Wiener-Khinchin th.

$$S_{\Delta x}(f) = \mathcal{F}_{\tau \rightarrow f} \langle \Delta x(t) \Delta x(t + \tau) \rangle_t = \left( \frac{\Delta x}{\Delta T} \right)^2 \mathcal{F}_{\tau \rightarrow f} \langle \Delta T(t) \Delta T(t + \tau) \rangle_t =$$

$$= \left( \frac{\Delta x}{\Delta T} \right)^2 S_{\Delta T}(f)$$

PSD of T - fluctuations in the coating

$$S_{\Delta T}(f) = S_{\Delta T}^{(\Theta)}(f) + S_{\Delta T}^{(\Phi)}(f)$$

Intrinsic fluctuations of thermodynamic origin

add in-coherently

Photo-thermal fluctuations arising from laser shot noise through optical absorption



$$S_{\Delta T}^{(\Theta)}(f) = \frac{k_B T^2}{\pi^{3/2} r_0^2 \sqrt{f \kappa_s C_s \rho_s}}$$

[V. Braginsky, Phys. Lett A264 (1999) 1]

single photon energy  
power abs. in coating

mass density  
specific heat capacity  
thermal conductivity } of substrate

$$S_{\Delta T}^{(\Phi)}(f) = \frac{P_{abs} E_\lambda}{4\pi^3 r_0^4 \kappa_s \rho_s C_s f}$$

[S. Rao, PhD Thesis, Caltech, 2003, etd-05092003-153759]

$$E_\lambda \cong 1.867 \cdot 10^{-19} J @ \lambda = 1064 nm$$

$$P_{abs} = 0.4 W \text{ for Adv LIGO}$$

(a different formula for  $S_{\Delta T}^{(\Phi)}$   
applies for sapphire substrates)



# Coating Thermal Noise Budget

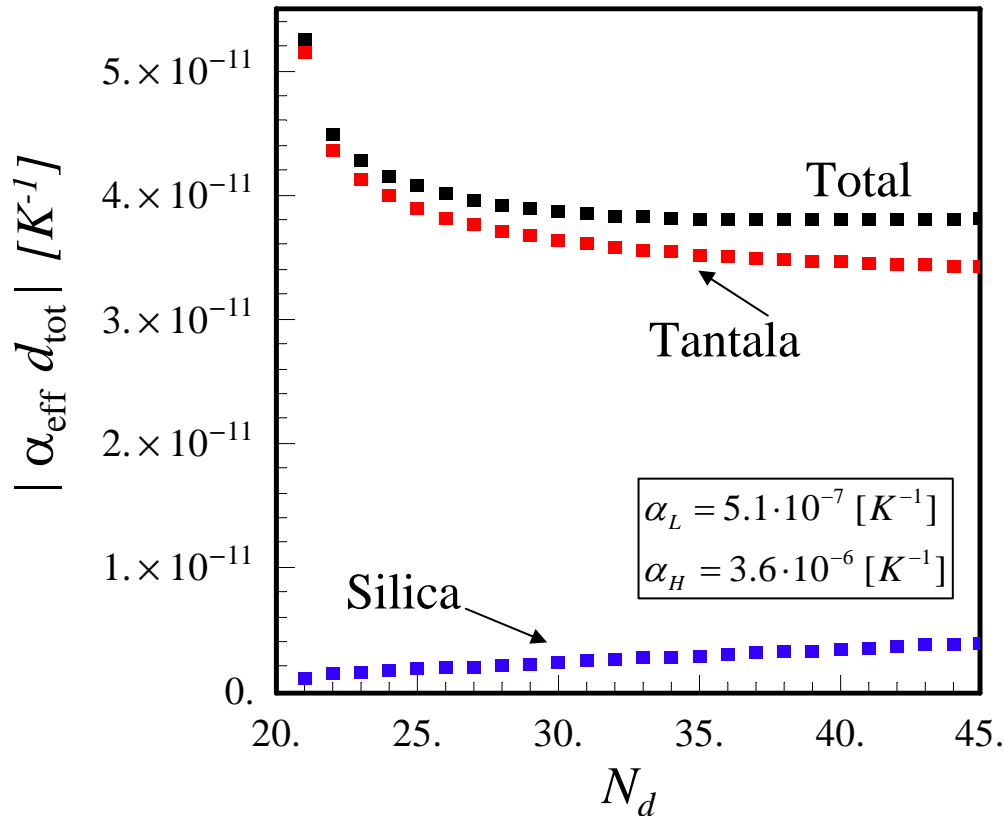


## Total Coating Noise PSD

$$S_{\Delta x}^{(tot)}(f) = S_{\Delta x}^{(B)}(f) + \left( \frac{\Delta x^{(TE)}}{\Delta T} + \frac{\Delta x^{(TR)}}{\Delta T} \right)^2 S_{\Delta T}(f)$$

The thermal - driven elastic and refractive fluctuations should add *coherently*. Indeed, the temperature in the coating does *not* fluctuate

- on the space-scale (thickness) of the coating,
- on the time scales whereby the field in the coating builds up.



[V.B. Braginsky and S.A. Vyatchanin, Phys. Lett. A312 (2003) 244; idem, cond-mat/0302617 contains important corrections]

[M. Fejer et al., PRD-70 (2004) 082003]

General formula available, OK for general SD coatings (also in the form of linear combination of  $z_L, z_H$ )

$$\bar{Y}_{in} = \bar{Y}_{in}^{(0)} + \Delta\bar{Y}_{in} \quad (\text{photorefractive change in coating input admittance})$$

$$\Gamma_{in} = \frac{1 - \bar{Y}_{in}}{1 + \bar{Y}_{in}} \approx \frac{1 - \bar{Y}_{in}^{(0)}}{1 + \bar{Y}_{in}^{(0)}} \left( 1 - \frac{2\Delta\bar{Y}_{in}}{1 - (\bar{Y}_{in}^{(0)})^2} \right) = \Gamma^{(0)} \left( 1 - \frac{2\Delta\bar{Y}_{in}}{1 - (\bar{Y}_{in}^{(0)})^2} \right)$$

$$\Gamma(\Delta x) = \Gamma(0) \exp \left[ i \frac{4\pi}{\lambda_0} \Delta x \right] \approx \Gamma(0) \left[ 1 + i \frac{4\pi}{\lambda_0} \Delta x \right] \quad (\text{Transport equation for reflection coeff.})$$

$$\Delta x = \beta_{eff} \lambda_0$$



$$\beta_{eff} = -\frac{1}{2\pi l} \frac{\Delta\bar{Y}_{in}}{1 - (\bar{Y}_{in}^{(0)})^2}$$

- High reflectivity coatings, thus  $N_d$  *very large*;
- For  $N_d$  very large ( $N_d \rightarrow \infty$ ), addition of a further doublet *does not* change the coating input admittance;
- For this *single added* doublet we accordingly have

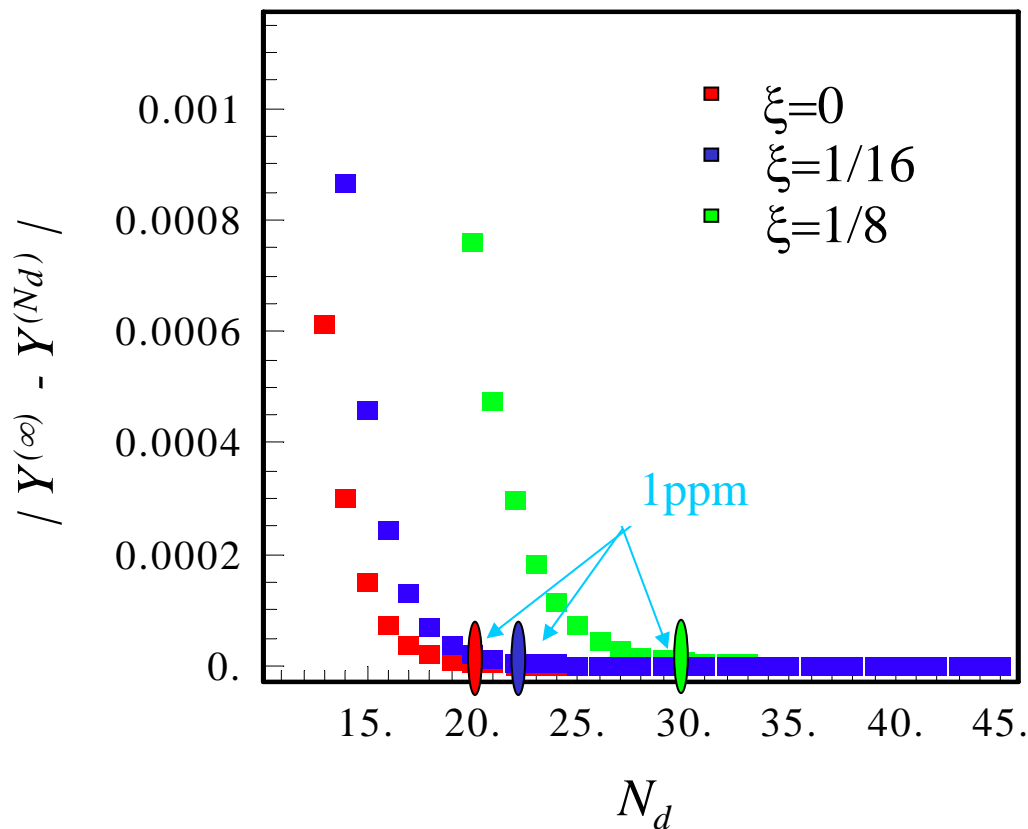
$$\frac{E_{in}}{Z_0 H_{in}} = \bar{Y}_{in} = \frac{E_{out}}{Z_0 H_{out}}$$

combined with single-doublet transmission matrix equation

$$\begin{pmatrix} E_{out} \\ Z_0 H_{out} \end{pmatrix} = \begin{pmatrix} \Theta_{11} & \Theta_{12} \\ \Theta_{21} & \Theta_{22} \end{pmatrix} \cdot \begin{pmatrix} E_{in} \\ Z_0 H_{in} \end{pmatrix}$$

gives an equation in  $Y_{in}$  . **Yields formula "1"** for QWL case.

[V. Galdi and I.M. Pinto, 2007]



$$\beta_{\text{eff}} = \frac{n_H^2 \beta_L + n_L^2 \beta_H}{4(n_L^2 - n_H^2)}$$

[Braginsky, Gorodetsky and Vyatchanin,  
Phys. Lett. A 271 (2000) 303]

- deduced for QWL coatings only
- deduction based on a consistency argument
- $\beta_{\text{eff}}$  does not vanish in the limit  $n_H \rightarrow \infty$

$$\beta_{\text{eff}} = \frac{n_L n_H (\beta_L + \beta_H)}{4(n_L^2 - n_H^2)}$$

[Braginsky and Vyatchanin, Phys. Lett.  
A312 (2003) 244]

- claimed to fix an error in previous formula
- no details given about deduction
- $\beta_{\text{eff}}$  **vanishes in the limit**  $n_H \rightarrow \infty$

$$\begin{pmatrix} E_{in} \\ Z_0 H_{in} \end{pmatrix} = \begin{pmatrix} T_{11}^{(N_d)} & T_{12}^{(N_d)} \\ T_{21}^{(N_d)} & T_{22}^{(N_d)} \end{pmatrix} \cdot \begin{pmatrix} E_{out} \\ Z_0 H_{out} \end{pmatrix}$$

$N_d$ - doublets coating QWL  
transmission matrix

$$T_{11}^{(N_d)} = \left( -\frac{n_H^{(0)}}{n_L^{(0)}} \right)^{N_d} \left[ 1 + \frac{2N_d}{\pi} (\Delta\psi_H - \Delta\psi_L) \right],$$

where:

$$T_{12}^{(N_d)} = i \left( \frac{\Delta\psi_H}{n_L^{(0)}} + \frac{\Delta\psi_L}{n_H^{(0)}} \right) S(N_d),$$

$$\Delta\psi_{L,H} = \frac{\pi}{2} \frac{\beta_{L,H}^{(0)}}{n_{L,H}^{(0)}} \Delta T$$

$$T_{21}^{(N_d)} = -i \left( n_L^{(0)} \Delta\psi_H + n_H^{(0)} \Delta\psi_L \right) S(N_d),$$

$$T_{22}^{(N_d)} = \left( -\frac{n_L^{(0)}}{n_H^{(0)}} \right)^{N_d} \left[ 1 - \frac{2N_d}{\pi} (\Delta\psi_H - \Delta\psi_L) \right],$$

(proven by complete induction)

$$S(N_d) = \begin{cases} \sum_{m=-P}^P \left( \frac{n_L^{(0)}}{n_H^{(0)}} \right)^{2m} & N_d = 2P \text{ (even)} \\ - \sum_{m=-P}^{P-1} \left( \frac{n_L^{(0)}}{n_H^{(0)}} \right)^{2m+1} & N_d = 2P + 1 \text{ (odd)} \end{cases}$$



$$E_{out} = n_S^{-1} Z_0 H_{out} \implies \bar{Y} = \frac{Z_0 H_{in}}{E_{in}} = \frac{T_{21}^{(N_d)} + n_S T_{22}^{(N_d)}}{T_{11}^{(N_d)} + n_S T_{12}^{(N_d)}}$$

$$\bar{Y} = n_S \left( \frac{n_L}{n_H} \right)^{2N_d} \left[ 1 - \frac{4N_d}{\pi} (\Delta\psi_H - \Delta\psi_L) \right]$$

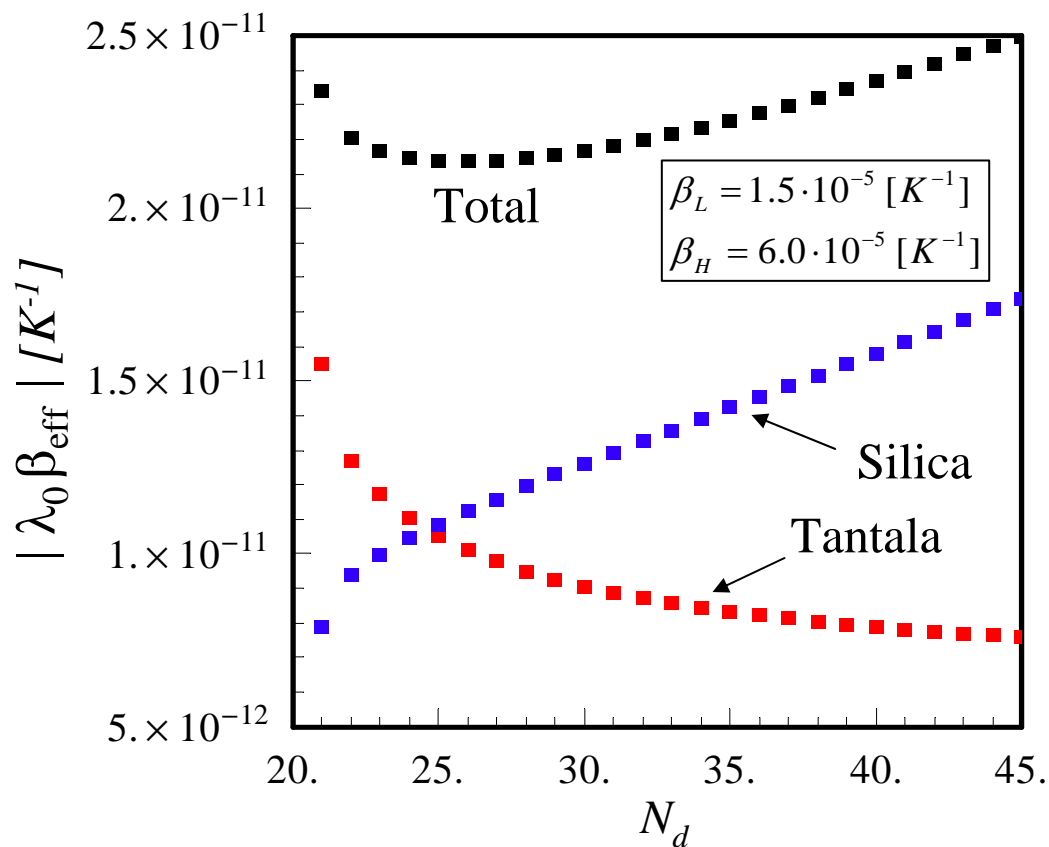
$$-i \frac{n_L^2 n_H \Delta\psi_H + n_H^2 n_L \Delta\psi_L}{n_L^2 - n_H^2}$$

returns Braginsky's formula "1"

$$+i \frac{(n_L^2 + n_S^2) n_H \Delta\psi_H + (n_H^2 + n_S^2) n_L \Delta\psi_L}{n_L^2 - n_H^2} \left( \frac{n_L}{n_H} \right)^{2N_d}$$

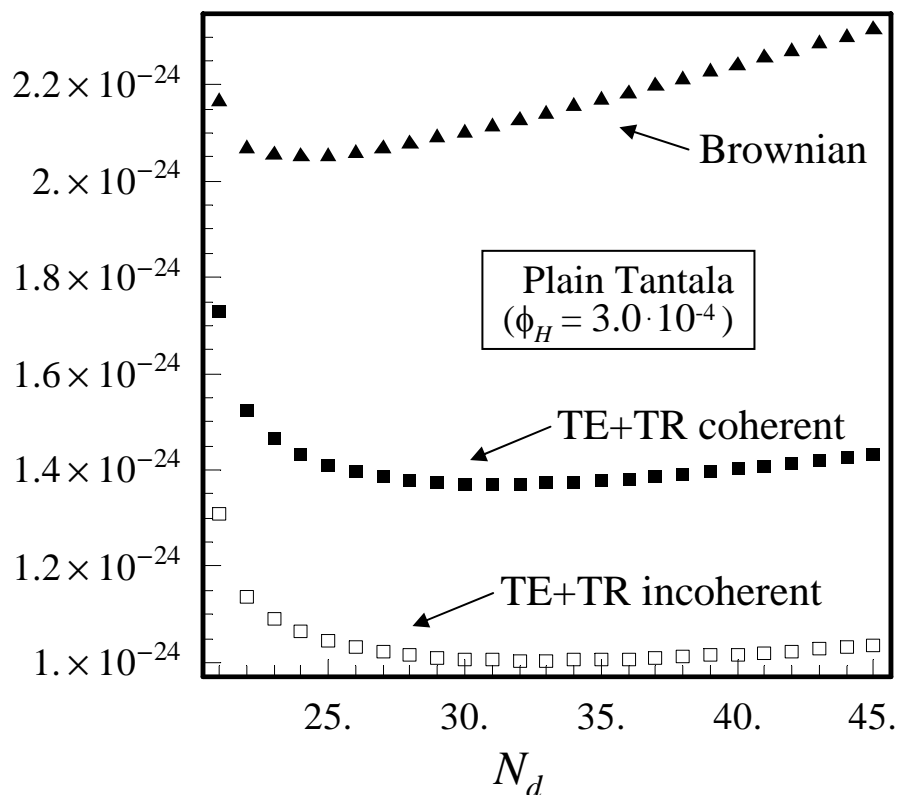
$$-i \frac{n_S^2 (n_H \Delta\psi_H + n_L \Delta\psi_L)}{n_L^2 - n_H^2} \left( \frac{n_L}{n_H} \right)^{4N_d}$$

These terms vanish  
as  $N_d \rightarrow \infty$ , since  $n_H > n_L$ .

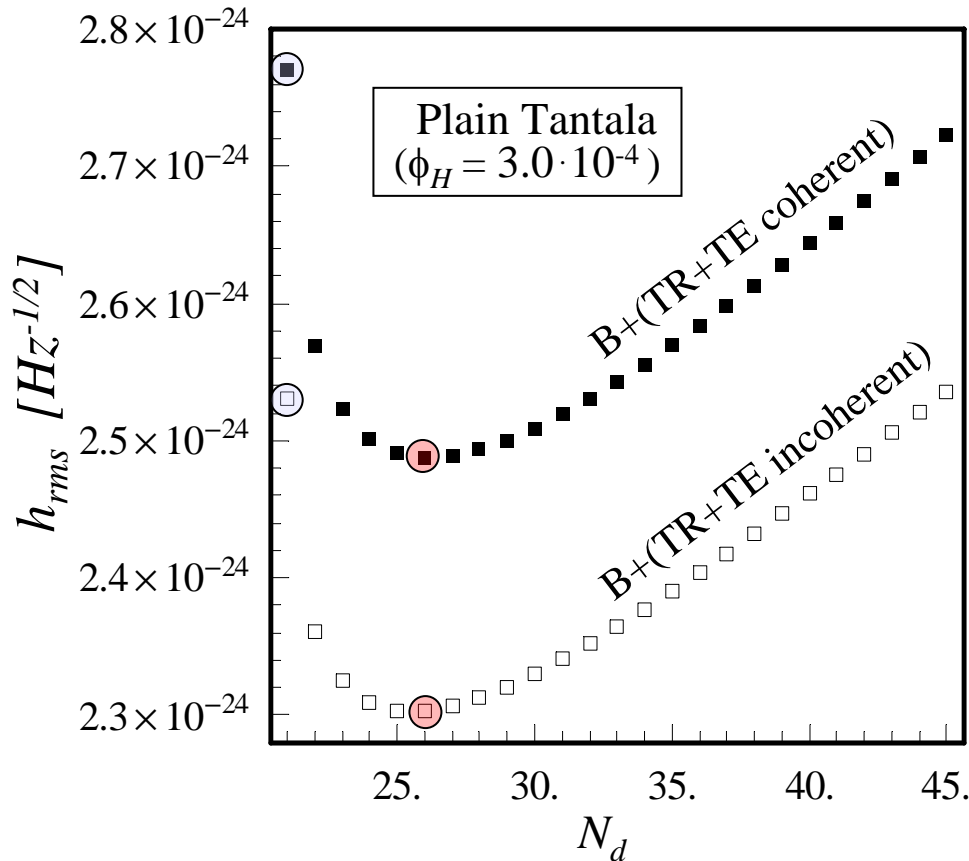


A non - QWL, minimum thermorefractive coefficient stacked doublet design exists, featuring the lowest combination of the low - high index material contributions

$$\tau = 0.9727 \text{ ppm}, f = 100 \text{ Hz}$$

$$h_{rms} [\text{Hz}^{-1/2}]$$


When using *plain* Tantala, in the standard QWL design, the *total* (TE+TR, coherent) non-Brownian coating noise term is *comparable* to the Brownian one...



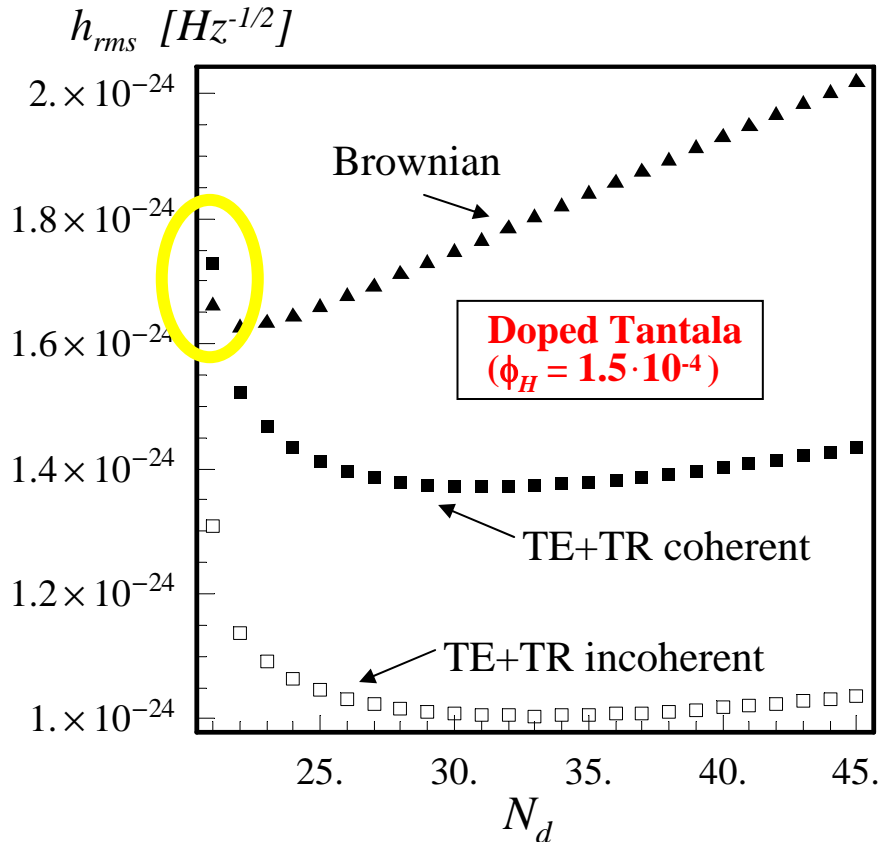
$$\tau = 0.9727 \text{ ppm}, f = 100 \text{ Hz}$$

● QWL coating

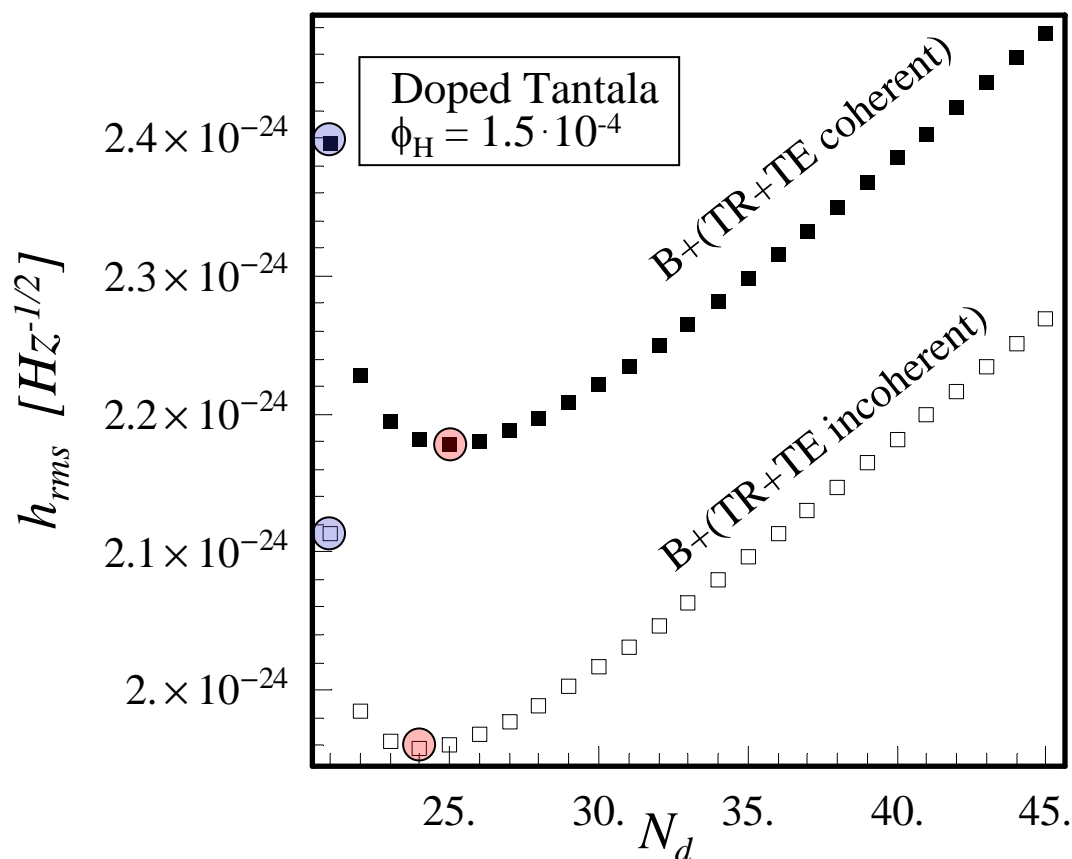
● Optimized coating

...The optimal SD coating design is *distinctly different* from QWL, and the related event rate boost is sensible (+ 38%)...

$$\tau = 0.9727 \text{ ppm}, f = 100 \text{ Hz}$$



When using **doped** Tantalum, in the standard QWL design, the **total** (TE+TR, coherent) non-Brownian coating noise term turns out to be **larger than** the Brownian one...



$$\tau = 0.9727 \text{ ppm}, f = 100 \text{ Hz}$$

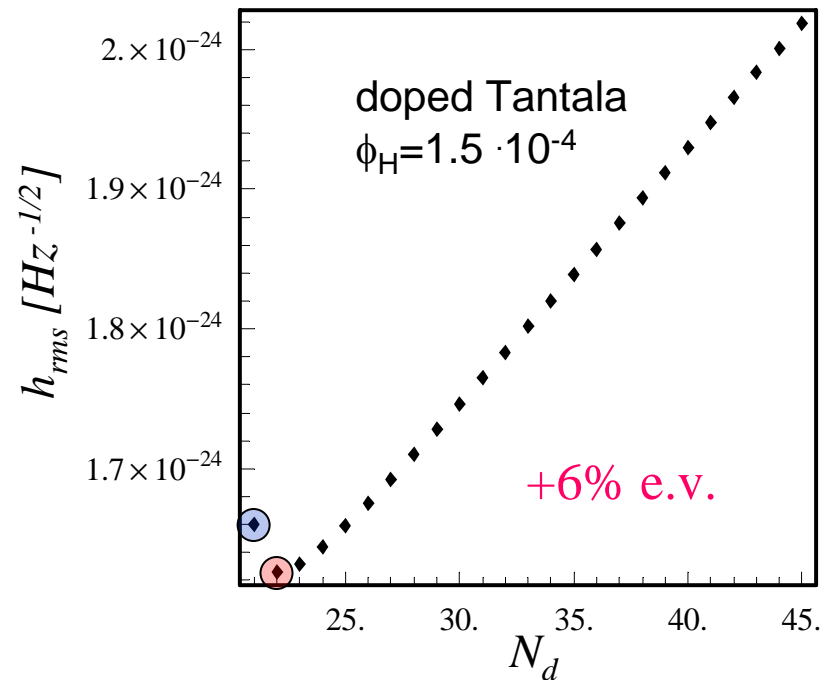
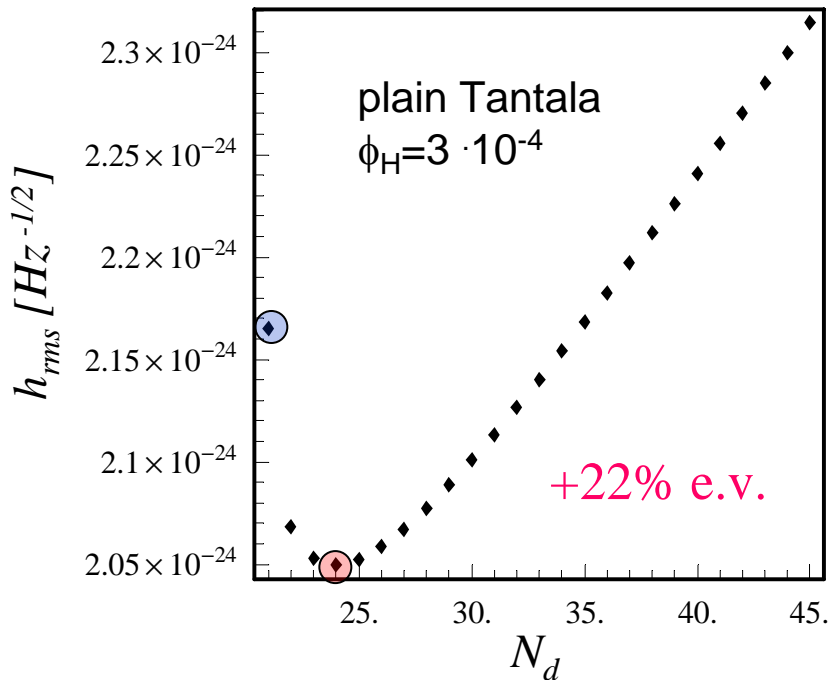
● QWL coating

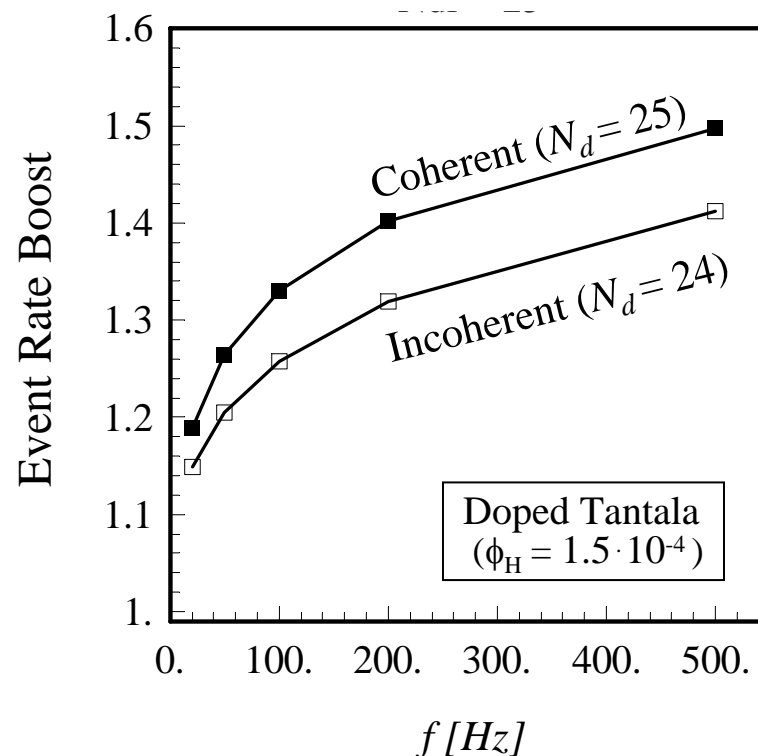
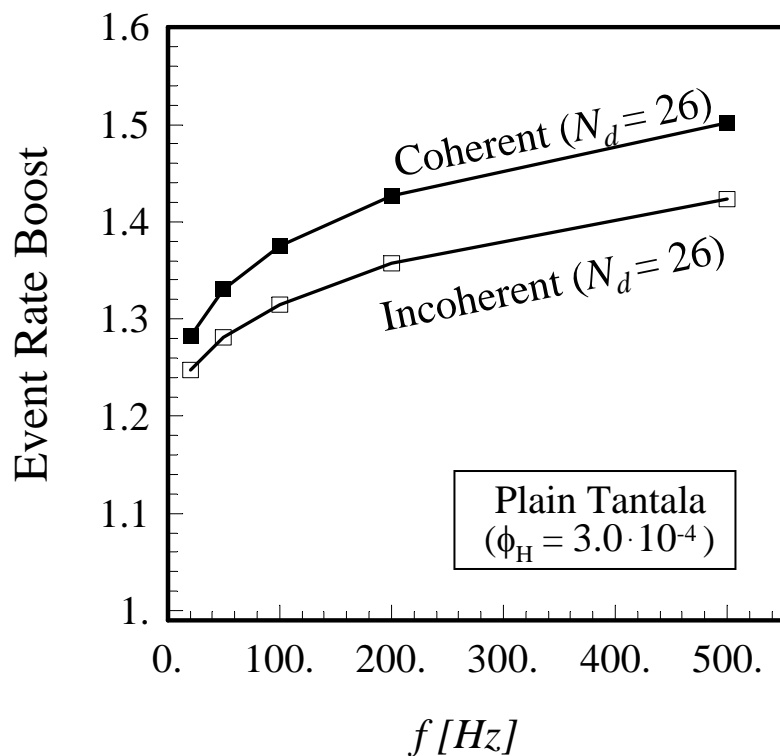
● Optimized coating

...The optimal SD coating design is *distinctly different* from QWL, and the related event rate boost is sensible (+ 33%)...

Brownian Noise Only.  $\tau = .9727\text{ppm}$ ,  $f = 100\text{Hz}$

● QWL coating ● Optimized coating









# Event Rate Boost (Isotropic/Homogeneous Source Distribution)



(Total Noise Budget))

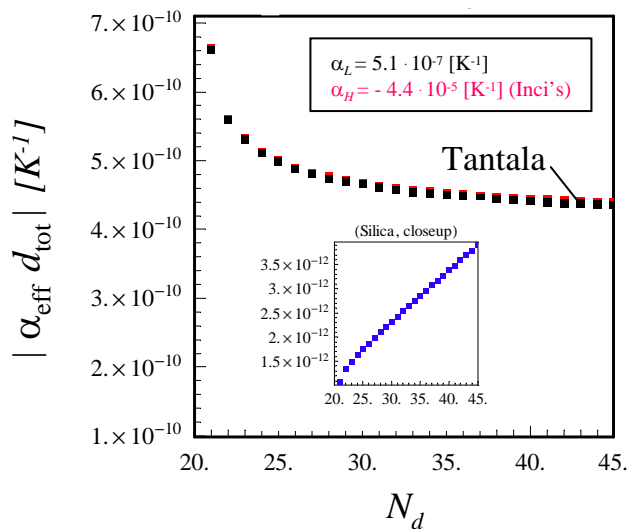
$\tau = 0.9727 \text{ ppm}$	Event rate boost@100Hz		
Plain Tantala, QWL	1		
Plain Tantala, OPT	1.38	1	
Doped Tantala, QWL	1.54	1.11	1
Doped Tantala, OPT	2.05	1.48	1.33

- Coating thickness optimization should be considered as *almost mandatory* to minimize coating noise, even more so when using doped Tantalum, yielding in all cases a *substantial increase* ( $> 30\%$  @  $100\text{Hz}$ ) in the expected event rate, as compared to the QWL design.
- Among all proposed coating noise reduction techniques (new materials, cryogenic mirrors, flat-top beams) thickness optimization offers a *cheapest reliable addition*
- Coating thickness optimization has been shown to be *effective* in reducing the total coating noise even when using the controversial Inci's values for  $\alpha_H$ ,  $\beta_H$ .
- Optimized coating prototypes are scheduled for testing at Caltech TNI; and (still) waiting for delivery from LMA.

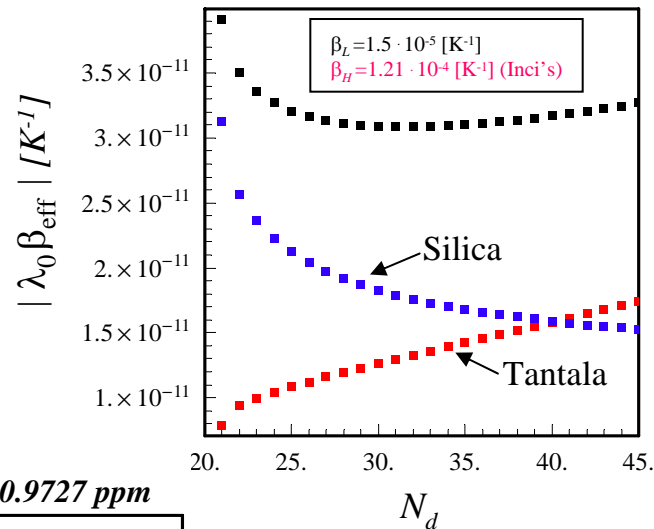
---

Special thanks are due to Gregg Harry (LIGO – LSC and MIT), Andri Gretarsson (LIGO - LSC and Embry Riddle Aeronautical University, Prescott AZ, and Sheila Rowan (LIGO – LSC and University of Glasgow), for many useful hints during the development of this work.

Work sponsored in part by INFN through a Group-V grant (COAT 2006, 2007)



(TE Coefficient)



(TR Coefficient)

