

# Observing Gravitational Waves from Spinning Neutron Stars

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for the LIGO Scientific Collaboration

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# Outline

- 1 **Astrophysical Motivation**
  - Gravitational Waves from Neutron Stars?
  - Emission Mechanisms (Mountains, Precession, Oscillations, Accretion)
  - Gravitational Wave Astronomy of NS
- 2 **Detecting Gravitational Waves from NS**
  - Status of LIGO (+GEO600)
  - The Data-analysis Problem
  - Observational Results

# Orders of Magnitude

Quadrupole formula (Einstein 1916).

GW luminosity ( $\epsilon$ : deviation from axisymmetry):

$$\mathcal{O}(10^{-53}) \mathcal{L}_{\text{GW}} \sim \left(\frac{G}{c^5}\right) \epsilon^2 \left(\frac{M V^3}{R}\right)^2$$

$$\mathcal{O}(10^{59}) \frac{\text{erg}}{\text{s}} = \frac{c^5}{G} \epsilon^2 \left(\frac{R_s}{R}\right)^2 \left(\frac{V}{c}\right)^6$$

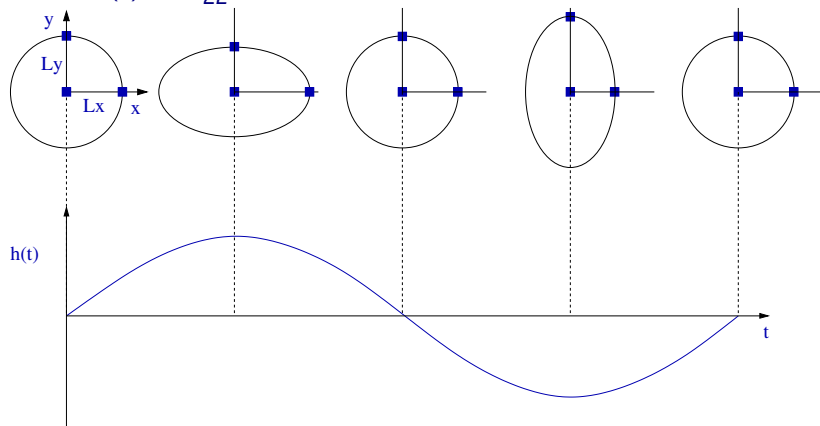
Schwarzschild radius  $R_s = 2GM/c^2$

 Need **compact objects** in **relativistic motion**:  
Black Holes, Neutron Stars, White Dwarfs

# Gravitational Wave Strain $h(t)$

Plane gravitational wave  $h_{\mu\nu}^+$  along  $z$ -direction:

Strain  $h(t) \equiv \frac{L_x - L_y}{2L}$ :



# Triaxial Spinning Neutron Stars

Rotating neutron star:

- non-axisymmetric  $\epsilon = \frac{I_{xx} - I_{yy}}{I_{zz}}$

- rotation rate  $\nu$

☞ GW with frequency  $f = 2\nu$

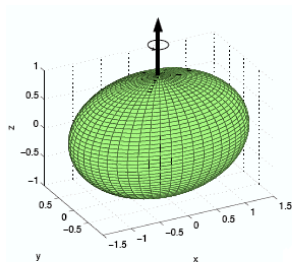
Strain-amplitude  $h_0$  on earth:

$$h_0 = \frac{16\pi^2 G}{c^4} \frac{\epsilon I_{zz} \nu^2}{d}$$

$$= 4 \times 10^{-25} \left( \frac{\epsilon}{10^{-6}} \right) \left( \frac{I_{zz}}{10^{45} \text{ g cm}^2} \right) \left( \frac{\nu}{100 \text{ Hz}} \right)^2 \left( \frac{100 \text{ pc}}{d} \right)$$

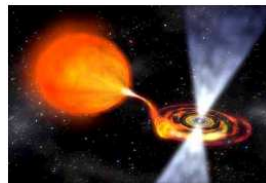
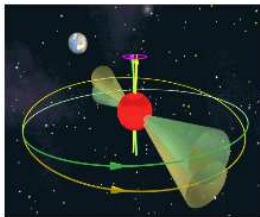
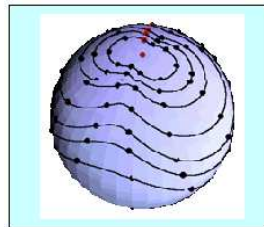
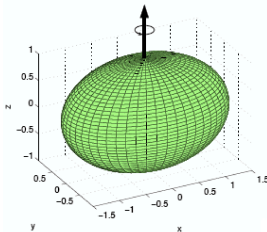
Current LIGO sensitivity (S5):  $\sqrt{S_n} \sim 4 \times 10^{-23} \text{ Hz}^{-1/2}$

☞ NS signals buried in the noise  $\implies$  need “**matched filtering**”



# Possible Emission Mechanisms

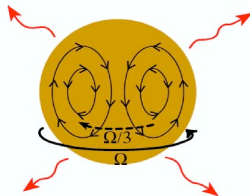
- “Mountains”
- Oscillations
- Free precession
- Accretion (driver)



# Neutron Star “Mountains”

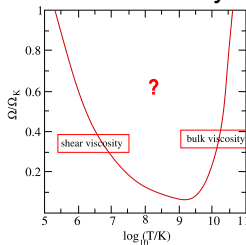
- Conventional NS crustal shear mountains:
  - ☞  $\epsilon_{\text{crust}} \lesssim 10^{-7} - 10^{-6}$  (Ushomirsky, Cutler, Bildsten)
- Superfluid vortices: Magnus-strain deforming crust
  - ☞  $\epsilon_{\text{Magnus}} \sim 5 \times 10^{-7}$  (D.I. Jones; Ruderman)
- Exotic EOS: strange-quark **solid cores**
  - ☞  $\epsilon_{\text{strange}} \lesssim 10^{-5} - 10^{-4}$  (B. Owen)
- Magnetic mountains:
  - large **toroidal** field  $B_t \sim 10^{15}$  G  $\perp$  to rotation:
    - ☞  $\epsilon_{\text{toroidal}} \sim 10^{-6}$  (C. Cutler)
  - accretion along **B**-lines  $\implies$  “bottled” mountains
    - ☞  $\epsilon_{\text{bottle}} \lesssim 10^{-6} - 10^{-5}$  (Melatos, Payne)

# Oscillation Modes



Chandrasekhar-Friedman-Schutz instability:  
counter-rotating mode “dragged forward”  
⇒ **negative** energy and angular momentum  
⇒ emission of GW **amplifies** the mode  
⇒ counteracted by dissipation

r-mode instability window:



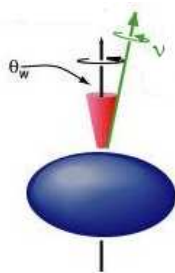
Open questions:

- Dissipation mechanisms: vortex friction, hyperons, crust-core coupling,...
- saturation amplitude, mode-mode coupling, evolution timescales



# Free Precession

“Most general motion of a rigid body” (Landau&Lifshitz 1976)

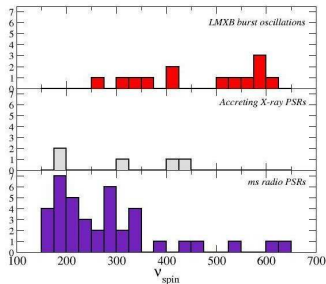


NS are **not** rigid: coupled crust - core  
(viscosity + superfluid vortex pinning)

- likely to be damped rapidly
- no obvious instability or “pumping mechanism”

$$h_0 \sim 10^{-26} \left( \frac{\theta_w}{0.1} \right) \left( \frac{100 \text{ pc}}{d} \right) \left( \frac{\nu}{500 \text{ Hz}} \right)^2$$

# Accretion





Breakup-limit  $\nu_K \sim 1.5$  kHz  $\Rightarrow$  What limits the NS-spin?

Bildsten, Wagoner: Accretion-torque = GW torque ( $\propto \nu^5$ )

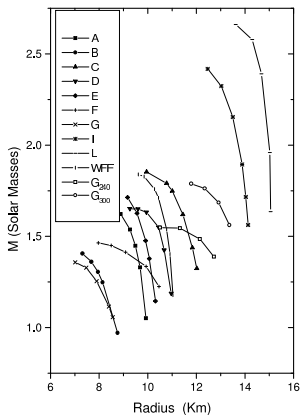
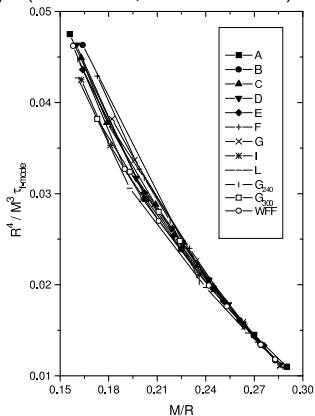
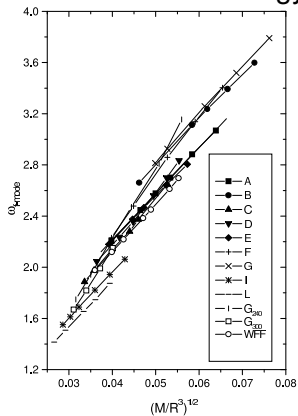
Observed X-ray flux  $\Rightarrow$  Sco X-1:  $h_0 \sim 3 \times 10^{-26} (270 \text{ Hz}/\nu)^{1/2}$

# Astrophysics Summary

- NS are **plausible** sources for GW detection with LIGO I or II
- Whether or not they are **detectable** depends on many poorly-understood aspects of NS physics
-  Any GW-detection from rotating NS will be extremely valuable for NS physics
-  Even the **absence** of detection can yield astrophysically interesting information (crust deformation,  $B_t$ , instabilities)
- NS physics producing GWs is **very different** and **complementary** to electromagnetic emission (bulk-mass motion vs magnetosphere-electron motion)

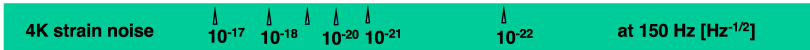
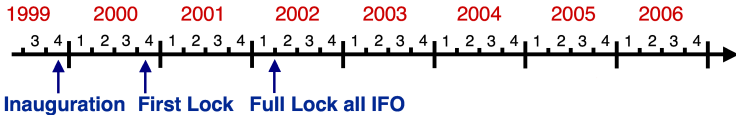
# Gravitational Wave Astronomy

## “Astero-Seismology” (Andersson, Kokkotas 1998): f-mode



👉 Measurement of  $(\omega_f, \tau_f) \Rightarrow$  deduce  $(M, R) \Rightarrow$  EOS

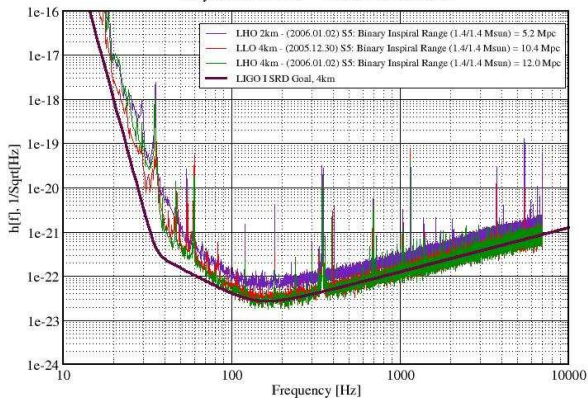
# LSC detectors: LIGO + GEO600



# Current LIGO noise performance

Best Strain Sensivities for the LIGO Interferometers

Early S5 Performance LIGO-G060010-01-Z



$$h_0 = \frac{\Delta L}{L} \sim 3 \times 10^{-23} \implies \Delta L \sim 10^{-19} \text{ m} = 10^{-4} \text{ fm!!}$$

# LSC Data Analysis

LIGO (H1, H2, L1) and GEO600 data analyzed within the  
**LIGO Scientific Collaboration** (LSC):

~ 40 institutions, ~ 320 authors (S3)

4 major search groups (different targets and methods):

- Binary inspirals: short inspiral signals (*modeled*)
- Bursts: short *unmodeled* signals (supernovae, merger)
- Stochastic background: cosmological background GWs
- **“Continuous waves”**: spinning NS signals (long-lived)

# Nature of GW from Rotating Neutron Stars

□ **NS frame**: monochromatic wave, slowly varying frequency

☞ Phase  $\Phi(\tau) = \phi_0 + 2\pi \left( f\tau + \frac{1}{2}\dot{f}\tau^2 + \dots \right)$

GW frequency for triaxial NS:  $f = 2\nu$ , r-modes:  $f = 4/3\nu$ , precession:  $f \approx \nu$

☞ Amplitudes (2 polarizations)

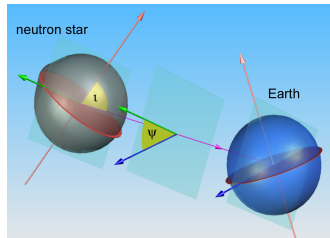
$$A_+ = h_0 \cos \iota$$

$$A_\times = \frac{1}{2} h_0 (1 + \cos^2 \iota)$$

⇒ Wave-components in **NS frame**:

$$h_\times(\tau) = A_\times \cos \Phi(\tau)$$

$$h_+(\tau) = A_+ \sin \Phi(\tau)$$



□ **Detector frame**  $t$ : sky-position  $(\alpha, \delta)$  dependent *modulations*:

● **Phase**: Doppler-effect due to earth's motion  $\tau = \tau(t; \alpha, \delta)$

● **Amplitude**: rotating Antenna-pattern  $F_{+,\times}(t, \psi; \alpha, \delta)$



# Signal Received at the Detector

GW strain at the detector:

$$h(t) = F_+(t) h_+(t) + F_\times(t) h_\times(t)$$

## Signal dependencies

$$h(t) = F_+(t, \psi; \alpha, \delta) A_+ \cos \Phi(t, \phi_0; \alpha, \delta, f, \dot{f}, \dots) \\ + F_\times(t, \psi; \alpha, \delta) A_\times \sin \Phi(t, \phi_0; \alpha, \delta, f, \dot{f}, \dots)$$

Signal parameters:

- “Amplitude parameters”:  $\mathcal{A} = \{h_0, \iota, \psi, \phi_0\}$
- “Doppler parameters”:  $\lambda = \{\alpha, \delta, f, \dot{f}, \dots$  (+ orbital parameters) }

# Matched filtering

Measured strain:  $x(t) = \underbrace{n(t)}_{\text{noise}} + \underbrace{h(t)}_{\text{signal}}$       noise power  $S_n(f)$


Optimal detection statistic:

“matched filtering” = correlate  $x(t)$  with signal-model  $h(t; \mathcal{A}, \lambda)$


⇒ find maximum “match”  $M = M(x; \mathcal{A}, \lambda)$

Signal-to-noise ratio @ perfect match

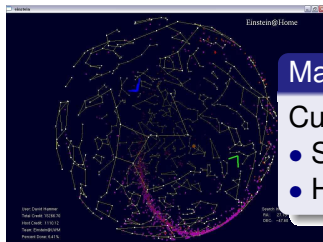
$$\text{SNR}^2 = \int_0^T \frac{h^2(t; \mathcal{A}, \lambda)}{S_n(f)} dt \quad \Rightarrow \quad \text{SNR} \propto \frac{h_0}{\sqrt{S_n}} \sqrt{T \mathcal{N}}$$

$h_0/\sqrt{S_n} \ll 1$   want long  $T$  (and many detectors  $\mathcal{N}$ )

# Search Strategies

- Wide-parameter searches for **unknown** NS:  
Need to **scan** space of Doppler-parameters  $\lambda$  (but not  $\mathcal{A}$ )  
e.g. isolated NS  $(\alpha, \delta, f, \dot{f})$ : number of templates  $N_p \propto T^5$ 
  - ① Fully coherent:  $\mathcal{F}$ -statistic (Einstein@Home  $T \lesssim 30$  hours)  
👉 optimal sensitivity @ *infinite* computing power
  - ② Semi-coherent: Hough, StackSlide, PowerFlux ( $T \sim$  data)  
👉 sub-optimal but fast
  - ③ **Hierarchical search**: combine 1 + 2, will run on E@H   
👉 optimal sensitivity @ *finite* computing power
- Targeted searches for **known** pulsars ( $f = 2\nu$ )  
👉 only *one* template  $\lambda_0 = \{\alpha, \delta, f, \dot{f}, \dots\}$  from radio/X-ray  
Fully coherent, not computationally limited ( $T \sim$  data),  
⇒ most sensitive search!

# Einstein@Home: Search for Unknown NS



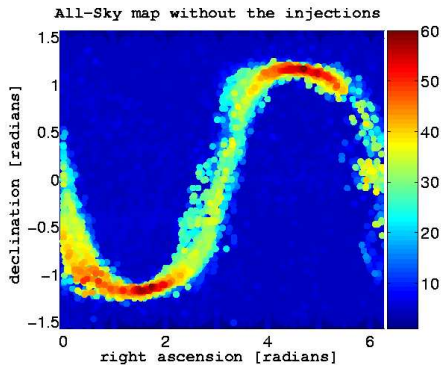
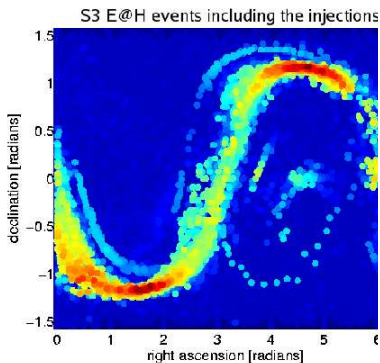
Maximize available computing power

Cut parameter-space  $\lambda$  in small pieces  $\Delta\lambda$

- Send workunits  $\Delta\lambda$  to participating hosts
- Hosts return finished work and request next

- Public distributed computing project, launched Feb. 2005
- Currently  $\sim 120,000$  active participants,  $\sim 50$ Tflops
- runs on GNU/Linux, Mac OSX, Windows,..
- Search for isolated neutron stars  $f \in [50, 1500]$  Hz
- Aiming for **detection**, not upper limits
- Analyzed data from S3, S4, next: S5

# Einstein@Home S3 results



- correctly identified injections ( $h_0 \sim 10^{-23}$ )
- all “outliers” either on  $\mathbf{r}(t) \cdot \mathbf{n} = 0$  circles (👉 stationary lines), or ruled out by follow-up studies (S4)

# Wide-Parameter Searches: (Best) Upper Limits

□ Fully coherent ( $\mathcal{F}$ -statistic) searches [gr-qc/0605028]:

S2 Sco X-1 (unknown  $f, a_p, \bar{T}$ ), using  $T = 6 h$  of S2

☞  $h_0^{95\%} \sim 2 \times 10^{-22}$

S2 All-sky, isolated NS, ( $f \in [160, 728]$  Hz), using  $T = 10 h$  of S2

☞  $h_0^{95\%} \sim 7 \times 10^{-23}$

□ Semi-coherent searches:

S2 Hough-transform: all-sky, isolated NS ( $f \in [200, 400]$  Hz)

☞  $h_0^{95\%} \sim 4.5 \times 10^{-23}$

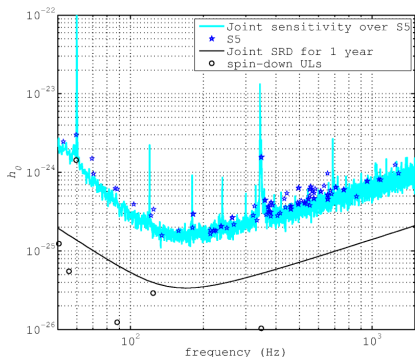
S4 StackSlide: all-sky, isolated NS ( $f \in [50, 225]$  Hz)

☞  $h_0^{95\%} \sim 4.5 \times 10^{-24}$  (preliminary)

Early S5 PowerFlux: all-sky, isolated NS ( $f \in [40, 700]$  Hz)

☞  $h_0^{95\%} \sim 2 \times 10^{-24}$  (preliminary)

## Targeted Pulsar Search: Early S5 (*preliminary*)



- Targeted 73 pulsars ( $f = 2\nu$ ):  
 32 isolated, 41 binary (29 in GCs)
- first 2 months of S5
- all 3 detectors: H1, H2, L1
- Best 95% upper limits:  
 $h_0 \lesssim 2 \times 10^{-25}$  (PSR J1603-7202)  
 $\epsilon \lesssim 4 \times 10^{-7}$  (PSR J2124-3358)

Upper-limits well above *spindown-limit* (except in GCs)

*But:* Crab-pulsar is only a factor **2.1** away from spindown-limit

👉 will (most likely) be able to beat spindown-limit during S5!

## Published results

### Published LSC results of neutron-star searches:

- S1 Setting upper limits on the strength of periodic gravitational waves from PSR J1939 + 2134 using the first science data from the GEO 600 and LIGO detectors, B. Abbott et al. (LSC), Phys. Rev. D 69, 082004 (2004)
- S2 Limits on gravitational wave emission from selected pulsars using LIGO data, B. Abbott et al. (LSC), Phys. Rev. Lett. 94, 181103 (2005)
- S2 First all-sky upper limits from LIGO on the strength of periodic gravitational waves using the Hough transform, B. Abbott et al. (LSC), Phys. Rev. D 72, 102004 (2005)
- S2 Coherent searches for periodic gravitational waves from unknown isolated sources and Scorpius X-1: results from the second LIGO science run, to be submitted, [gr-qc/0605028]
- S3 Online report on Einstein@Home results for S3 search:  
<http://einstein.phys.uwm.edu/PartialS3Results/>



## Summary and outlook

- No GW detection so far, but none expected
  - ➡ setting upper limits on  $h_0$  and  $\epsilon$
- S5 upper-limits are approaching astrophysically relevant regimes (➡ Crab, EOS-limits on  $\epsilon$ )
- LIGO S5 operating at design-sensitivity, will collect one year's worth of data (duration  $\sim 1.5$  years)
- Einstein@Home: Currently analyzing S4, S5.  
Developing a fully hierarchical search ➡ most sensitive possible search for unknown NS
- NS detection with LIGO-I not very likely, but not impossible (“Expect the unexpected!”)
- The future is bright: S6, VIRGO, LIGO-II, GEO-HF, ...

# You can help us find Gravitational Waves!

