



Measuring Kerrness in Binary Black Hole Simulation Ringdowns

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Acknowledgments

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Acknowledgments

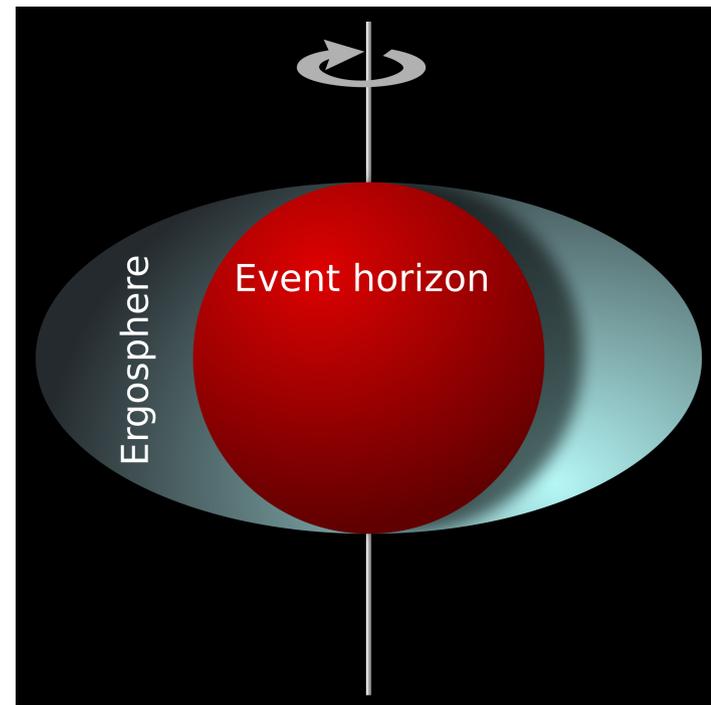
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Kerr

- Kerr is the spacetime of a single black hole which is
 - Axially symmetric
 - Uncharged
 - Asymptotically flat
 - Spinning

<https://upload.wikimedia.org/wikipedia/commons/thumb/0/0c/Ergosphere.svg/2000px-Ergosphere.svg.png>



“Kerrness”

- **Local** scalar representing closeness to Kerr

Local quantity: At a given (4D) point in spacetime, we may compute the similarity to Kerr

- These quantities are **invariant of choice of coordinates**

Important for implementing in a numerical code

Why measure Kerrness?

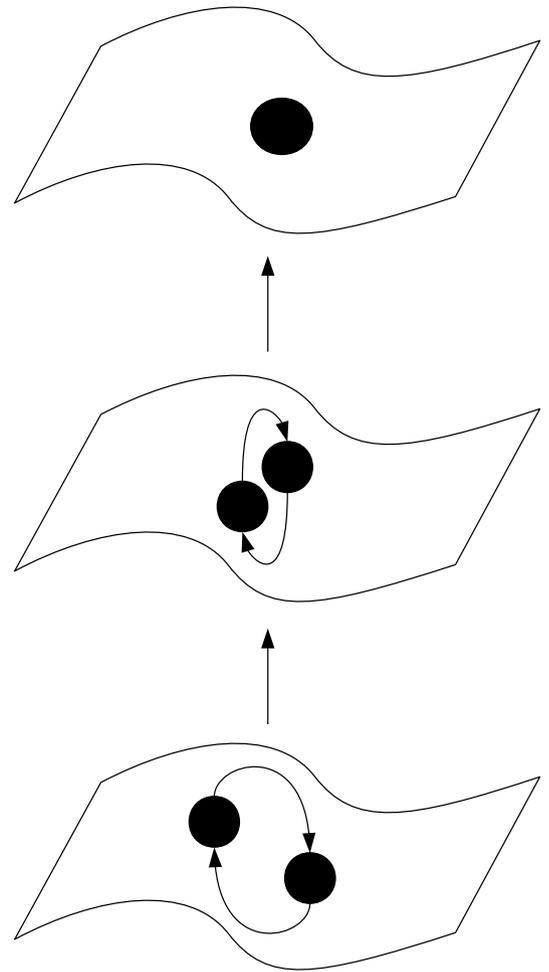
- Analyses of binary black hole mergers assume a Kerr remnant (or a perturbation) [4, 5]
- LIGO data analyses may assume that the resulting spacetime is Kerr (or a perturbation) [4, 5]
- Measures of Kerrness applied to binary black hole merger simulations may help to **quantify** when it is valid to make these assumptions

SpEC

- The simulations were run using the **Spectral Einstein Code (SpEC)**

Spectral methods: compute coefficients of basis functions

- Codebase in **C++** (with Perl for parsing input files)
- Simulations consist of **slices** (three dimensional spacelike hypersurfaces) evolved in time
- **3 + 1** formalism



Speciality Index [1]

$$S = \frac{27J^2}{I^3}, \quad I = \tilde{C}_{abcd}\tilde{C}^{abcd}, \quad J = \tilde{C}_{abcd}\tilde{C}^{cd}_{mn}\tilde{C}^{mnab} \quad (1)$$

$$\tilde{C}_{abcd} = C_{abcd} + (i/2)\epsilon_{abmn}C^{mn}_{cd} \quad (2)$$

- Computed from contractions of the self-dual **Weyl Tensor**
- **Complex** quantity
- $\text{Re}[S] \rightarrow 1, \text{Im}[S] \rightarrow 0$
- **Necessary** but **not sufficient** condition

This quantity is 1 for any algebraically special spacetime
Kerr \subset Algebraically special [1]

García-Parrado 2015 [2]

$$\mathcal{L} \equiv F_1 + F_2 + F_3 + F_4 + F_5 + F_6 \quad (1)$$

- $F_1 - F_6$ are expressions involving contractions and covariant derivatives
- **Directly** measures similarity to Kerr
- **Real, non-negative** quantity which vanishes for Kerr spacetime

Each term **independently** vanishes

It is possible to consider each term independently of the others during debugging and analysis

García-Parrado 2015

$$\mathcal{L} \equiv \frac{(\mathbf{r}(A) + \mathbf{r}(B))^2 + (\mathbf{j}(A)_i + \mathbf{j}(B)_i)(\mathbf{j}(A)^i + \mathbf{j}(B)^i) + (\mathbf{t}(A)_{ij} + \mathbf{t}(B)_{ij})(\mathbf{t}(A)^{ij} + \mathbf{t}(B)^{ij})}{\sigma^{14}} + \frac{\mathbf{a}_{ij}\mathbf{a}^{ij} + \mathbf{b}_{ij}\mathbf{b}^{ij}}{\sigma^4} + \frac{((1 - 3\lambda^2)\beta + \lambda(3 - \lambda^2)\alpha)^2}{\sigma^2} + \frac{(\mathfrak{B}_{ij}\mathfrak{B}^{ij})^3}{\sigma^4} + \frac{(\mathfrak{C}_{ij}\mathfrak{C}^{ij})^3}{\sigma^7} + \frac{\Omega}{\sigma^2}, \quad (1)$$

$$\begin{aligned} \mathbf{r}(A) &\equiv A_{\nu\mu}n^\nu n^\mu = \mathcal{E}(Q)_{\mu\nu}D^\mu AD^\nu A, \\ \mathbf{j}(A)_\mu &\equiv A_{\nu\rho}n^\nu h^\rho_\mu = -\varepsilon_{\mu\kappa\pi}\mathcal{B}(Q)_\lambda{}^\pi D^\lambda AD^\kappa A - \mathcal{E}(Q)_{\mu\lambda}\mathcal{A}(A)D^\lambda A, \\ \mathbf{t}(A)_{\mu\nu} &\equiv A_{\kappa\pi}h^\kappa_\mu h^\pi_\nu = \\ &D^\kappa A(-2D_{(\nu}A\mathcal{E}(Q)_{\mu)\kappa} + h_{\mu\nu}\mathcal{E}(Q)_{\kappa\pi}D^\pi A + 2\mathcal{A}(A)\mathcal{B}(Q)_{(\mu}{}^\pi\varepsilon_{\nu)\kappa\pi}) + \\ &+\mathcal{E}(Q)_{\mu\nu}(D_\kappa AD^\kappa A + (\mathcal{A}(A))^2). \\ \mathbf{a}_{\mu\nu} &\equiv -B_\mu{}^\lambda B_{\nu\lambda} + E_\mu{}^\lambda E_{\nu\lambda} - \frac{1}{3}Ah_{\mu\nu} - E_{\mu\nu}\alpha - B_{\mu\nu}\beta = 0, \\ \mathbf{b}_{\mu\nu} &\equiv B_\mu{}^\lambda E_{\nu\lambda} + B_\nu{}^\lambda E_{\mu\lambda} - \frac{1}{3}Bh_{\mu\nu} - B_{\mu\nu}\alpha + E_{\mu\nu}\beta = 0. \\ A &\equiv \frac{1}{8}C^{\mu\nu\lambda\rho}C_{\mu\nu\lambda\rho}, \\ B &\equiv \frac{1}{8}C^{\mu\nu\lambda\rho}C_{\mu\nu\lambda\rho}^*, \\ D &\equiv \frac{1}{16}C_{\mu\nu\lambda\rho}C_{\sigma\pi\mu\nu}C_{\lambda\rho\sigma\pi}, \end{aligned}$$

$$E \equiv \frac{1}{16}C_{\mu\nu\lambda\rho}C_{\sigma\pi\mu\nu}C_{*\lambda\rho\sigma\pi}.$$

$$T \equiv \frac{\beta}{\alpha} = \frac{BD - AE}{AD + BE},$$

$$\lambda \equiv \frac{K(KT + 1)}{K^2 - 3KT - 2},$$

$$K = \frac{2R\|\Theta\| - 2R^\perp{}^\mu\Theta_\mu^\perp}{R\|\Theta\|^2 - \Theta\|\Theta\|^2 - R_\mu^\perp R^{\perp\mu} + \Theta_\mu^\perp\Theta^{\perp\mu}},$$

$$\sigma = \frac{1}{4}(-2 + K^2 - 3KT)^2 \times$$

$$\left(\frac{R\|\Theta\|^2 - \Theta\|\Theta\|^2 - R_\mu^\perp R^{\perp\mu} + \Theta_\mu^\perp\Theta^{\perp\mu}}{K^2(1 + KT)^2 - (K^2 - 3KT - 2)^2} - \frac{8\alpha}{(K^2 - 3KT - 2)^2 - 3K^2(1 + KT)} \right)$$

$$\alpha = -\frac{AD + BE}{A^2 + B^2},$$

$$\beta = \frac{AE - BD}{A^2 + B^2}.$$

García-Parrado 2015 [2]

$$\mathcal{L} \equiv F_1 + F_2 + F_3 + F_4 + F_5 + F_6 \quad (1)$$

- Can be computed on an individual slice

The computations involving the pulled-back tensors **do not involve time derivatives**

Can be computed using **3D quantities** implemented in SpEC

García-Parrado 2015 (continued)

$$\mathcal{L} \equiv F_1 + F_2 + F_3 + F_4 + F_5 + F_6 \quad (1)$$

- The first three terms have been implemented and evaluated on single and binary black hole simulations

Fourth and fifth terms: Equivalent terms from the 2016 paper [3] have been implemented

Simulations

- New code was written for the García-Parrado 2015 quantity
- Single black hole simulations
 - Kerr black hole with mass 1, spin vector (0., 0., 0.4)
 - Sanity check that the quantities behave
 - Allows checking the **convergence** of quantities with respect to resolution
- Binary black hole ringdown
 - The quantities were computed on the ringdown phase of a simulation of the GW150914 event at a single resolution (obtained from California State University, Fullerton)



Evaluation of Kerrness Quantities

Analysis of Convergence

- “Error” measured by L2 norm of deviation from theoretical values ($\text{Re}[S] = 1, \text{Im}[S] = 0$)

$$\sqrt{\sum_i^N (x_i - x)^2 / N}$$

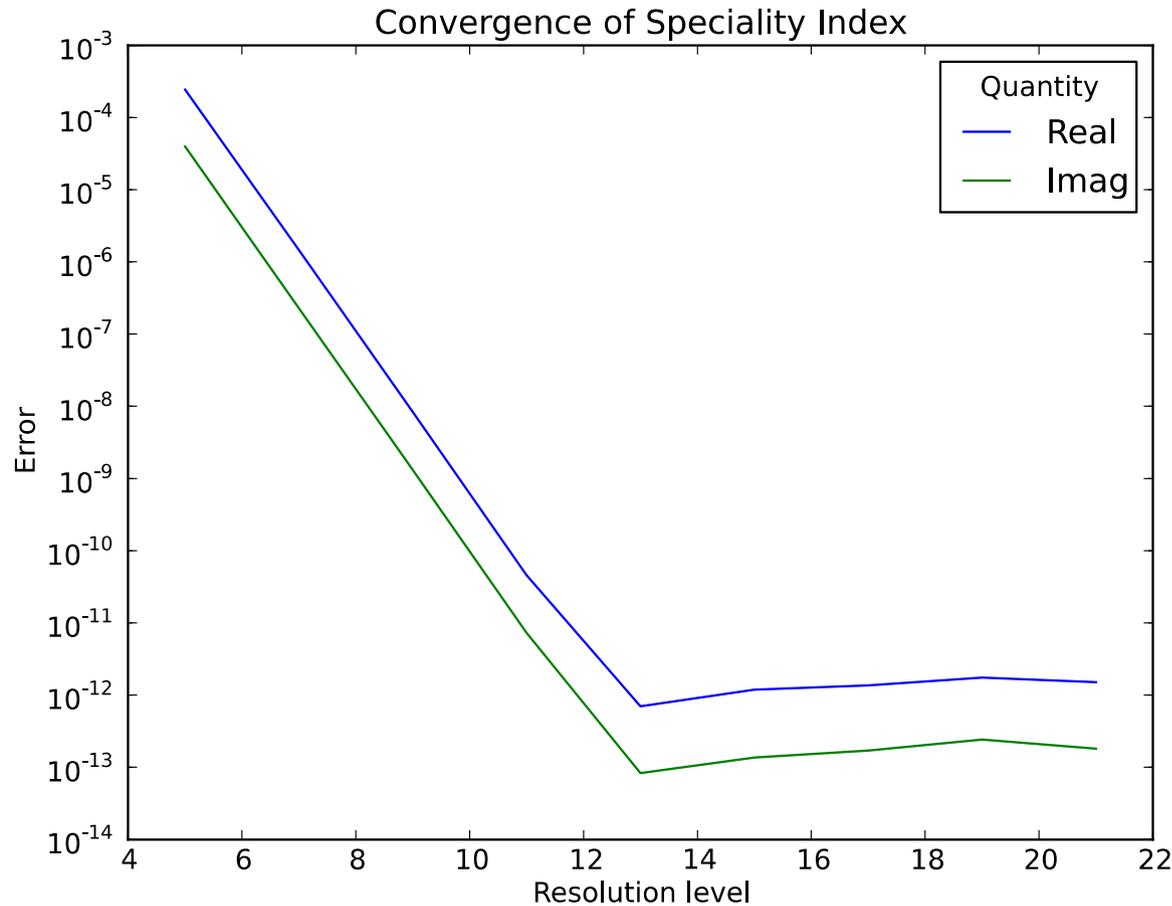
- “Resolution” refers to simulation angular resolution

Spherical harmonics: $Y_{l m}$

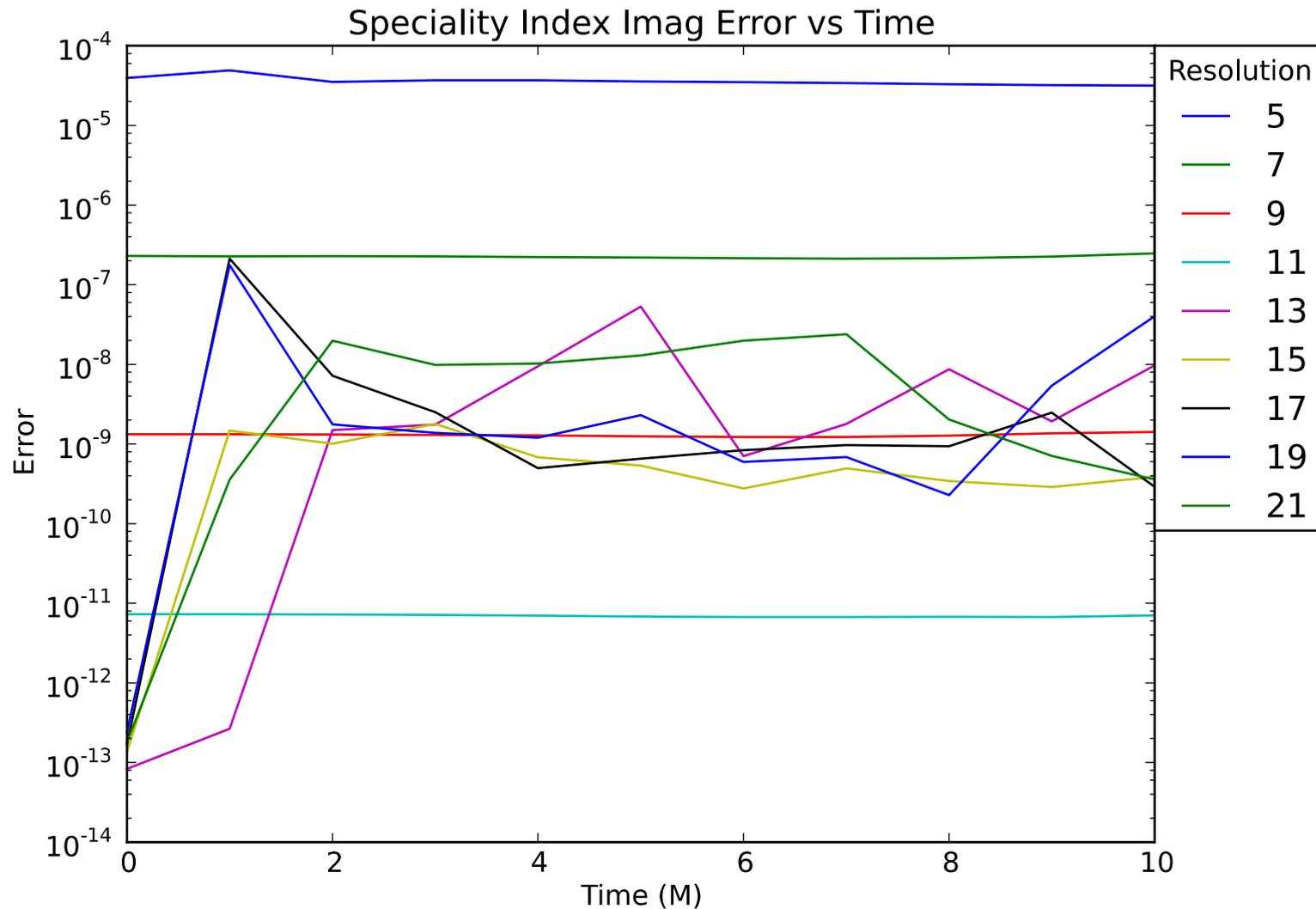
Analysis of Convergence

- The error was computed for points on a spherical shell located 12M from the origin.
- Other spheres (other than the innermost) exhibit similar convergence patterns.

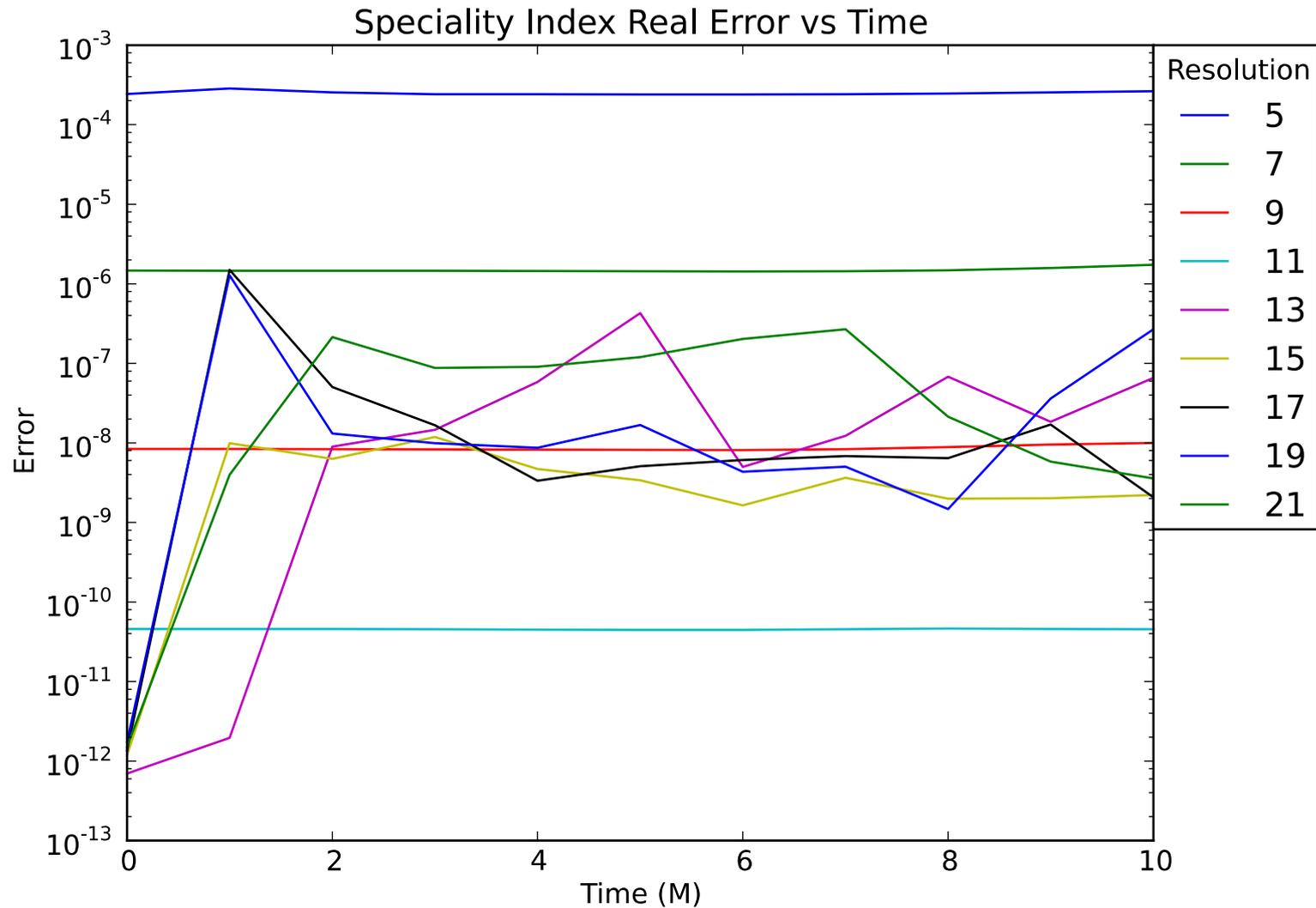
Speciality Index: Single Black Hole



Speciality Index: Single Black Hole



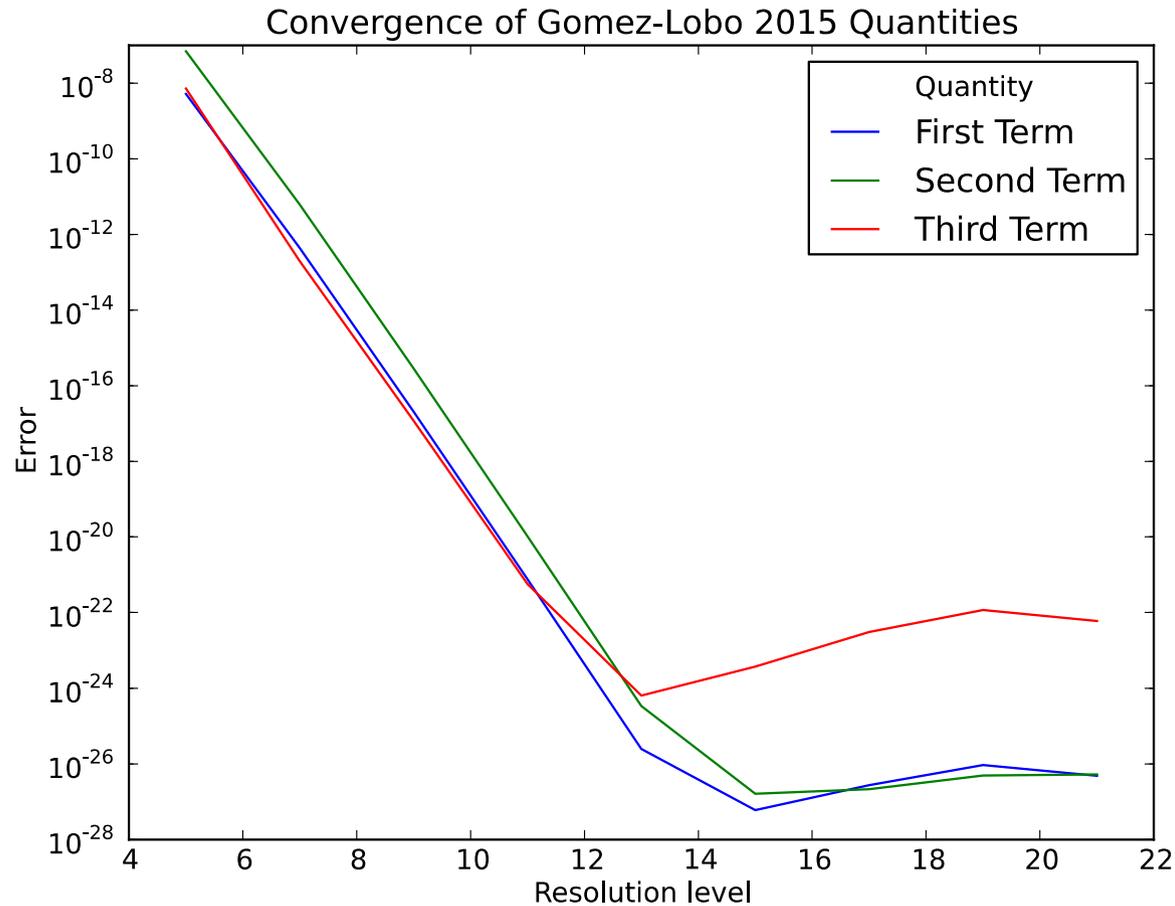
Speciality Index: Single Black Hole



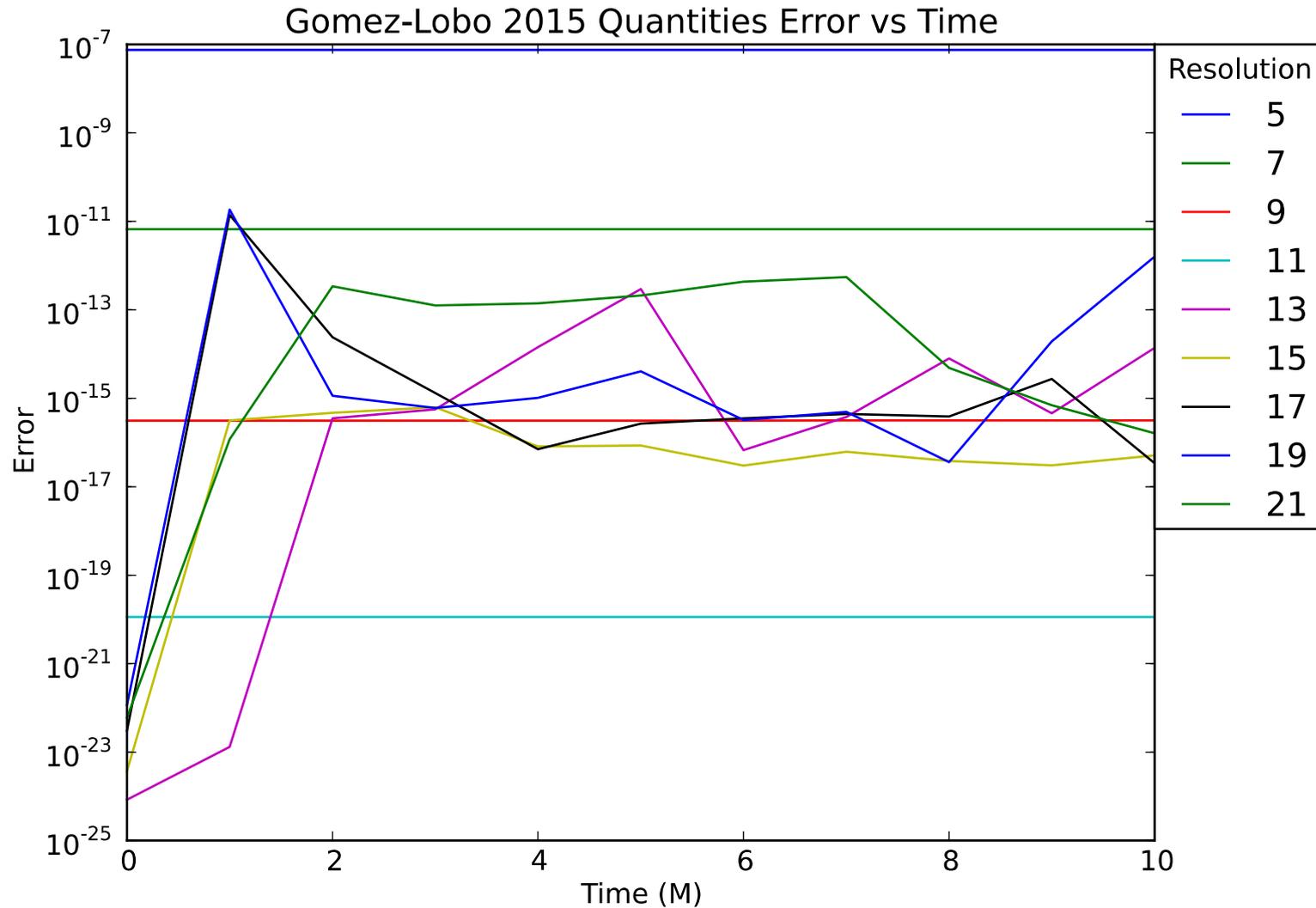
Speciality Index: Binary Black Hole

- Qualitatively, $\text{Re}[S] \rightarrow 1$ and $\text{Im}[S] \rightarrow 0$ as the ringdown progresses
- Computations of the quantities at time resolutions to generate plots are ongoing

García-Parrado 2015: Single Black Hole



García-Parrado 2015: Single Black Hole



García-Parrado 2015: Binary Black Hole

- Qualitatively, $F_{1-3}[S] \rightarrow 0$ as the ringdown progresses
- Computations of the quantities at time resolutions to generate plots are ongoing

Further Work

- Short term:

Finish implementation of remaining terms and compare with corresponding terms from the 2016 paper [3]

- Long term:

Apply Kerrness measures to quantify similarity to Kerr during ringdowns

References

- [1] J. Baker and M. Campanelli. Making use of geometrical invariants in black hole collisions. *Physical Review D*, 62(12):127501, December 2000.
- [2] A. García-Parrado Gómez-Lobo. Local non-negative initial data scalar characterization of the Kerr solution. *Physical Review D*, 92(12):124053, December 2015.
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- [4] B. P. Abbott et al. Properties of the binary black hole merger GW150914. 2016.
- [5] B. P. Abbott et al. Tests of general relativity with GW150914. 2016.