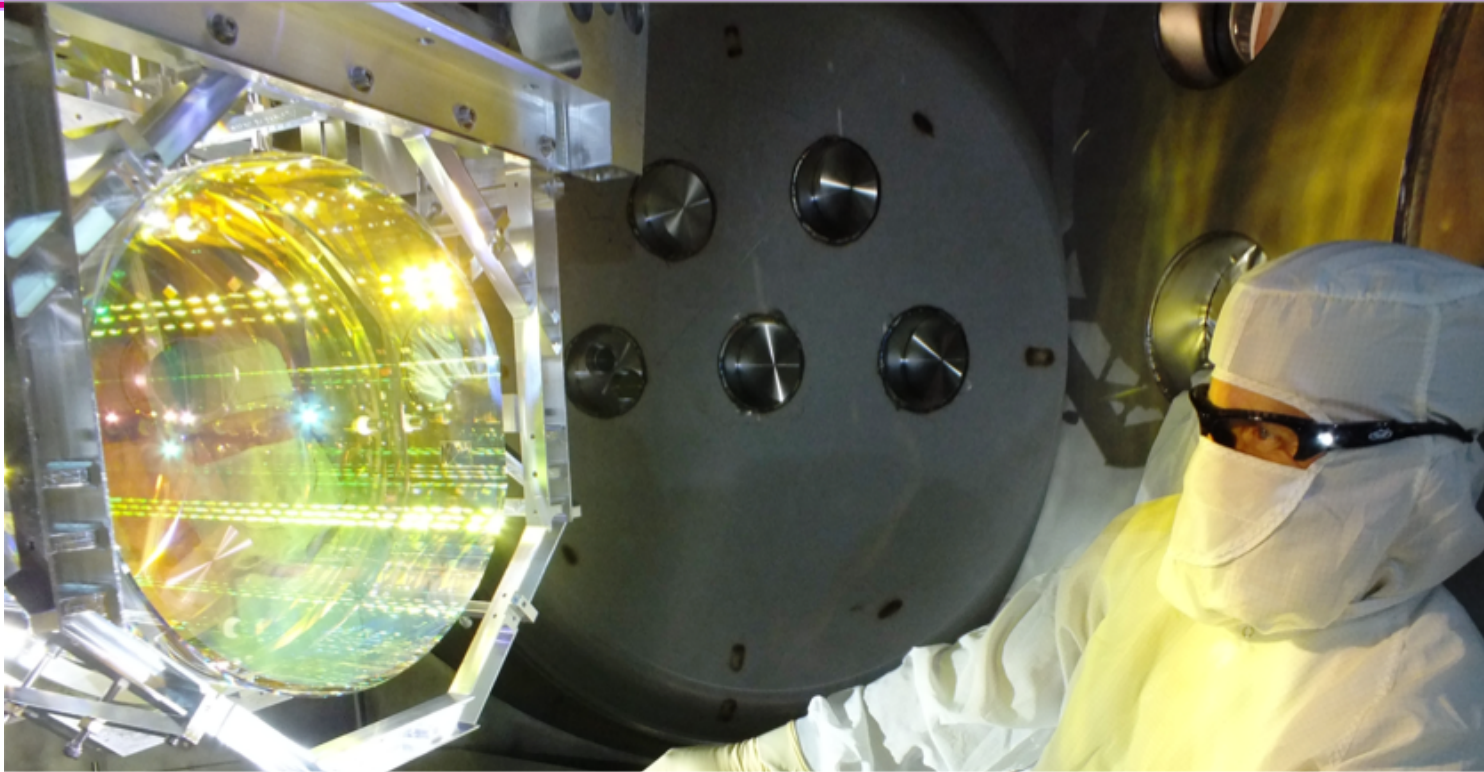




Optical Challenges in LIGO: Past and Future



Stan Whitcomb
CREOL IA Symposium
15 March 2019

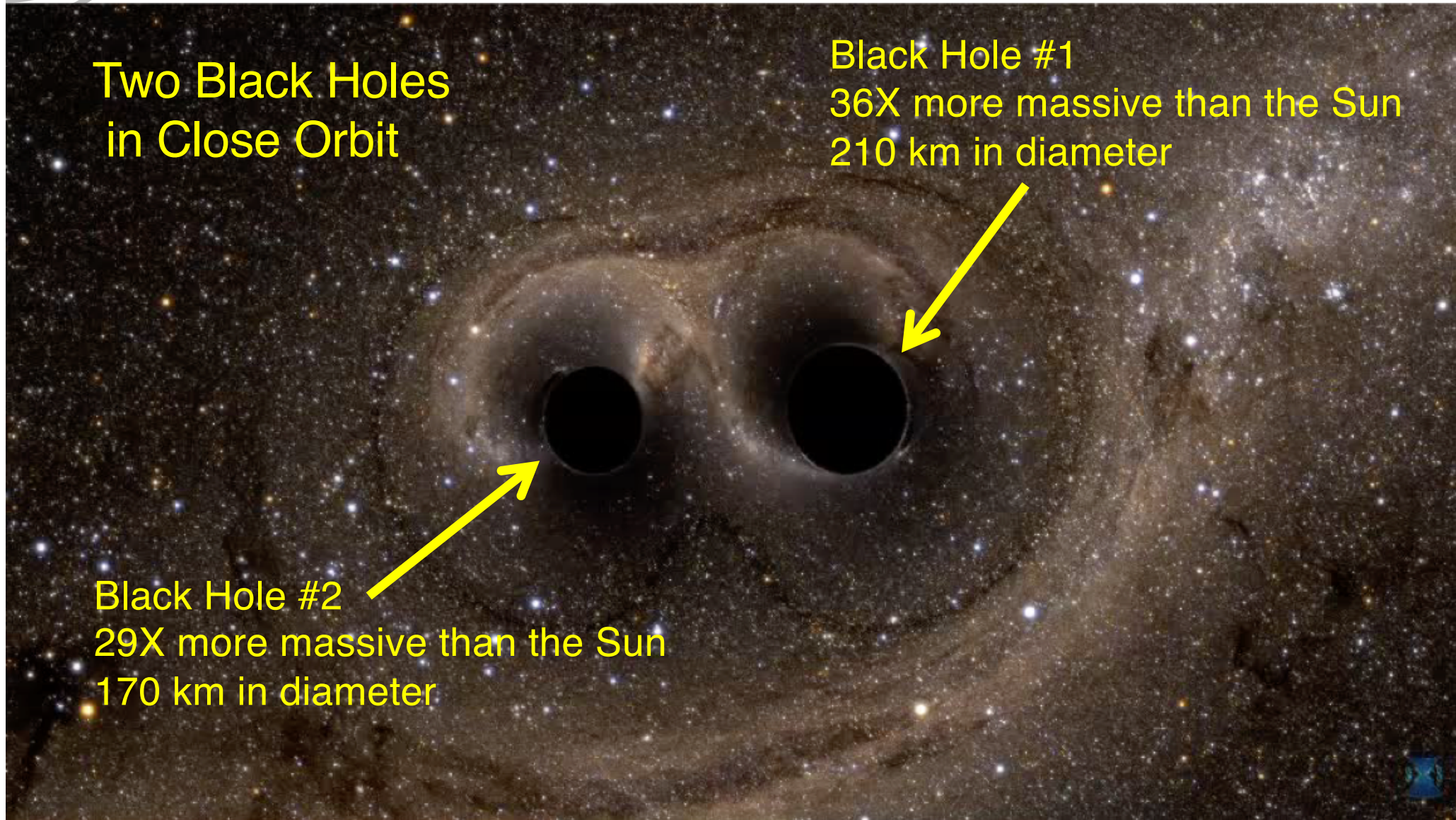


This story begins 1.3 Billion years ago,
in a distant galaxy...

Two Black Holes
in Close Orbit

Black Hole #1
36X more massive than the Sun
210 km in diameter

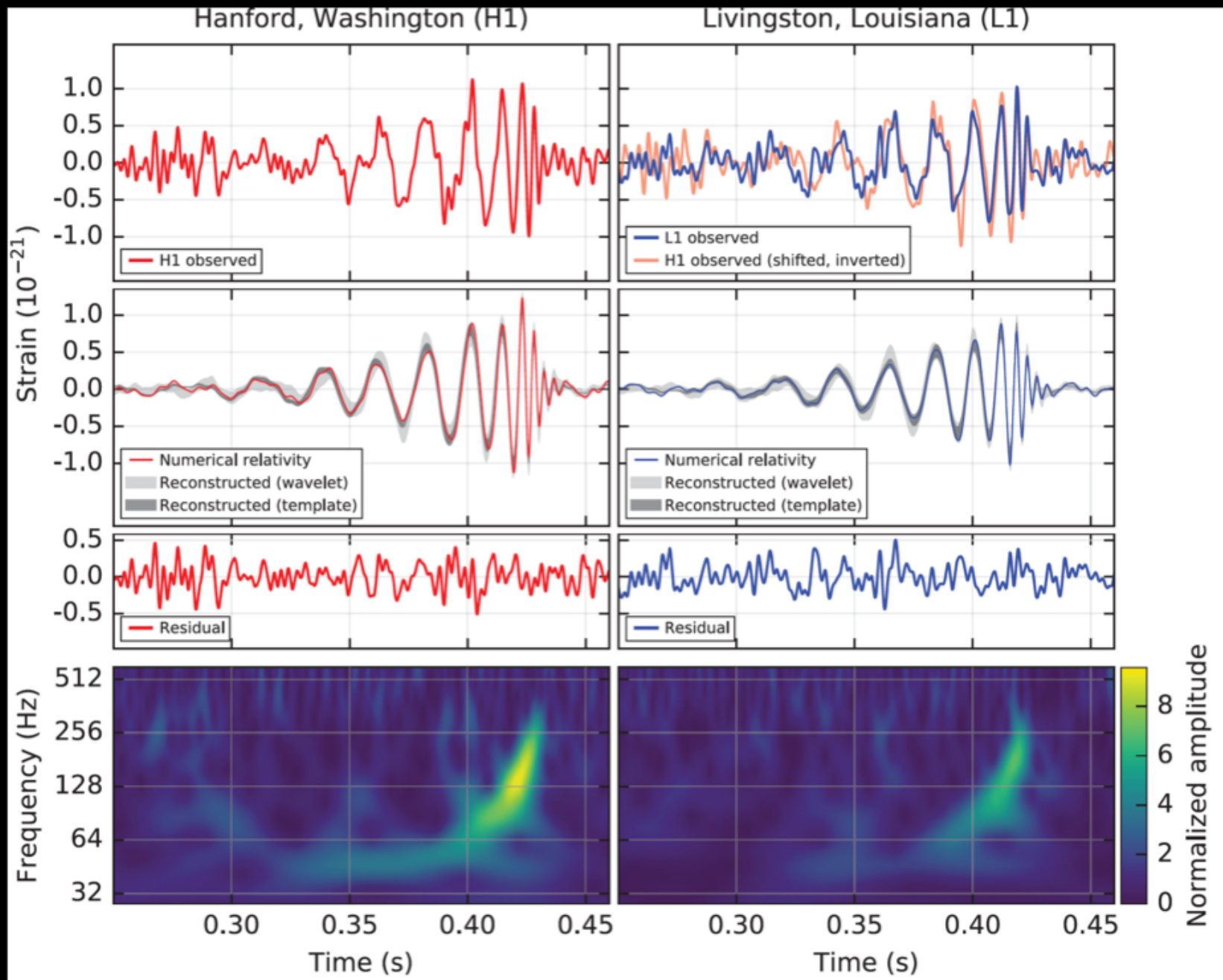
Black Hole #2
29X more massive than the Sun
170 km in diameter





Then on 14 September 2015, at the LIGO sites...



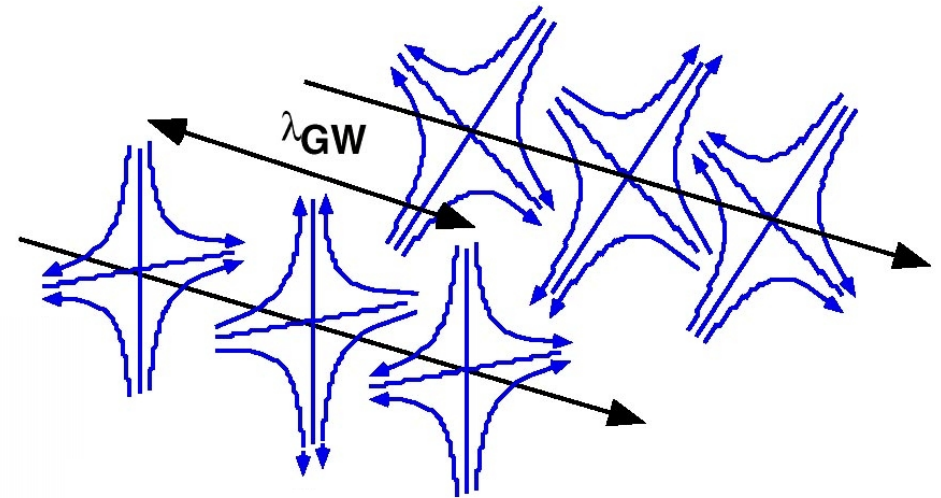


B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration), *Observation of Gravitational Waves from a Binary Black Hole Merger*, Phys. Rev. Lett. 116, 061102 (2016)



Gravitational Wave Basics

- Einstein (in 1916) recognized gravitational waves in his theory of General Relativity
 - » Necessary consequence of Special Relativity with its finite speed for information transfer
 - » Most distinctive departure from Newtonian theory
- Time-dependent distortions of space-time created by the acceleration of masses
 - » Propagate away from the sources at the speed of light
 - » Pure transverse waves
 - » Two orthogonal polarizations

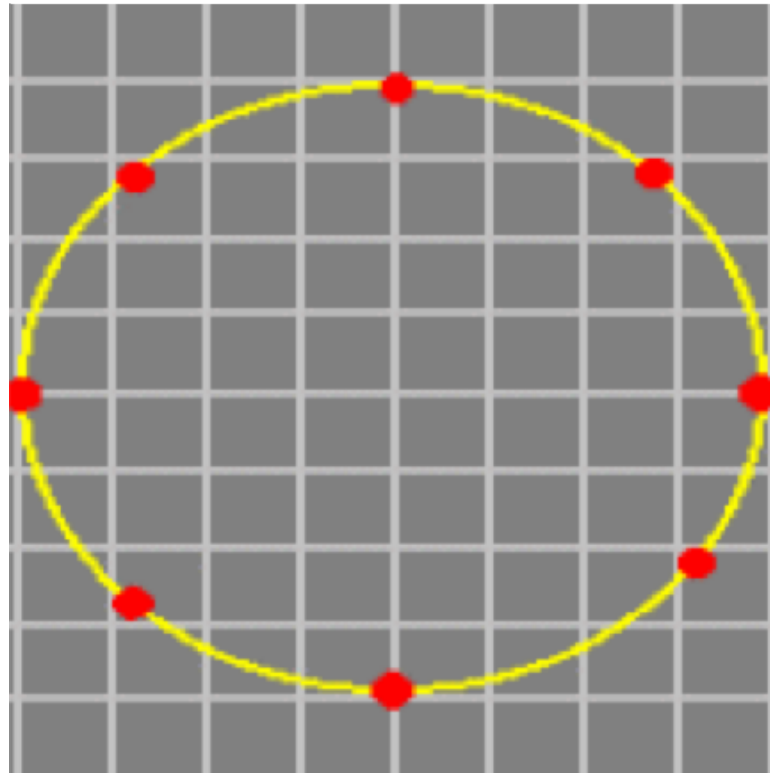


$$h = 2(\Delta L / L)$$

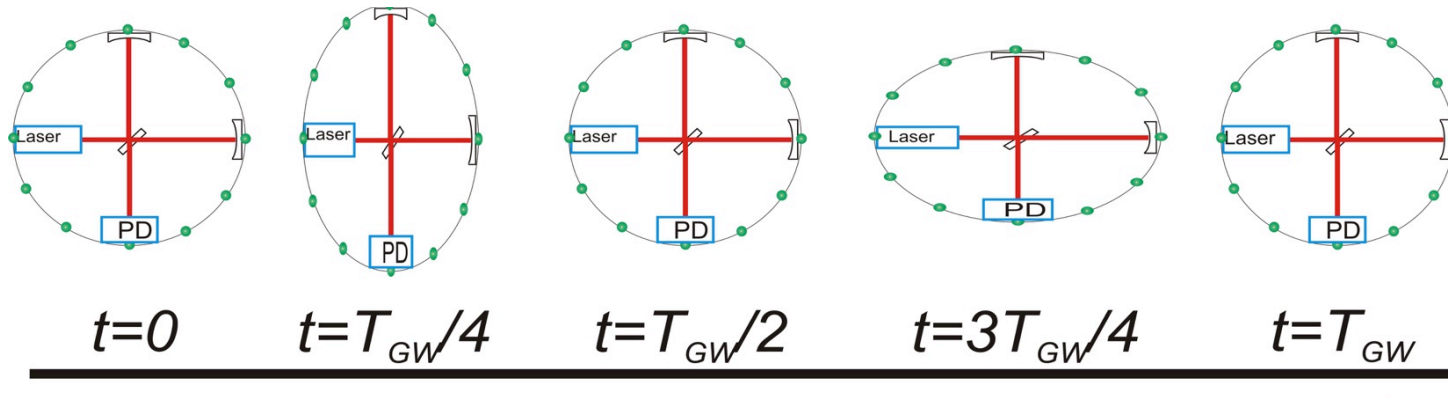
- Requires huge masses and relativistic accelerations
 - $h \sim 10^{-21}$ for plausible astrophysical sources

Effect of Gravitational Wave

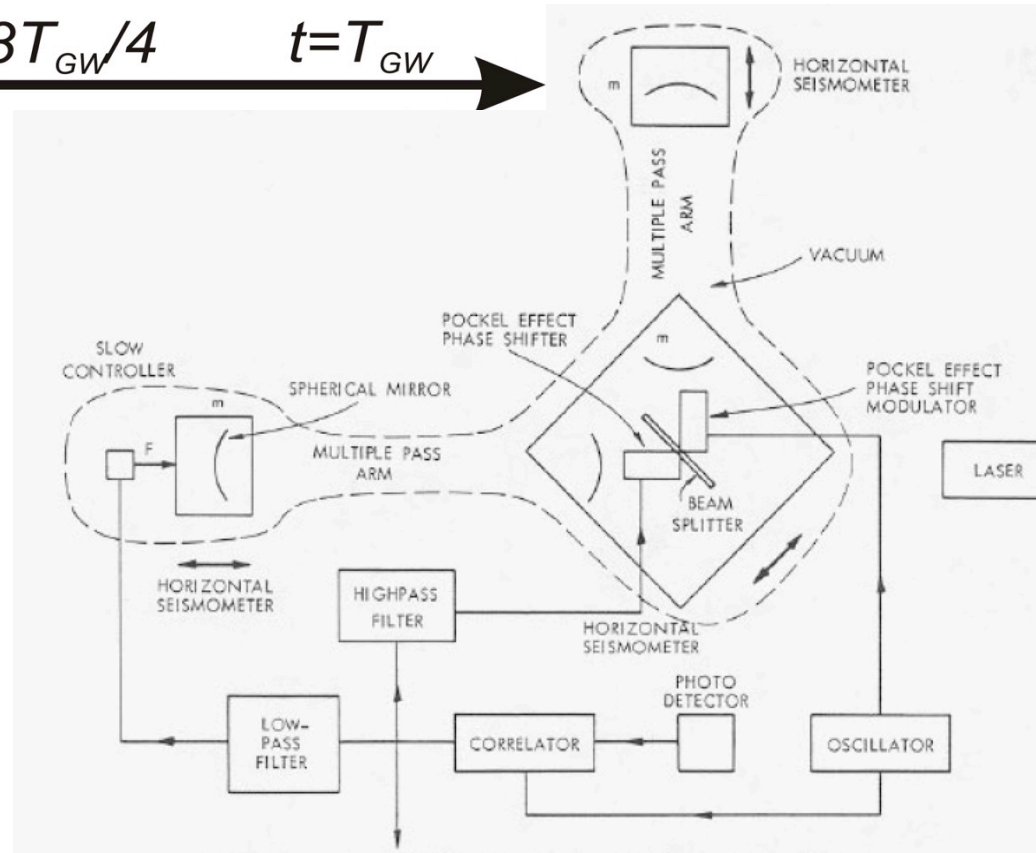
- Gravitational wave travelling INTO the plane of this slide
- Changes the separations of “free masses”



Detecting GWs with Interferometry



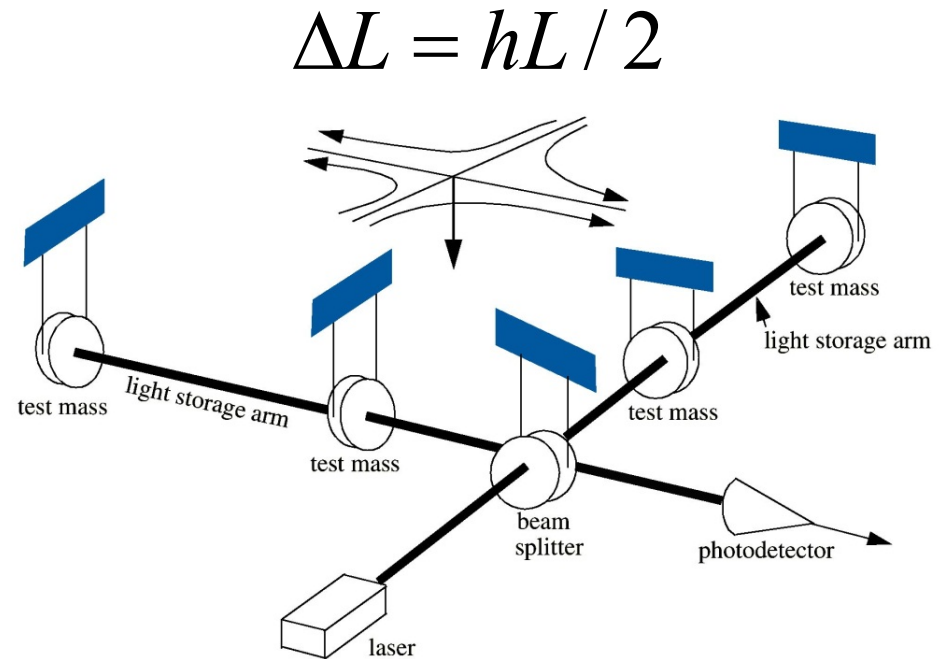
Earliest concepts show interferometer mirrors mounted on a “free test mass”



Weiss, 1972

Suspended mirrors act as “freely-falling” test masses in horizontal plane for frequencies $f \gg f_{\text{pend}}$

For a LIGO detector,
 $L \sim 4 \text{ km}$, $h \sim 10^{-21}$
 $\Delta L \sim 10^{-18} \text{ m}$





Must act like a free mass in GR

Ultra-high purity fused silica
~35 cm diameter, 40 kg

Test Mass

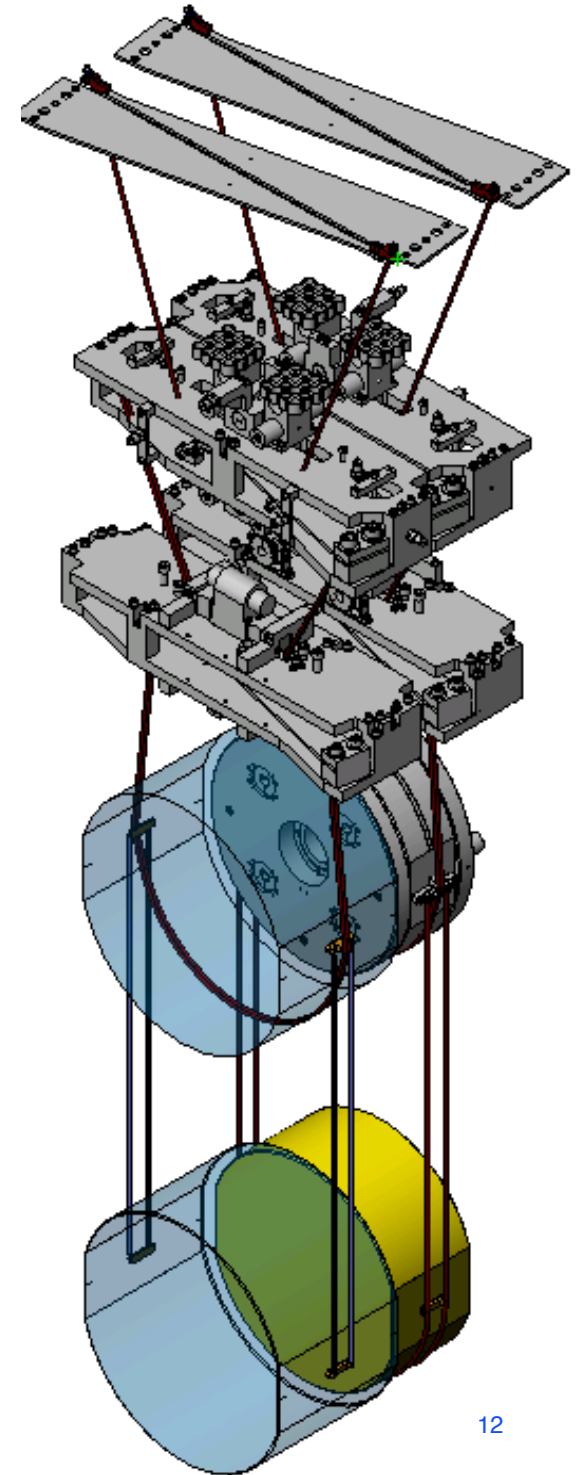
Test Mass Suspension

Four-stage pendulum suspension

Monolithic fused silica suspension fibers

Low thermal noise

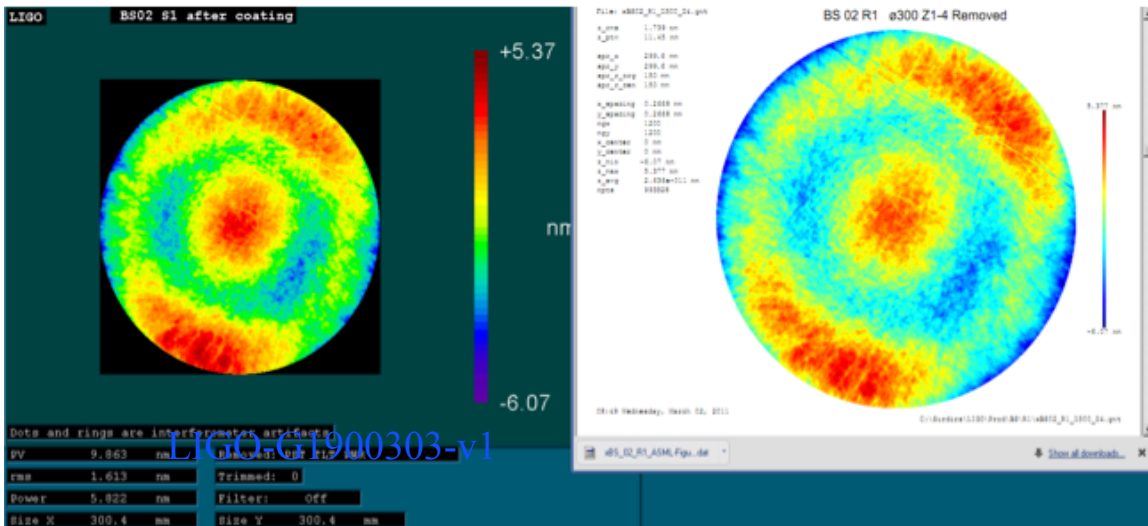
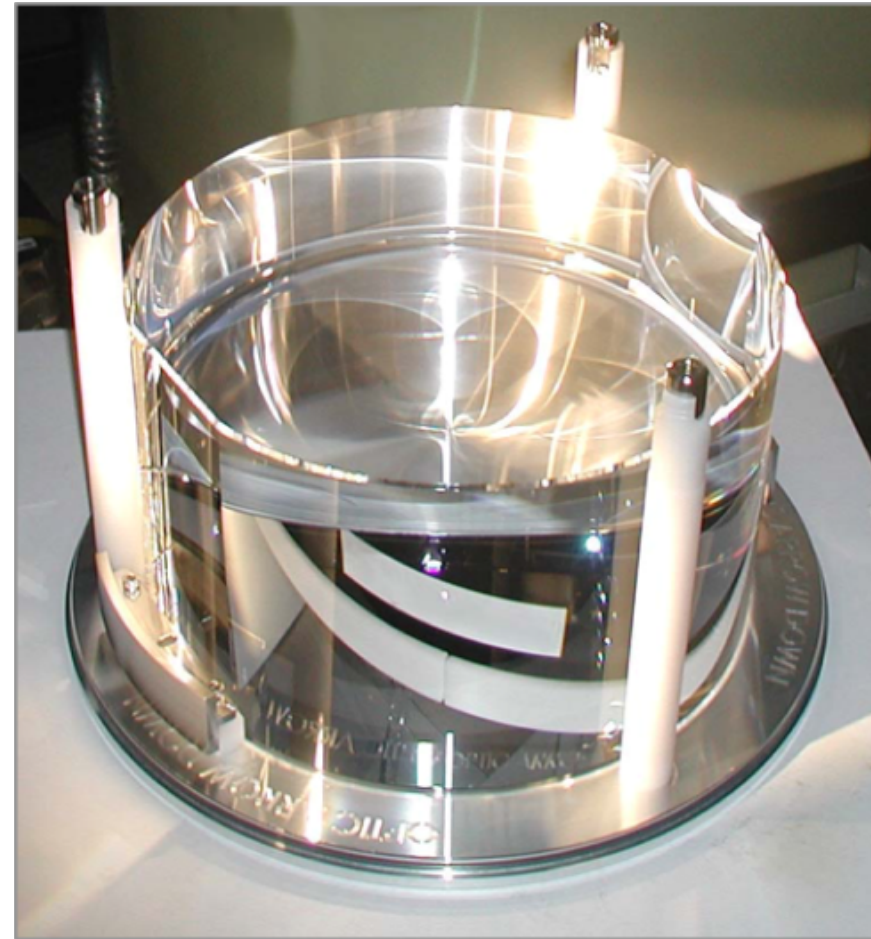
Low noise actuation



Core Optics Specifications

Challenging optical requirements:

- ROC match to <1%
- $\lambda/1000$ surface figure
- < 0.5 ppm absorption
- ~10 ppm scatter
- 0.1 % coating uniformity



Precision Interferometry = Controlling Measurement Noises

Displacement Noise

- Seismic noise
- Radiation Pressure
- Thermal noise
 - Suspensions
 - Optics

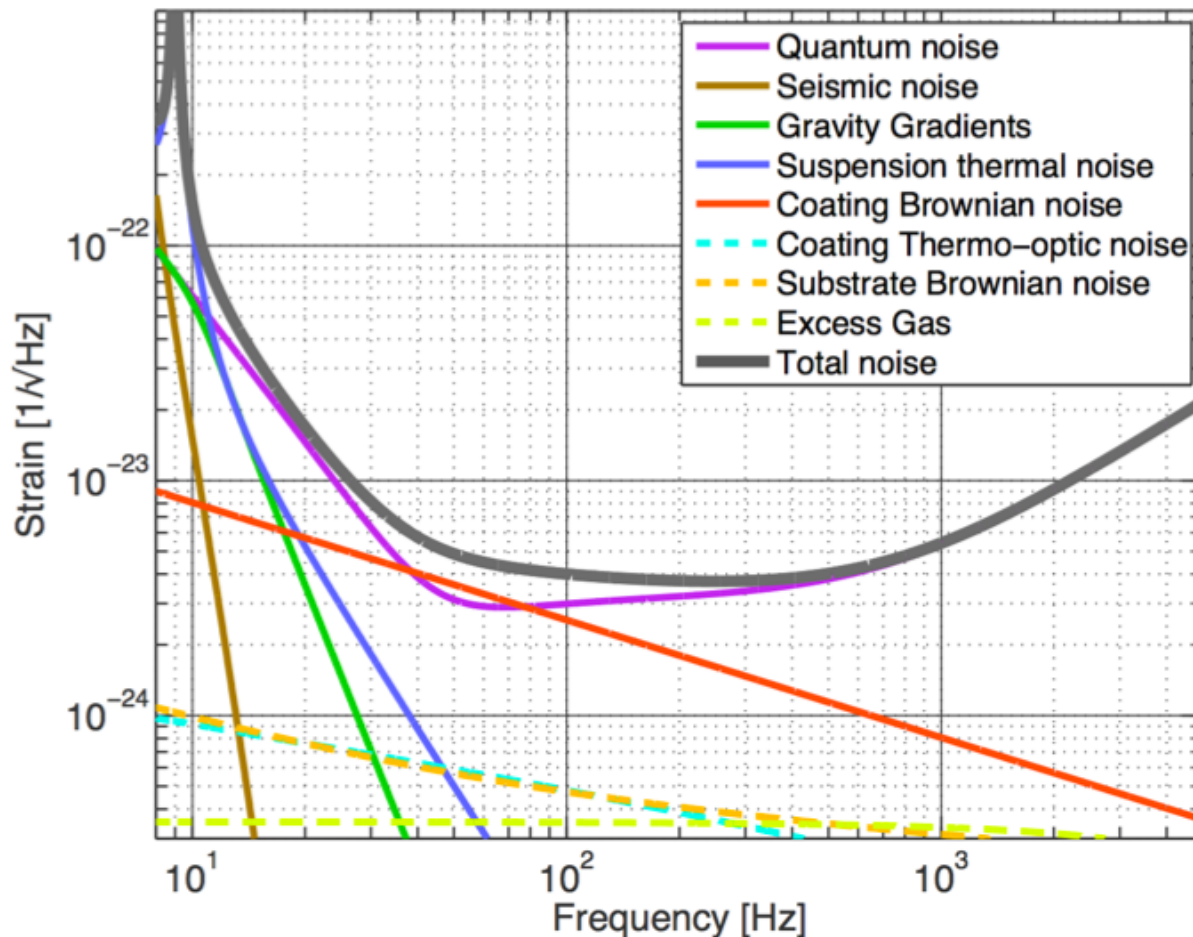
Sensing Noise

- Shot Noise
- Residual Gas

Technical Noises:

Hundreds of them...

Broadband tuning, full input power (125 W)

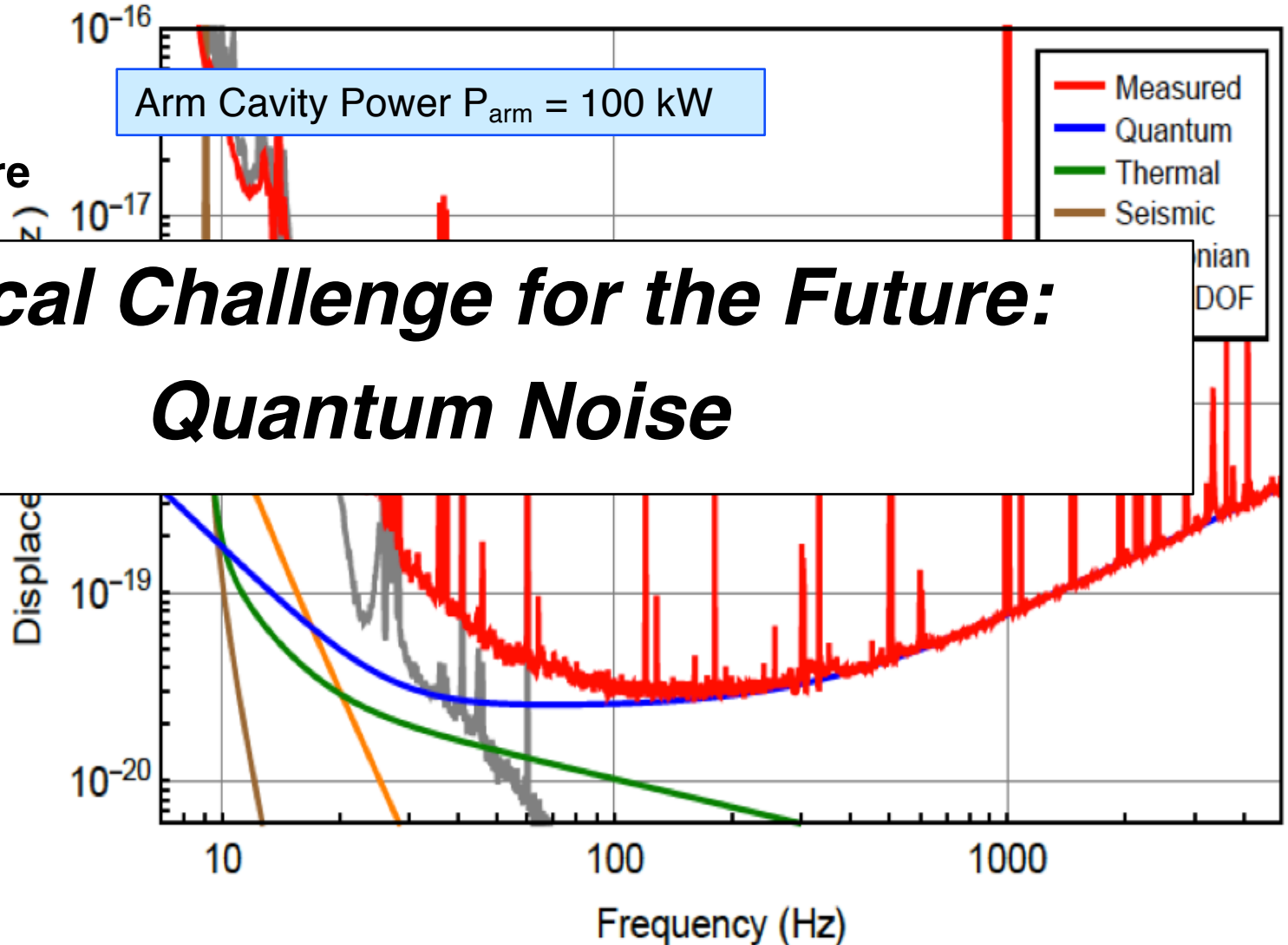




Advanced LIGO Detector Sensitivity During O1

Displacement Noise

- Seismic noise
- Radiation Pressure
- Thermal noise



**Optical Challenge for the Future:
Quantum Noise**

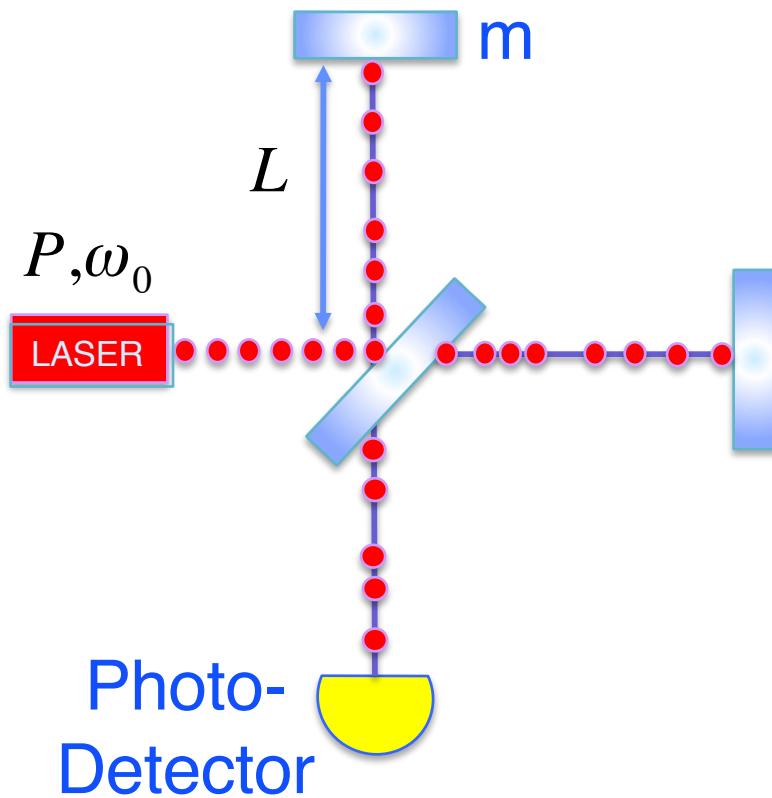
Sensitivity

- Shot Noise
- Residual Gas

Technical Noises:

Hundreds of them...

Abbott, et al. ,LIGO Scientific Collaboration and Virgo Collaboration, "GW150914: The Advanced LIGO Detectors in the Era of First Discoveries", Phys. Rev. Lett. 116, 131103 (2016).



$$h_{\text{quantum}} = \sqrt{h_{\text{rad}}^2 + h_{\text{shot}}^2}$$

SHOT NOISE:

Photon counting noise

$$h_{\text{shot}} \propto \frac{1}{L} \sqrt{\frac{1}{P}}$$

RADIATION PRESSURE NOISE:

Noise caused by photon pressure on mirrors

$$h_{\text{rad}} \propto \frac{1}{f^2 L} \frac{\sqrt{P}}{m}$$

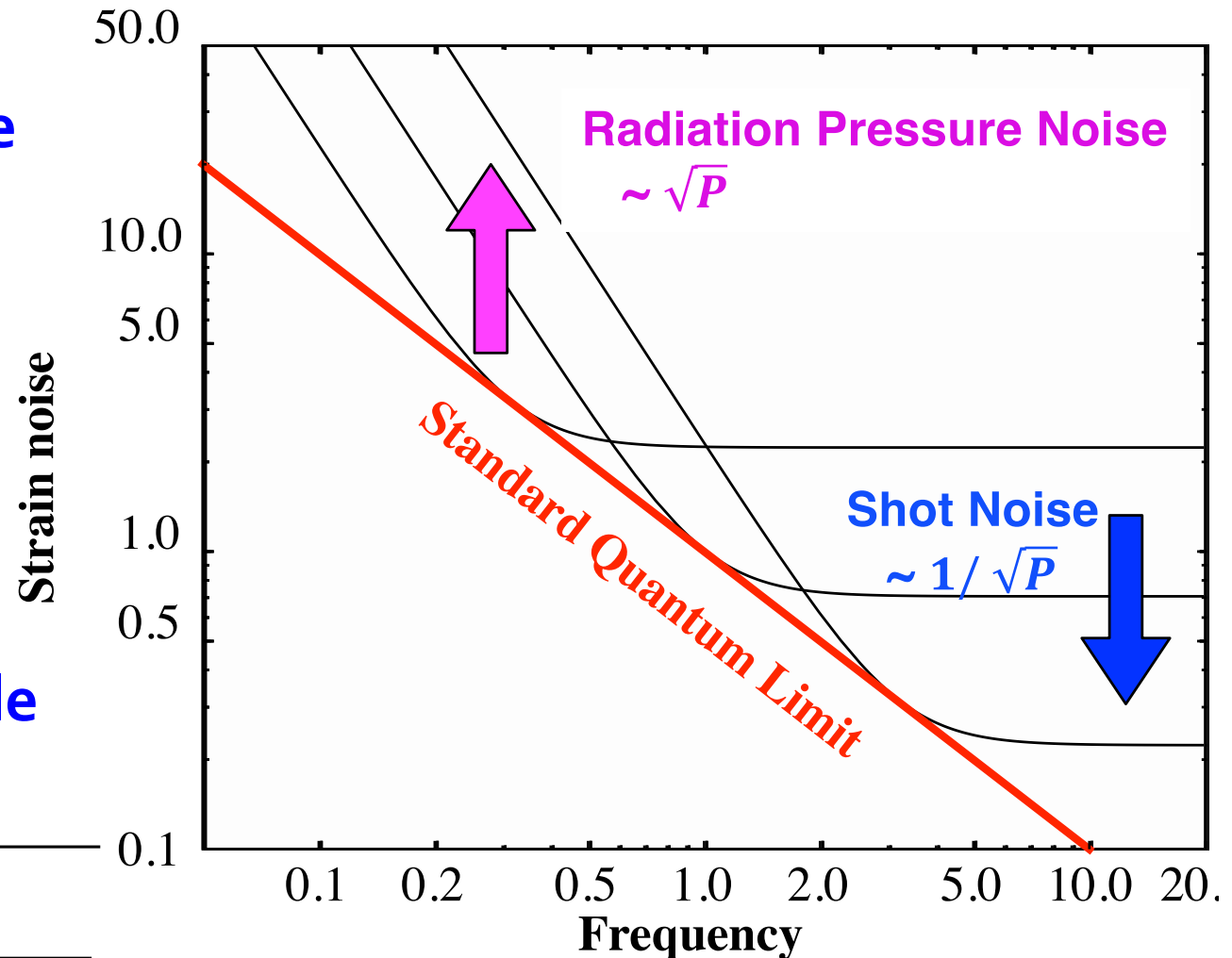


Measurement frequency

Standard Quantum Limit

- Trade-off in Power Between Shot Noise and Radiation-Pressure Noise
- Standard Quantum Limit (SQL): Uncertainty of test mass position due to Heisenberg Uncertainty Principle

$$\sqrt{S_h^{\text{SQL}}} = \sqrt{\frac{8\hbar}{M(2\pi f)^2 L^2}}$$



Reduce Quantum Noise?

- Quantum noise

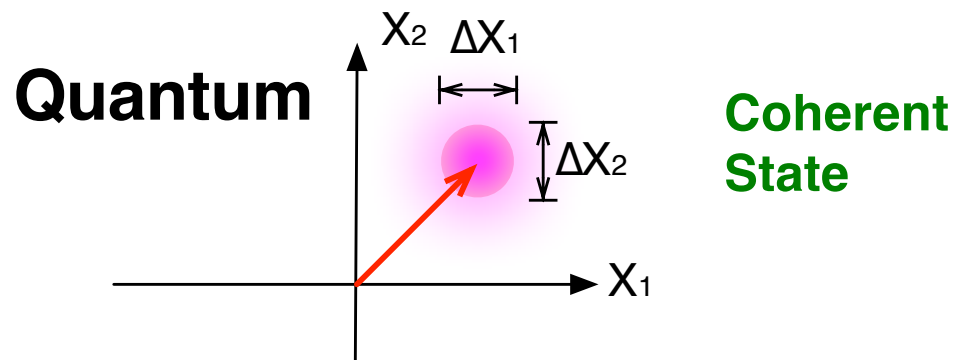
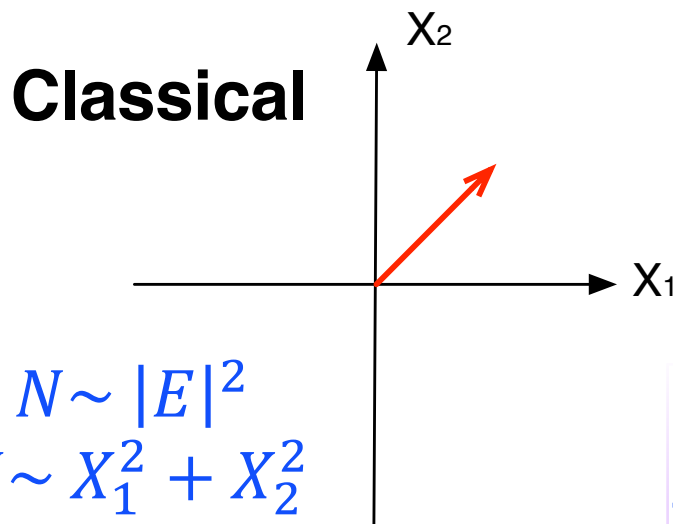
$$h_{Quantum} = \sqrt{\frac{4\hbar}{m\Omega^2 L^2}} \sqrt{\frac{1}{2} \left(\overset{\text{RPN}}{\underset{\downarrow}{\text{K}}} + \overset{\text{SN}}{\underset{\swarrow}{\frac{1}{\text{K}}}} \right)}, \quad \text{K} = \frac{4P\omega_0}{c^2 m\Omega^2}$$

- » Make the interferometer longer
- » Heavier test masses & more optical power
- » Inject of squeezed states of light

Ball and Stick Picture of Quantum Optical Noise

- Quantized Electromagnetic Fields
Quadrature Field Amplitudes

$$\hat{E} = \hat{X}_1 \cos \omega t + i\hat{X}_2 \sin \omega t$$



$$N \sim (X_1 + \Delta X_1)^2 + (X_2 + \Delta X_2)^2$$

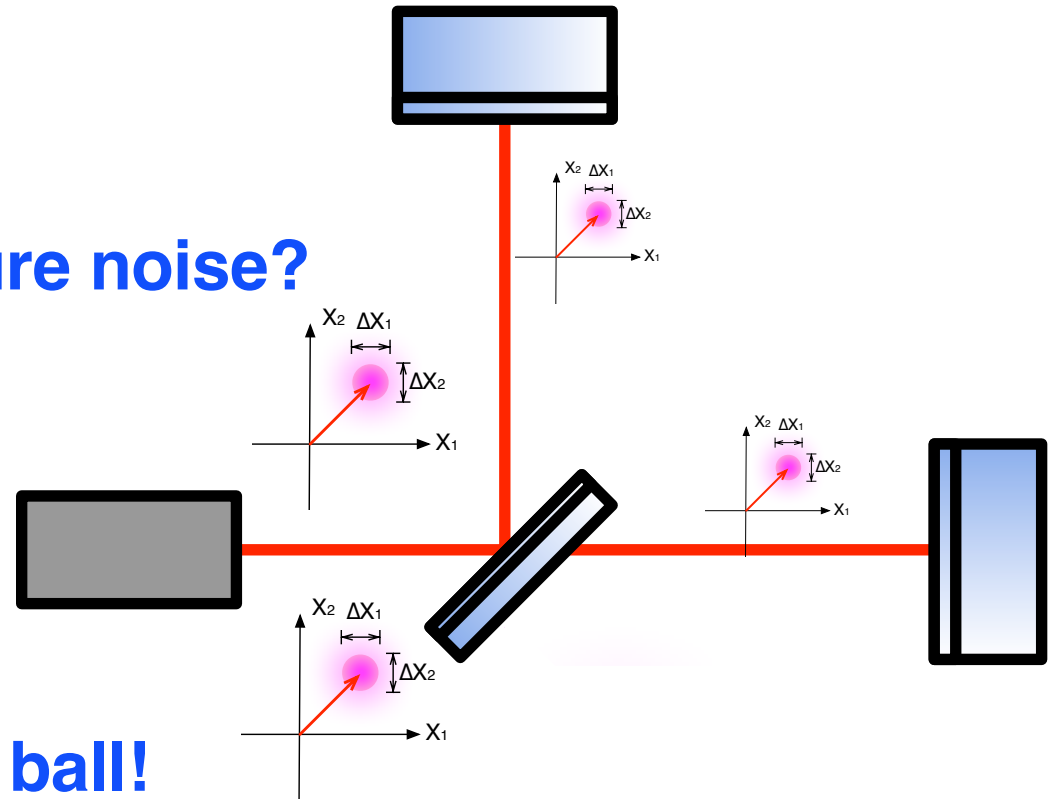
$$\langle N \rangle \sim X_1^2 + X_2^2 + \Delta X_1^2 + \Delta X_2^2$$

Heisenberg's uncertainty principle
 $\Delta X_1 \Delta X_2 \geq 1$

Equivalent to $\frac{1}{2}$ photon

How Does Quantum Noise Enter the Measurement?

- Beamsplitter acts the same on the stick and the ball!
-> no radiation pressure noise?
- Interferometer arms set to interfere the light back toward laser— both the stick and the ball!
-> no shot noise either?
- No such thing as Quantum Noise?

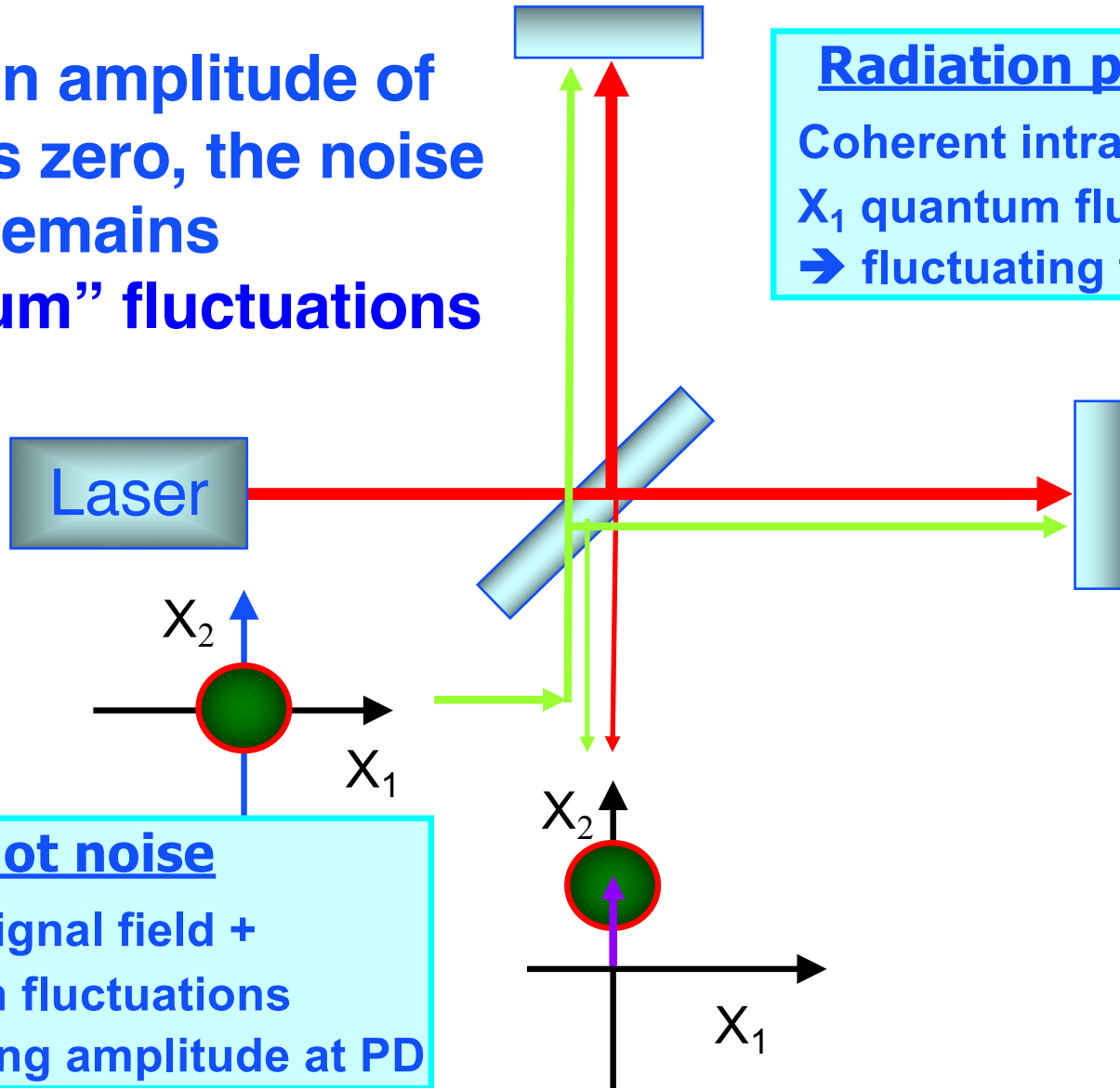


LIGO Quantum Noise in an Interferometer

Even when amplitude of the field is zero, the noise fuzz ball remains
=> “vacuum” fluctuations

Radiation pressure noise

Coherent intracavity field +
 X_1 quantum fluctuations
→ fluctuating force on mirrors



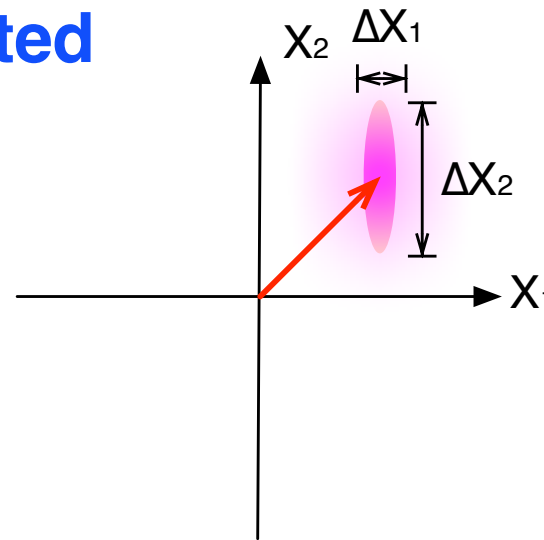
Shot noise

Coherent signal field +
 X_2 quantum fluctuations
→ fluctuating amplitude at PD

- No quantum constraints individually on ΔX_1 or ΔX_2 , only their product
- The noise can be redistributed while keeping the minimum uncertainty product

$$\Delta X_1 \Delta X_2 \geq 1$$

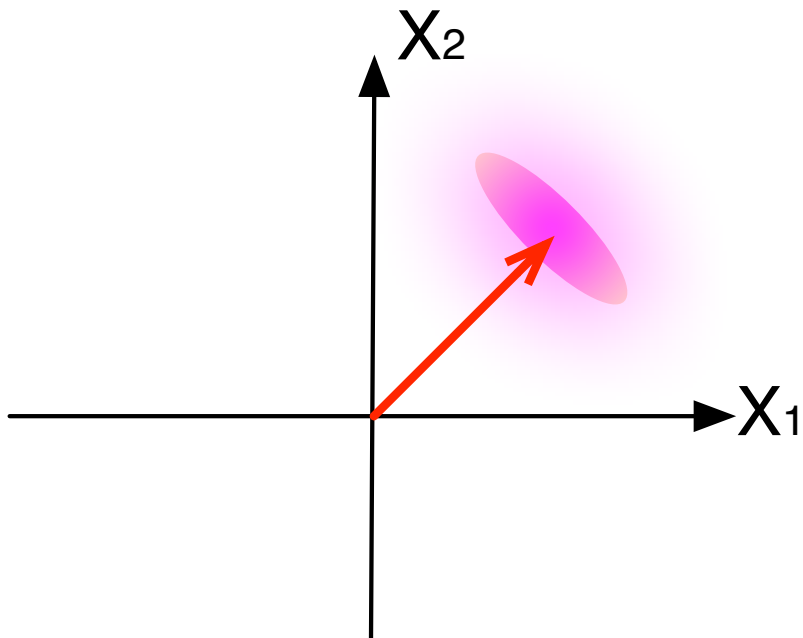
= Squeezed light



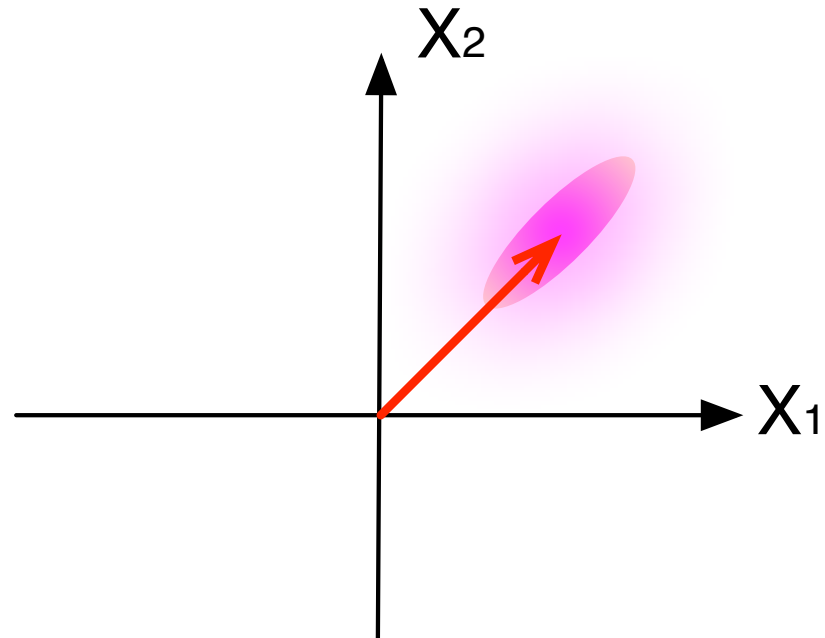
Squeezing

- Particularly useful two states

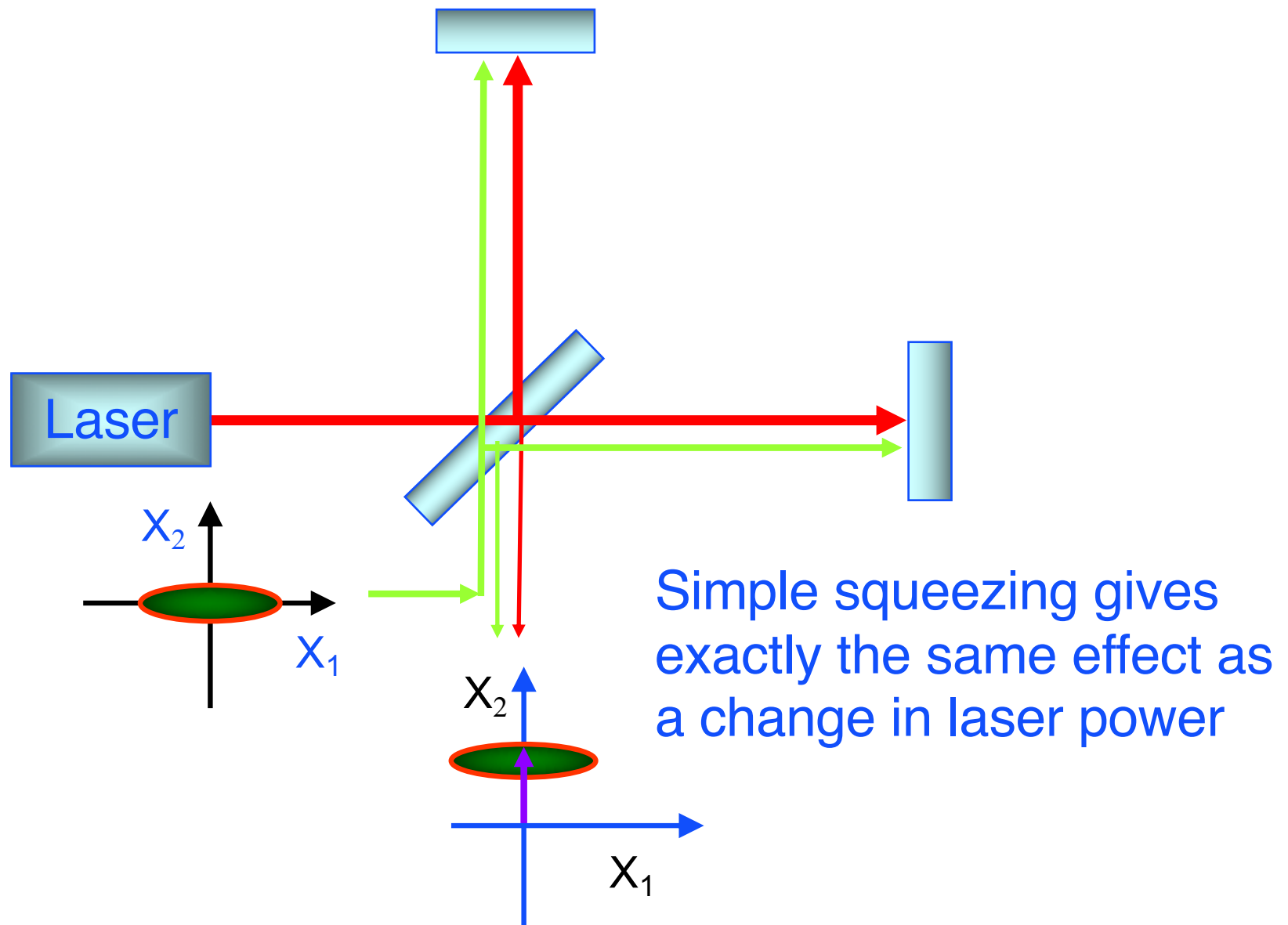
**Amplitude squeezing
(Phase anti-squeezing)**



**Phase squeezing
(Amplitude anti-squeezing)**

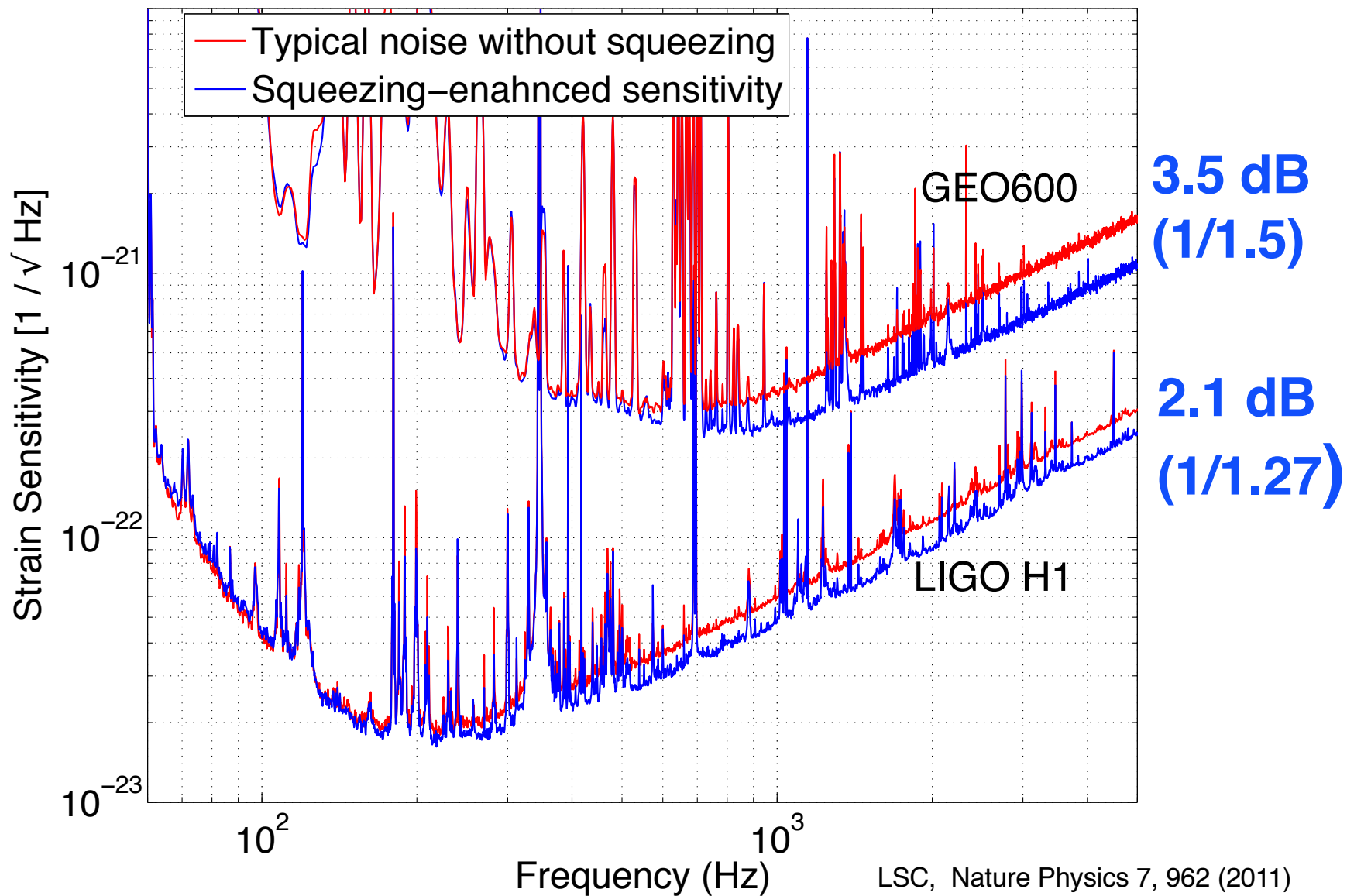


Squeezing in an Interferometer



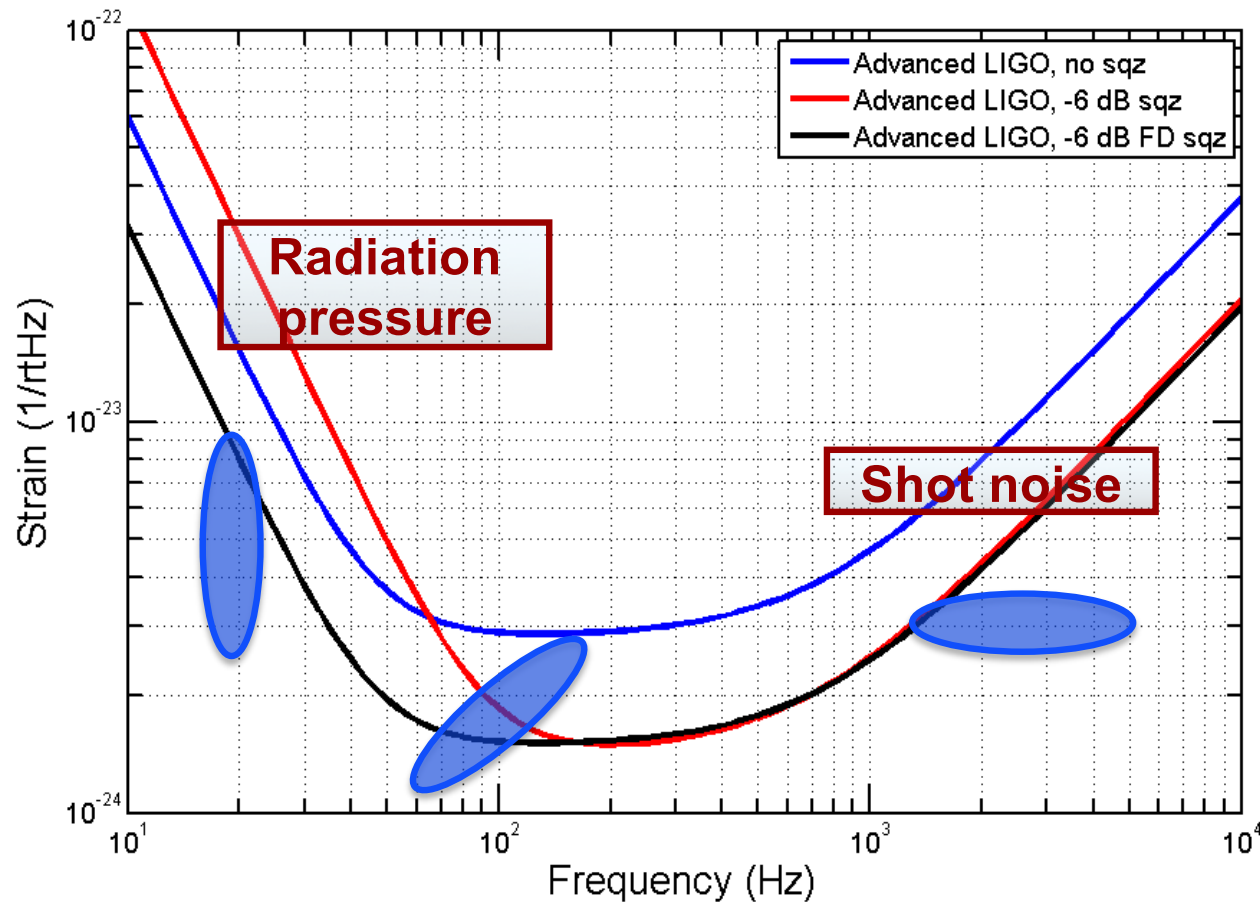


Does it Really Work? Squeezing in Action



Best of Both Worlds?

Rotate the squeezing quadrature as a function of frequency



Summing Up

- Detection of gravitational waves is already giving important new information about the Universe on its largest scales
- Made possible by advances in optics and precision interferometry
- Future advances will require us to confront the quantum nature of measurement